IMPERIAL

Constrained instantons in scalar field theory

KINGA GAWRYCH

BASED ON 2510.XXXXX WITH BEN ELDER AND ARTTU RAJANTIE

SCALARS 2025, UNIVERSITY OF WARSAW 24.09.2025

Vacuum decay – very important (electroweak vacuum stability, phase transitions)

Metastable vacuum (local minimum of V) → Stable vacuum (global minimum of V)

- Vacuum decay very important (electroweak vacuum stability, phase transitions)
 - Metastable vacuum (local minimum of V) → Stable vacuum (global minimum of V)
- Vacuum decay rate

Vacuum decay – very important (electroweak vacuum stability, phase transitions)

Metastable vacuum (local minimum of V) → Stable vacuum (global minimum of V)

■ Vacuum decay rate – standard computation method relies on <u>instantons</u>

saddle points of the **Euclidean action**

Vacuum decay – very important (electroweak vacuum stability, phase transitions)

Metastable vacuum (local minimum of V) → Stable vacuum (global minimum of V)

■ Vacuum decay rate — standard computation method relies on <u>instantons</u>

HOWEVER

saddle points of the **Euclidean action**

Vacuum decay – very important (electroweak vacuum stability, phase transitions)

Metastable vacuum (local minimum of V) → Stable vacuum (global minimum of V)

Vacuum decay rate – standard computation method relies on <u>instantons</u>

saddle points of the **Euclidean action**

HOWEVER

Sometimes we have a theory with a metastable vacuum but no instantons

Vacuum decay – very important (electroweak vacuum stability, phase transitions)

Metastable vacuum (local minimum of V) → Stable vacuum (global minimum of V)

Vacuum decay rate – standard computation method relies on <u>instantons</u>

saddle points of the Euclidean action

HOWEVER

Sometimes we have a theory with a metastable vacuum but no instantons

>>>> standard calculation method inapplicable

Vacuum decay – very important (electroweak vacuum stability, phase transitions)

Metastable vacuum (local minimum of V) → Stable vacuum (global minimum of V)

■ Vacuum decay rate – standard computation method relies on instantons

saddle points of the Euclidean action

HOWEVER

Sometimes we have a theory with a metastable vacuum but no instantons



IN THIS TALK: a method for computing the vacuum decay rate in theories with no instanton solutions

Vacuum decay rate
$$\Gamma = \frac{2}{\mathcal{V}} \operatorname{Im} \log \int \mathcal{D} \phi \, e^{-S_E[\phi]}$$

Vacuum decay rate
$$\Gamma = \frac{2}{\mathcal{V}} \operatorname{Im} \log \int \mathcal{D} \phi \, e^{-S_E[\phi]}$$

$$\int \mathcal{D}\phi \, e^{-S_E[\phi]} \approx Z[\phi^0] + \sum_{\hat{\phi}} Z[\hat{\phi}]$$

Vacuum decay rate
$$\Gamma = \frac{2}{\mathcal{V}} \operatorname{Im} \log \int \mathcal{D} \phi \, e^{-S_E[\phi]}$$

$$\int \mathcal{D}\phi \, e^{-S_E[\phi]} \approx Z[\phi^0] + \sum_{\hat{\phi}} Z[\hat{\phi}]$$

$$Z[\phi] = (\operatorname{Det} M[\phi])^{-1/2} e^{-S_E[\phi]}$$

Vacuum decay rate
$$\Gamma = \frac{2}{\mathcal{V}} \operatorname{Im} \log \int \mathcal{D} \phi \, e^{-S_E[\phi]}$$

$$\int \mathcal{D}\phi \, e^{-S_E[\phi]} \approx Z[\phi^0] + \sum_{\hat{\phi}} Z[\hat{\phi}]$$
 Non-trivial saddle points
$$Z[\phi] = (\mathrm{Det} M[\phi])^{-1/2} \, e^{-S_E[\phi]}$$

Vacuum decay rate
$$\Gamma = \frac{2}{\mathcal{V}} \operatorname{Im} \log \int \mathcal{D} \phi \, e^{-S_E[\phi]}$$

$$\int \mathcal{D}\phi \, e^{-S_E[\phi]} \approx Z[\phi^0] + \sum_{\hat{\phi}} Z[\hat{\phi}]$$
 Non-trivial saddie points
$$Z[\phi] = (\mathrm{Det} M[\phi])^{-1/2} \, e^{-S_E[\phi]}$$

Base idea due to Affleck [Affleck, '81] – perturbative

- Base idea due to Affleck [Affleck, '81] perturbative
- Our work non-perturbative

- Base idea due to Affleck [Affleck, '81] perturbative
- Our work non-perturbative

PROBLEM

No saddle points of the action in the space of all functions

- Base idea due to Affleck [Affleck, '81] perturbative
- Our work non-perturbative

PROBLEM

No saddle points of the action in the space of all functions



SOLUTION

Split the path integral into sectors, such that in a given sector a saddle point exists + sum over the sectors

$$\int \mathcal{D}\phi \, e^{-S_E[\phi]}$$

$$\int \mathcal{D}\phi \, e^{-S_E[\phi]} = \int d\bar{\xi} \int \mathcal{D}\phi \, \delta \left(\xi[\phi] - \bar{\xi}\right) \, e^{-S_E[\phi]}$$

$$\int \mathcal{D}\phi \, e^{-S_E[\phi]} = \int d\bar{\xi} \int \mathcal{D}\phi \, \delta \left(\xi[\phi] - \bar{\xi}\right) \, e^{-S_E[\phi]} \qquad \xi[\phi] = \int d^4x \, \mathcal{O}(\phi)$$

$$\int \mathcal{D}\phi \, e^{-S_E[\phi]} = \int d\bar{\xi} \int \mathcal{D}\phi \, \delta \left(\xi[\phi] - \bar{\xi}\right) \, e^{-S_E[\phi]}$$

$$\int \mathcal{D}\phi \, \delta \left(\xi[\phi] - \bar{\xi}\right) e^{-S_E[\phi]}$$

$$\int \mathcal{D}\phi \, e^{-S_E[\phi]} = \int d\bar{\xi} \int \mathcal{D}\phi \, \delta \left(\xi[\phi] - \bar{\xi}\right) \, e^{-S_E[\phi]}$$

$$\int \mathcal{D}\phi \, \delta \left(\xi[\phi] - \bar{\xi}\right) \, e^{-S_E[\phi]} \approx Z_{\bar{\xi}}[\phi^0] + \sum_{\hat{\phi}} Z_{\bar{\xi}}[\hat{\phi}]$$

$$\int \mathcal{D}\phi \, e^{-S_E[\phi]} = \int d\bar{\xi} \int \mathcal{D}\phi \, \delta \left(\xi[\phi] - \bar{\xi}\right) \, e^{-S_E[\phi]}$$

$$\int \mathcal{D}\phi \, \delta \left(\xi[\phi] - \bar{\xi}\right) \, e^{-S_E[\phi]} \approx Z_{\bar{\xi}}[\phi^0] + \sum_{\hat{\phi}} Z_{\bar{\xi}}[\hat{\phi}]$$
Non-trivial solutions subject to the constraint

$$\int \mathcal{D}\phi \, e^{-S_E[\phi]} = \int d\bar{\xi} \int \mathcal{D}\phi \, \delta \left(\xi[\phi] - \bar{\xi}\right) \, e^{-S_E[\phi]}$$

$$\int \mathcal{D}\phi \, \delta \left(\xi[\phi] - \bar{\xi}\right) e^{-S_E[\phi]} \approx Z_{\bar{\xi}}[\phi^0] + \sum_{\hat{\phi}} Z_{\bar{\xi}}[\hat{\phi}]$$

$$\sum_{\hat{\phi}} \sum_{\substack{\text{Non-trivial solutions subject to the constraint}}} \sum_{\substack{\text{Non-trivial solutions subject to the constraint}}} Z_{\bar{\xi}}[\phi] = \left(\operatorname{Det}_{\bar{\xi}} M_{\bar{\xi}}[\phi]\right)^{-1/2} e^{-S_E[\phi]}$$

$$\int \mathcal{D}\phi \, e^{-S_E[\phi]} = \int d\bar{\xi} \int \mathcal{D}\phi \, \delta \left(\xi[\phi] - \bar{\xi}\right) \, e^{-S_E[\phi]}$$

$$\int \mathcal{D}\phi \, \delta \left(\xi[\phi] - \bar{\xi}\right) e^{-S_E[\phi]} \approx Z_{\bar{\xi}}[\phi^0] + \sum_{\hat{\phi}} Z_{\bar{\xi}}[\hat{\phi}]$$

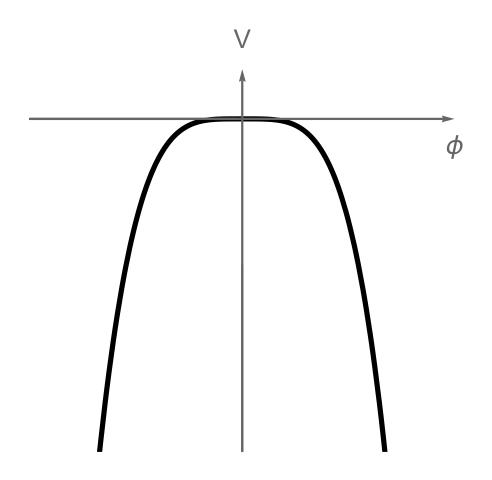
$$\text{Non-trivial solutions subject to the constraint}$$

$$\frac{\text{CAUTION!}}{\text{Det}_{\bar{\xi}}\text{"}} \neq \text{"Det} \quad M_{\bar{\xi}}[\phi] \neq M[\phi]$$

$$Z_{\bar{\xi}}[\phi] = \left(\text{Det}_{\bar{\xi}} M_{\bar{\xi}}[\phi]\right)^{-1/2} e^{-S_E[\phi]}$$

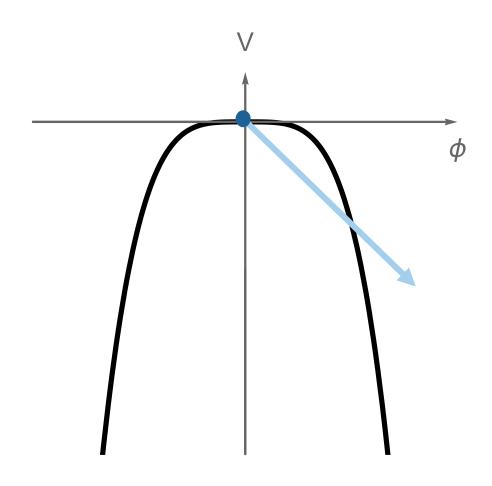
$$S_E[\phi] = \int d^4x \left(\frac{1}{2} (\partial_\mu \phi)^2 + \frac{1}{2} m^2 \phi^2 - \frac{\lambda}{4!} \phi^4 \right)$$

$$S_E[\phi] = \int d^4x \left(\frac{1}{2} (\partial_\mu \phi)^2 + \frac{1}{2} m^2 \phi^2 - \frac{\lambda}{4!} \phi^4 \right)$$



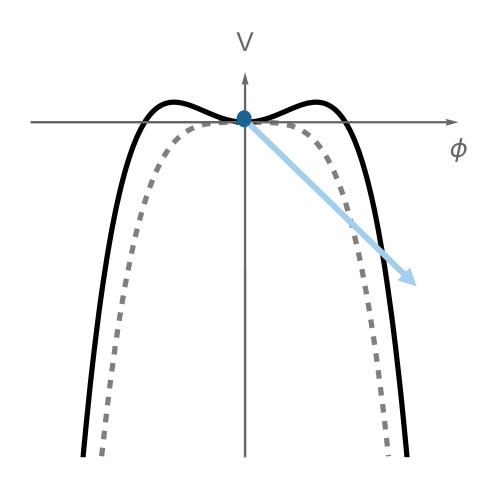
$$S_{E}[\phi] = \int d^{4}x \left(\frac{1}{2} (\partial_{\mu}\phi)^{2} + \frac{1}{2} m^{2} \phi^{2} - \frac{\lambda}{4!} \phi^{4} \right)$$

$$m^{2} = 0$$



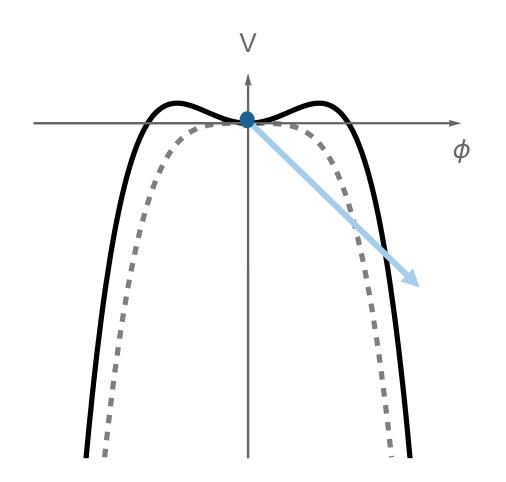
$$S_{E}[\phi] = \int d^{4}x \left(\frac{1}{2} (\partial_{\mu}\phi)^{2} + \frac{1}{2} m^{2} \phi^{2} - \frac{\lambda}{4!} \phi^{4} \right)$$

$$m^{2} = 0$$



$$S_{E}[\phi] = \int d^{4}x \left(\frac{1}{2} (\partial_{\mu}\phi)^{2} + \frac{1}{2} m^{2} \phi^{2} - \frac{\lambda}{4!} \phi^{4} \right)$$

$$m^{2} > 0$$

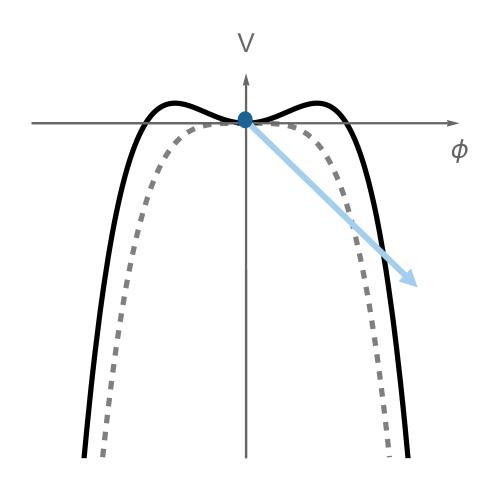


$$S_{E}[\phi] = \int d^{4}x \left(\frac{1}{2} (\partial_{\mu}\phi)^{2} + \frac{1}{2} m^{2} \phi^{2} - \frac{\lambda}{4!} \phi^{4} \right)$$

$$m^{2} > 0$$

$$\phi_{inst}(x) = ?$$
 $S_E[\phi_{inst}] = ?$

Example $--\lambda \phi^4$

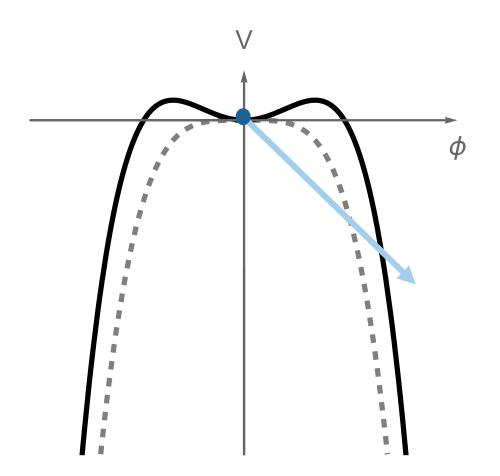


$$S_{E}[\phi] = \int d^{4}x \left(\frac{1}{2} (\partial_{\mu}\phi)^{2} + \frac{1}{2} m^{2} \phi^{2} - \frac{\lambda}{4!} \phi^{4} \right)$$

$$m^{2} > 0$$

$$\phi_{inst}(x) = 2$$
 $S_{E}[\phi_{inst}] = 2$

Example $--\lambda \phi^4$



$$S_{E}[\phi] = \int d^{4}x \left(\frac{1}{2} (\partial_{\mu}\phi)^{2} + \frac{1}{2} m^{2} \phi^{2} - \frac{\lambda}{4!} \phi^{4} \right)$$

$$m^{2} > 0$$

$$\phi_{inst}(x) = 2$$
 $S_{E}[\phi_{inst}] = 2$

CULPRIT

$$\phi(x) \to a\phi(ax)$$

AIM: find the saddle points of the constrained path integral and compute their action

AIM: find the saddle points of the constrained path integral and compute their action

■ In practice — use Lagrange multipliers to find the constrained instantons

AIM: find the saddle points of the constrained path integral and compute their action

■ In practice — use Lagrange multipliers to find the constrained instantons

Simplification:

```
Theory has O(4) symmetry \longrightarrow <u>assume</u> the solution has the same symmetry 4 PDEs \longrightarrow 1 ODE
```

AIM: find the saddle points of the constrained path integral and compute their action

■ In practice — use Lagrange multipliers to find the constrained instantons

Simplification:

```
Theory has O(4) symmetry \longrightarrow <u>assume</u> the solution has the same symmetry 4 PDEs \longrightarrow 1 ODE
```

Find solutions numerically:

BVP --> shooting in a finite interval

AIM: find the saddle points of the constrained path integral and compute their action

■ In practice — use Lagrange multipliers to find the constrained instantons

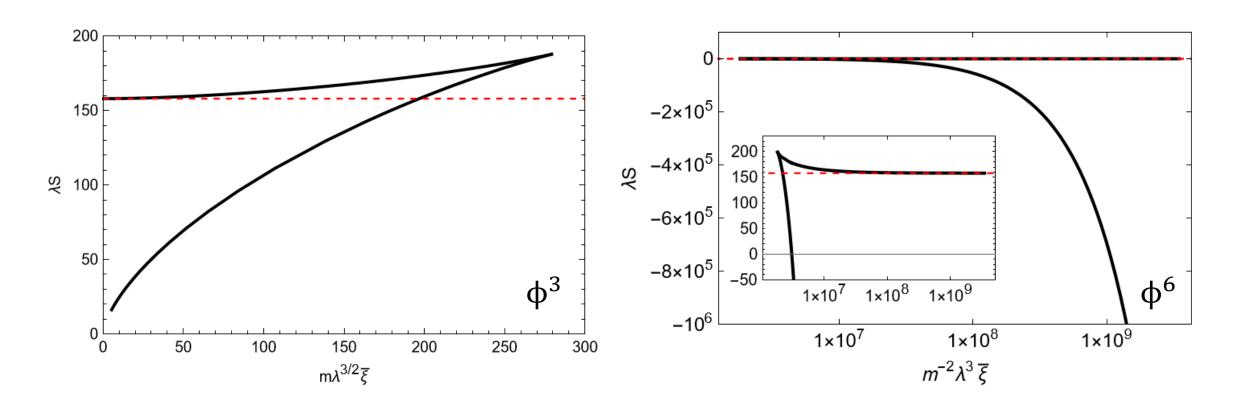
Simplification:

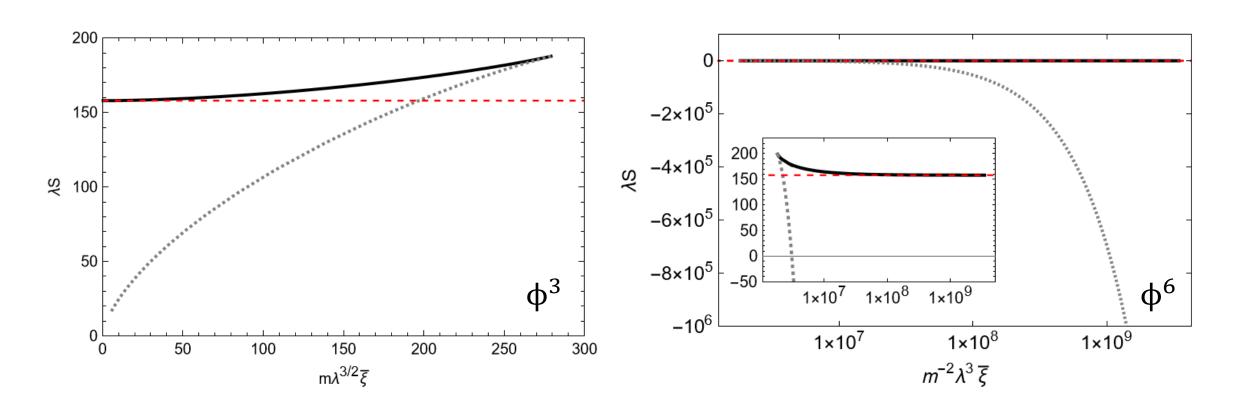
Theory has
$$O(4)$$
 symmetry \longrightarrow assume the solution has the same symmetry 4 PDEs \longrightarrow 1 ODE

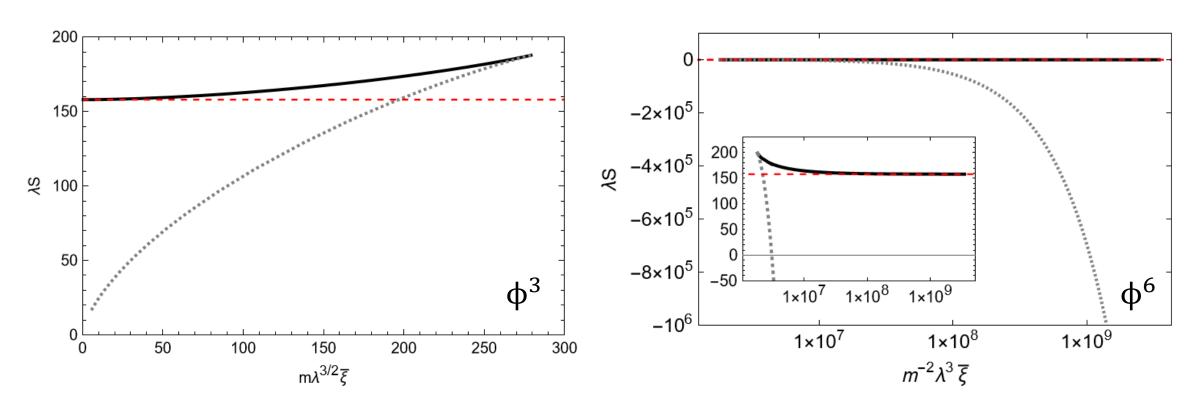
Find solutions numerically:

BVP --> shooting in a finite interval

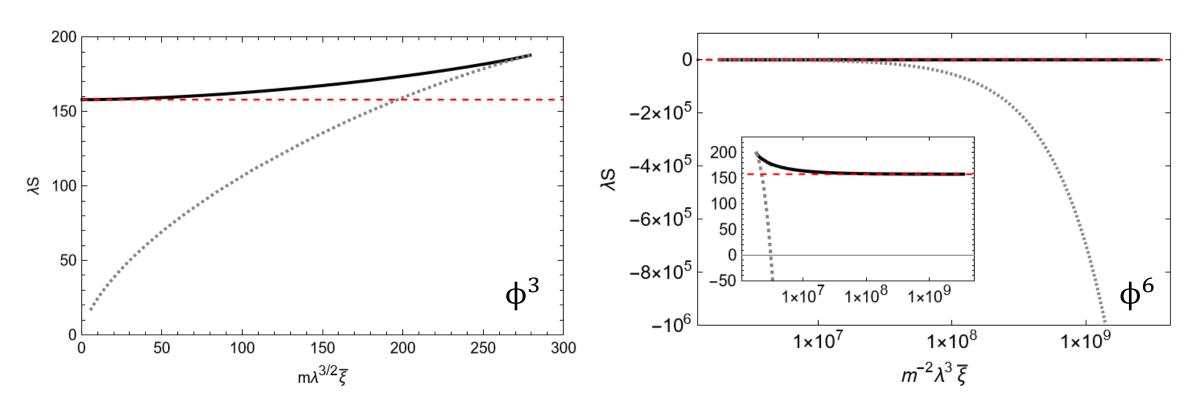
• Choice of the constraint operator $\mathcal{O}(\varphi)$: two different constraints - φ^3 and φ^6



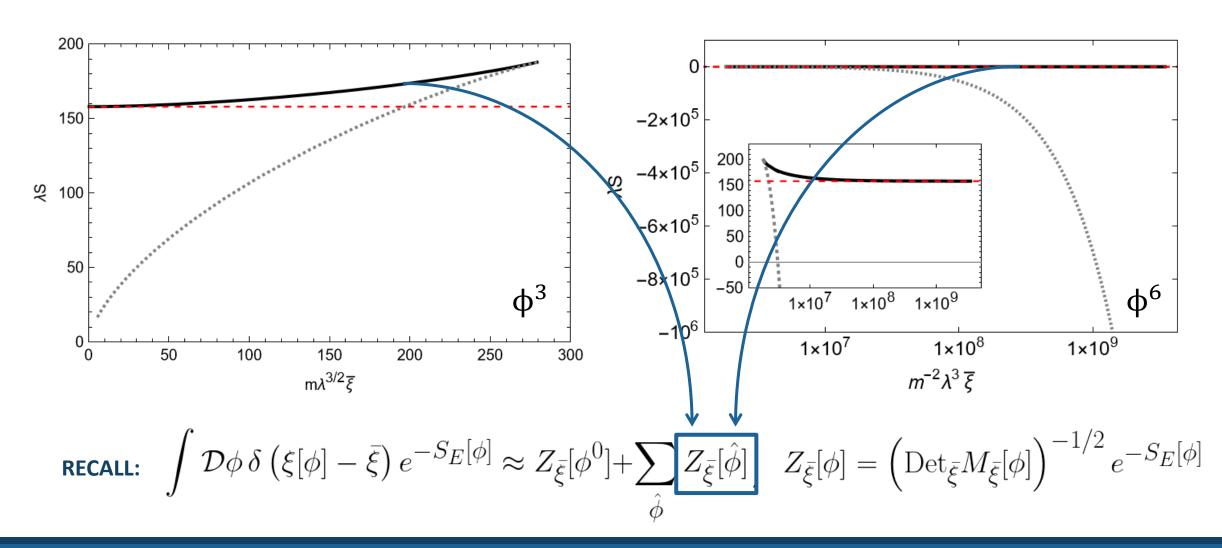


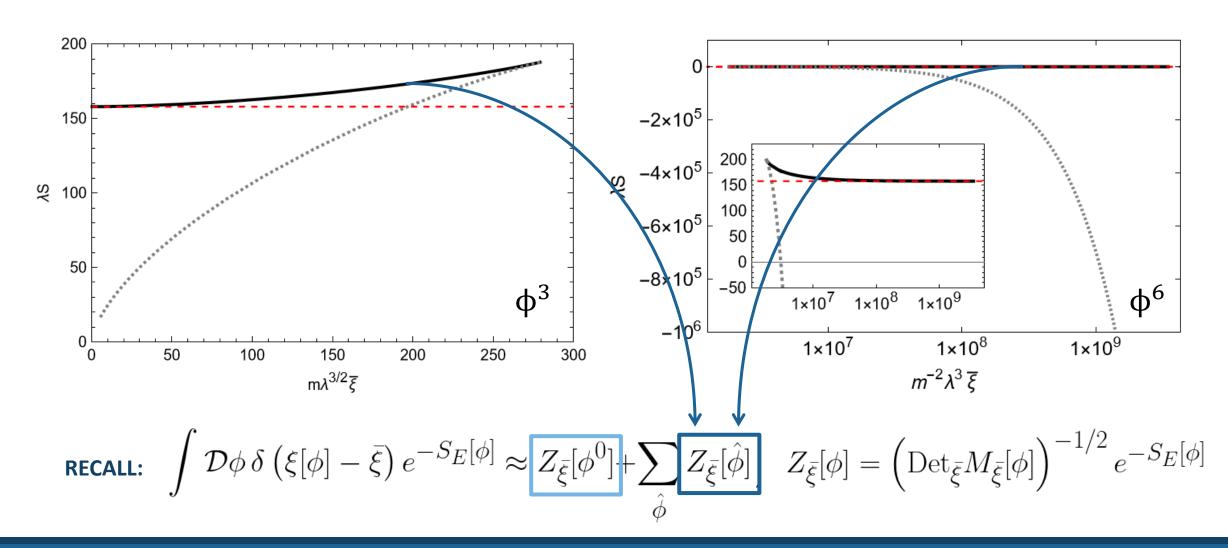


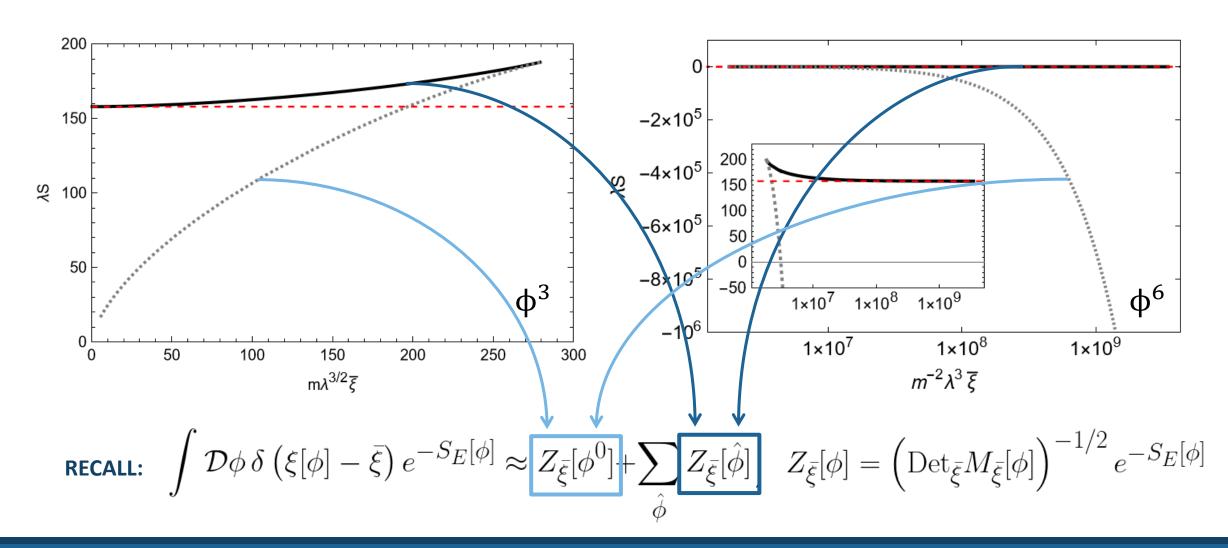
$$\text{RECALL:} \quad \int \mathcal{D}\phi \, \delta \left(\xi[\phi] - \bar{\xi} \right) e^{-S_E[\phi]} \approx Z_{\bar{\xi}}[\phi^0] + \sum_{\hat{\phi}} Z_{\bar{\xi}}[\hat{\phi}] \, , \quad Z_{\bar{\xi}}[\phi] = \left(\mathrm{Det}_{\bar{\xi}} M_{\bar{\xi}}[\phi] \right)^{-1/2} e^{-S_E[\phi]} \, .$$



$$\text{RECALL:} \quad \int \mathcal{D}\phi \, \delta \left(\xi[\phi] - \bar{\xi} \right) e^{-S_E[\phi]} \approx Z_{\bar{\xi}}[\phi^0] + \sum_{\hat{\phi}} \boxed{Z_{\bar{\xi}}[\hat{\phi}]} \quad Z_{\bar{\xi}}[\phi] = \left(\mathrm{Det}_{\bar{\xi}} M_{\bar{\xi}}[\phi] \right)^{-1/2} e^{-S_E[\phi]}$$

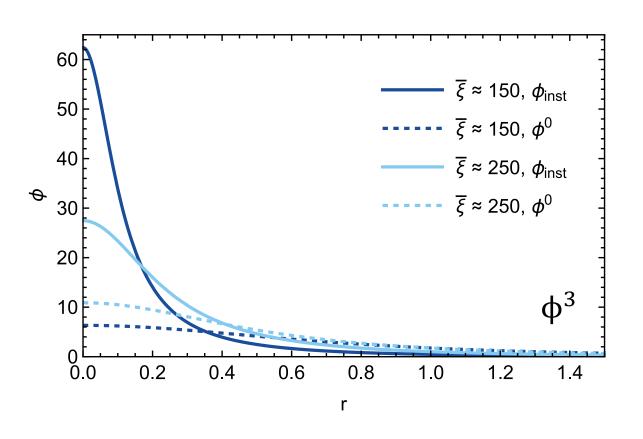


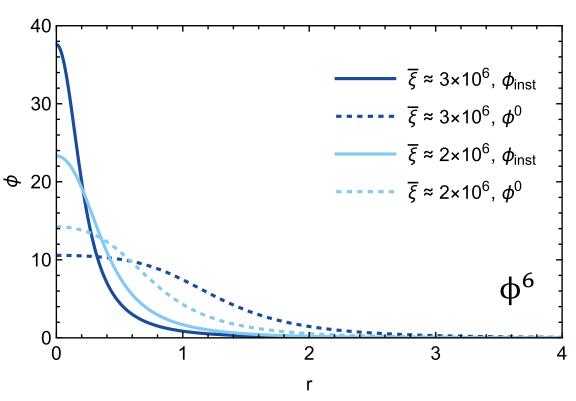




Results – field profiles

Results – field profiles





Summary and further research

SUMMARY

- Introduced a non-perturbative constrained instanton method for computing vacuum decay rates in theories with no explicit instanton solution
- Showed you an application of this method to a particular scalar field theory
- Omitted a lot of technical details these can be found in the paper (out soon)

or just ask me!

FURTHER RESEARCH

- Computation of the functional determinant (work in progress)
- Comparison of the results (full rate) for different constraint operators
- Analytic construction of constrained instantons for extreme values of $\bar{\xi}$ (work in progress)
- Extension of the method to theories with gauge fields (partial results) and application to the electroweak sector of the Standard Model

Thank you

Backup slides

Lagrange multipliers

In practice? → Lagrange multipliers

Saddle point of
$$S_E[\phi]$$
 for some $\xi[\phi] = \overline{\xi}$

- - 1

- Note different than the standard textbook approach to Lagrange multipliers
- Look for solutions for fixed κ (easy) instead of fixed ξ (hard) \longrightarrow Solve EOM following from \tilde{S}_{κ}

MODIFIED ACTION

$$\tilde{S}_{\kappa}[\phi] = S_{E}[\phi] + \underset{\kappa}{\kappa} \xi[\phi]$$
 Constant Lagrange multiplier

Equation of motion

$$\tilde{S}_{\kappa}[\Phi] = \frac{1}{\lambda} \int d^4X \left(\frac{1}{2} (\partial_{\mu} \Phi)^2 + \frac{1}{2} \Phi^2 - \frac{1}{4!} \Phi^4 + \kappa_n \Phi^n \right)$$

$$n = 3, 6$$

+
$$O(4)$$
 symmetry
$$\Phi'' + \frac{3}{R}\Phi' - \Phi + \frac{1}{6}\Phi^3 + n\kappa_n\Phi^{n-1} = 0$$

 $R \equiv mr$

$$\Phi \equiv \frac{\lambda^{1/2}}{m} \phi$$

$$\kappa_n = \frac{m^{n-4}}{\lambda^{n/2-1}} \kappa$$

Boundary conditions: $\Phi'(R_{\min}) = 0$ and $\Phi(R_{\max}) = 0$



Results - constraints

