

NNLO QCD corrections to $\bar{B} \rightarrow X_s \gamma$ without interpolation in m_c

Mikołaj Misiak

University of Warsaw

“Scalars 2025”, Warsaw, September 22-25th, 2025

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M. Czaja, M. Czakon, T. Huber, MM, M. Niggetiedt, A. Rehman, K. Schönwald and M. Steinhauser
[to appear in arXiv:2510.nnnnn]

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3. The current SM prediction for $\mathcal{B}_{s\gamma}$ and updated bounds on 2HDM-II

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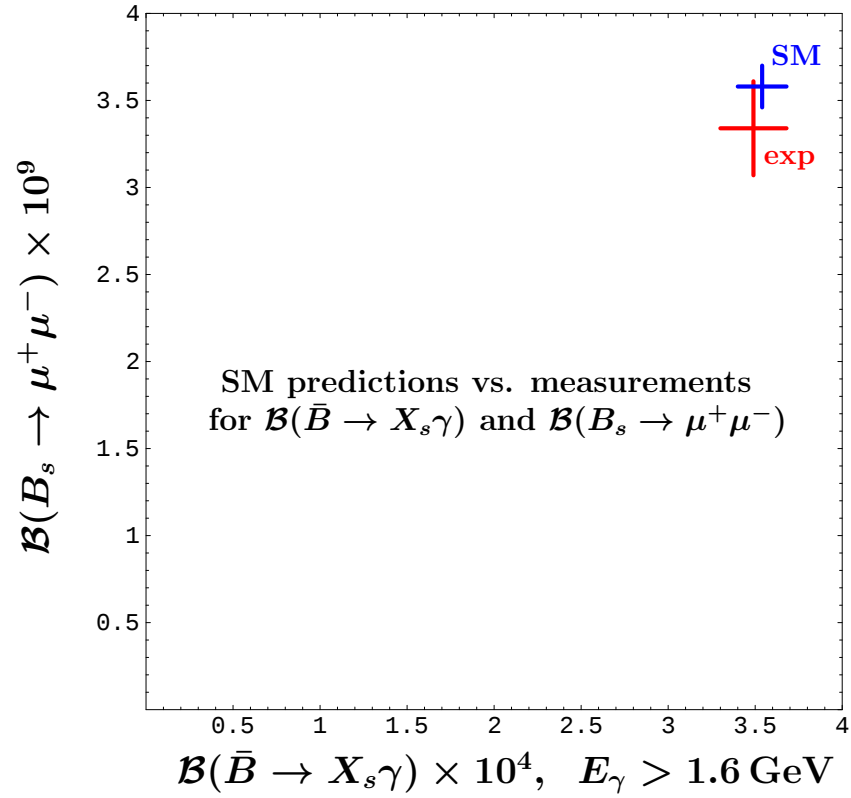
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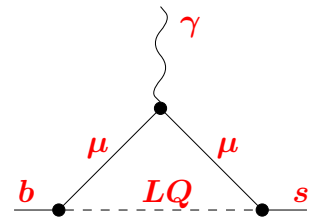
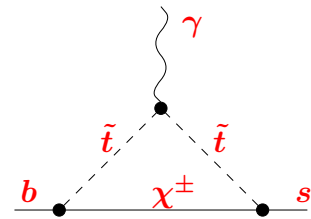
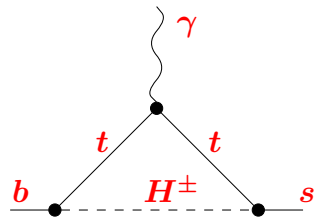
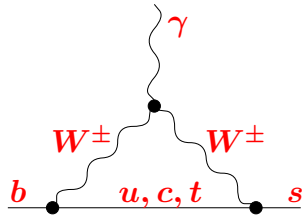
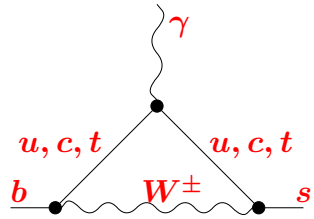
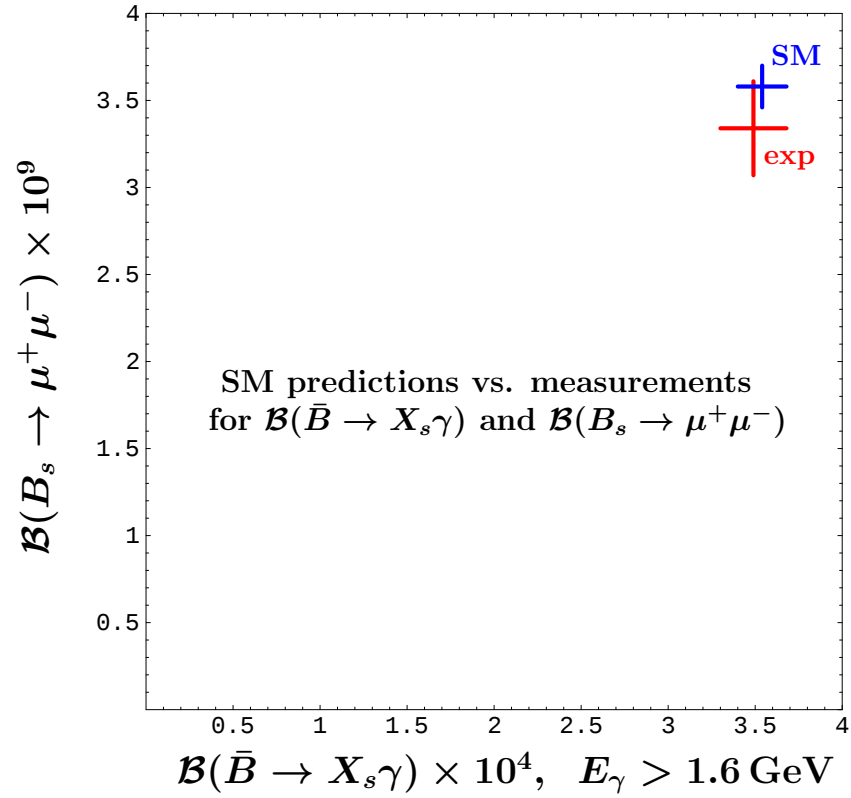
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5. Summary and outlook

Introductory conclusions

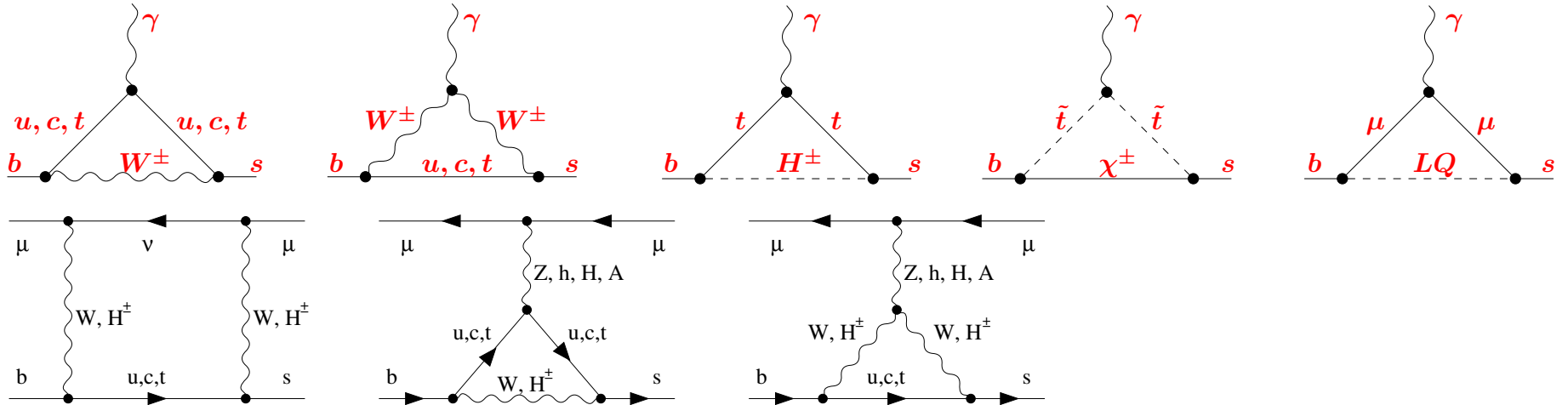
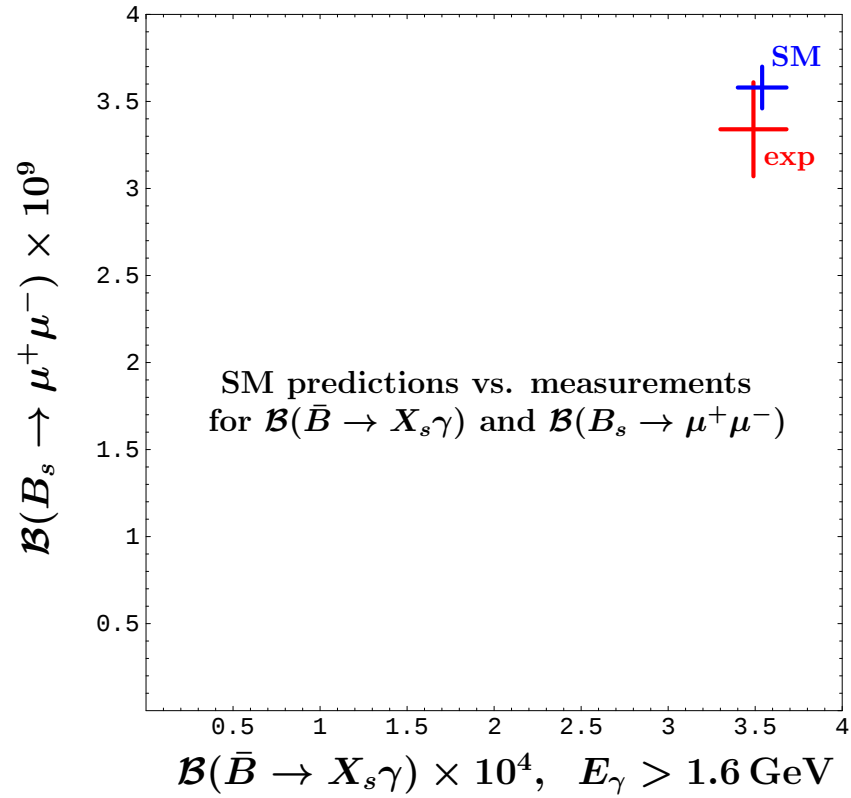
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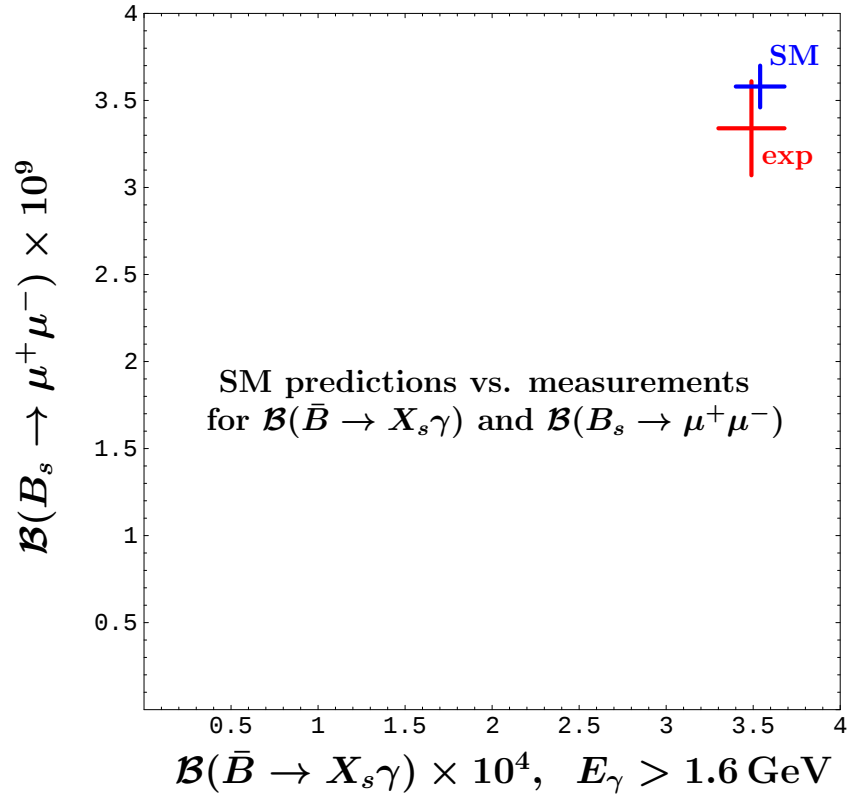
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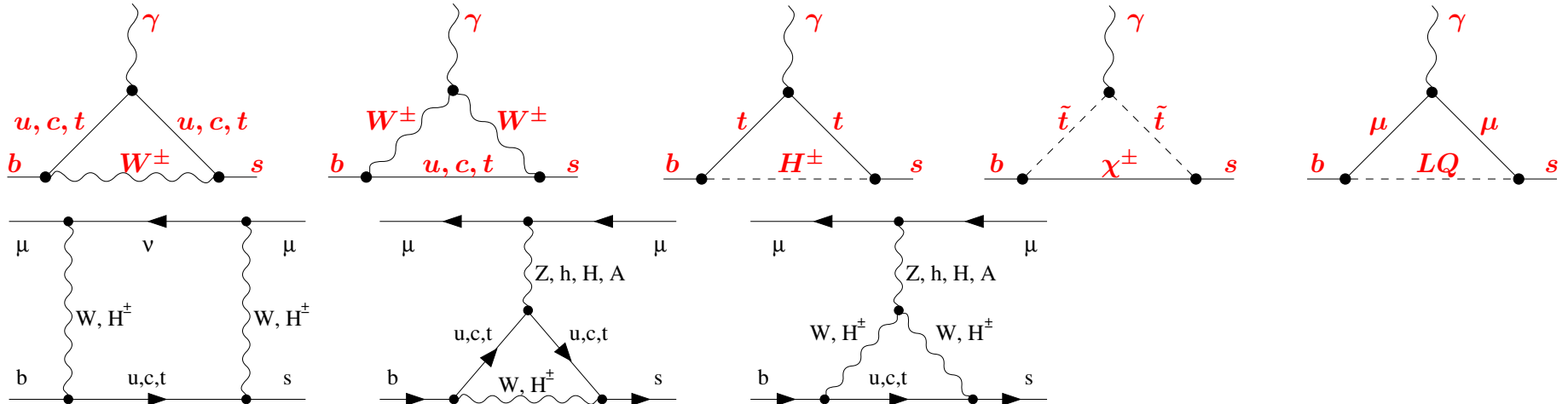


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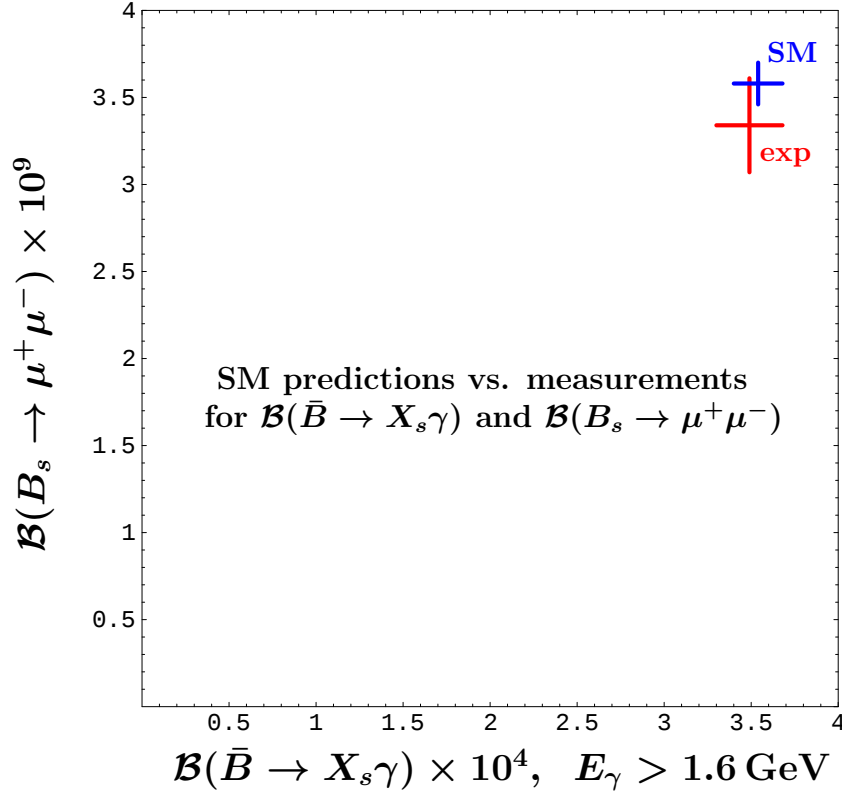


$$\mathcal{B}(\bar{B} \rightarrow X_s \gamma)_{E_\gamma > 1.6}^{\text{exp}} \times 10^4 = 3.49 \pm 0.19 \quad (\pm 5.5\%)$$

CLEO, BaBar and Belle measurements combined
by PDG [2022,2024] and HFLAV [arXiv:2206.07501].



Introductory conclusions

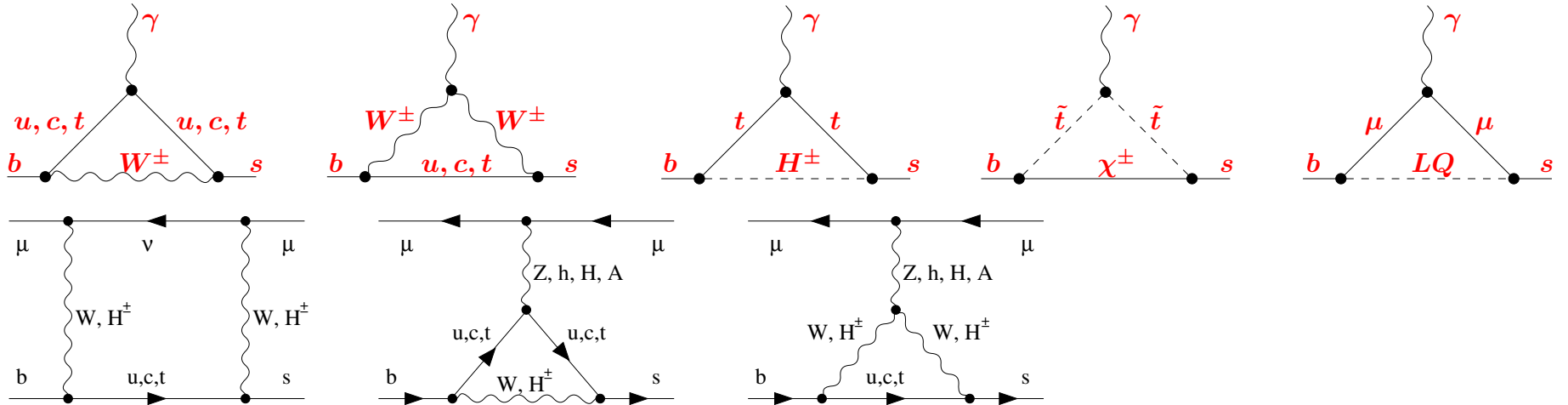


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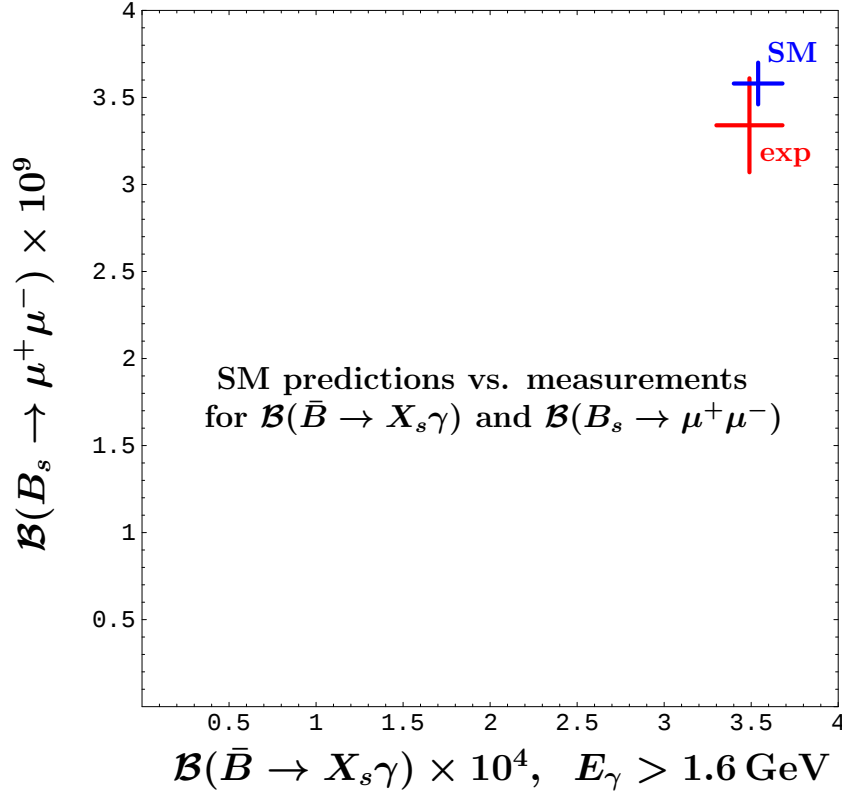
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$$\mathcal{B}(\bar{B} \rightarrow X_s \gamma)_{E_\gamma > 1.6}^{\text{SM}} \times 10^4 = 3.54 \pm 0.14 \quad (\pm 4.0\%)$$

Final for 2510.nnnnn.



Introductory conclusions



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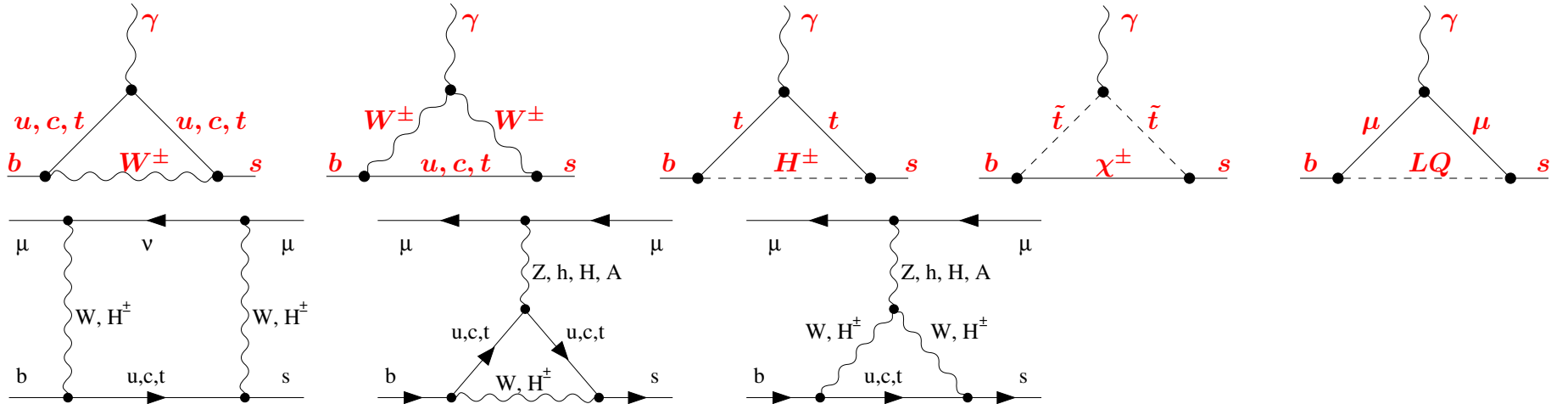
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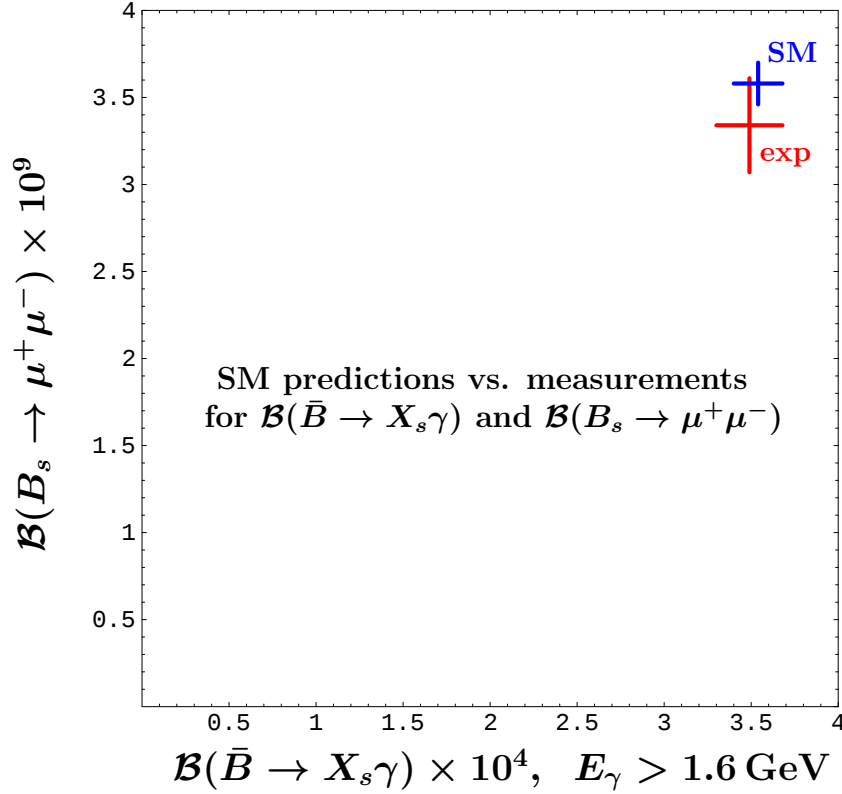
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$$\mathcal{B}(B_s \rightarrow \mu^+ \mu^-)^{\text{exp}} \times 10^9 = 3.34 \pm 0.27 \quad (\pm 8.1\%)$$

LHCb, CMS and ATLAS measurements combined by PDG [2024].



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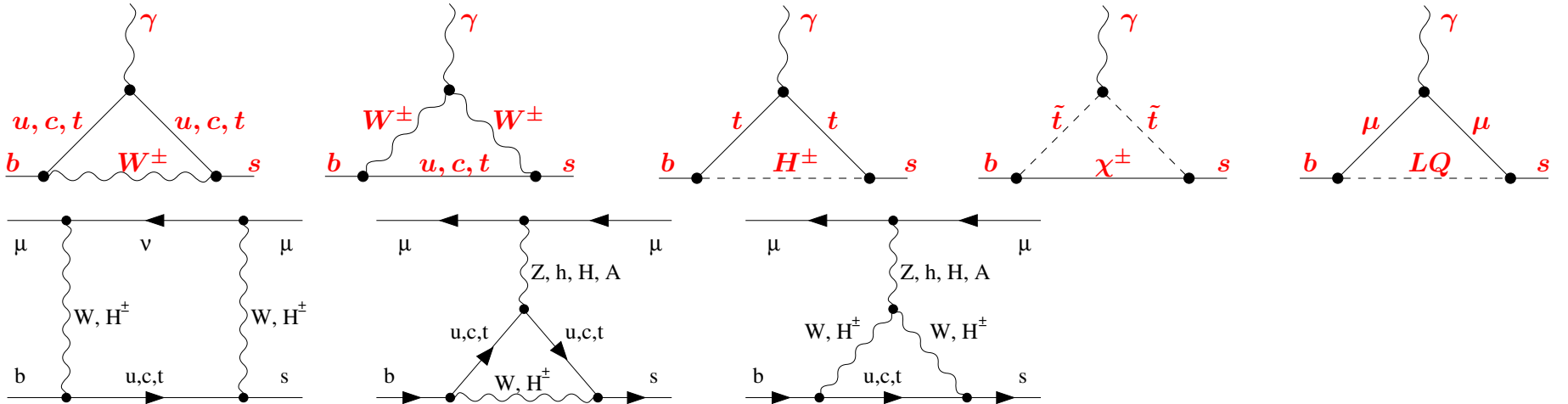
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$$\mathcal{B}(B_s \rightarrow \mu^+ \mu^-)^{\text{SM}} \times 10^9 = 3.58 \pm 0.12 \quad (\pm 3.4\%)$$

arXiv:1311.0903 by C. Bobeth, M. Gorbahn, T. Hermann, MM,
E. Stamou, M. Steinhauser [with parameter updates and](#) -0.5%
QED correction from arXiv:1907.07011 by M. Beneke, C. Bobeth
and R. Szafron.



$\mathcal{B}_{s\gamma}$ in the Two-Higgs-Doublet Model II

$\mathcal{B}_{s\gamma}$ in the Two-Higgs-Doublet Model II

$$\mathcal{L}_{\text{Yukawa}} = -\bar{q}_i Y_{ij}^u u_j H_2 - \bar{q}_i Y_{ij}^d d_j H_1 - \bar{l}_i Y_{ij}^e e_j H_1 + \text{h.c.},$$

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$$\langle H_2 \rangle = \frac{v \sin \beta}{\sqrt{2}} \begin{pmatrix} 0 \\ 1 \end{pmatrix}, \quad \langle H_1 \rangle = \frac{v \cos \beta}{\sqrt{2}} \begin{pmatrix} 1 \\ 0 \end{pmatrix}.$$

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$$\mathcal{L}_{H^\pm \bar{q} q} = \frac{g_{\text{ew}}}{M_W \sqrt{2}} \left[H^+ \bar{u}_i (m_{u_i} V_{ij} P_L \cot \beta + V_{ij} m_{d_j} P_R \tan \beta) d_j + \text{h.c.} \right].$$

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$$\delta_{H^\pm A}(b \rightarrow s\gamma) \quad \sim \quad F_1 \left(\frac{m_t^2}{M_{H^\pm}^2} \right) + F_2 \left(\frac{m_t^2}{M_{H^\pm}^2} \right) \cot^2 \beta$$

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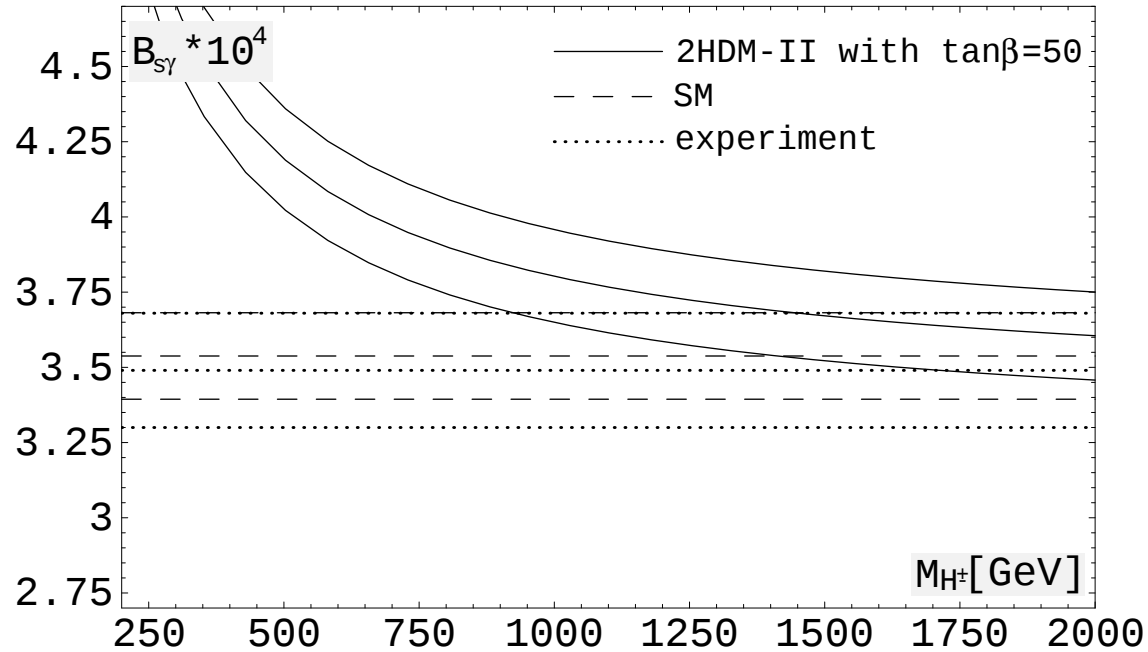
$$\delta_{H^\pm} A(b \rightarrow s\gamma) \quad \sim \quad \underset{\mathcal{B}_{s\gamma} \uparrow}{F_1 \left(\frac{m_t^2}{M_{H^\pm}^2} \right)} + \underset{\mathcal{B}_{s\gamma} \uparrow}{F_2 \left(\frac{m_t^2}{M_{H^\pm}^2} \right)} \cot^2 \beta$$

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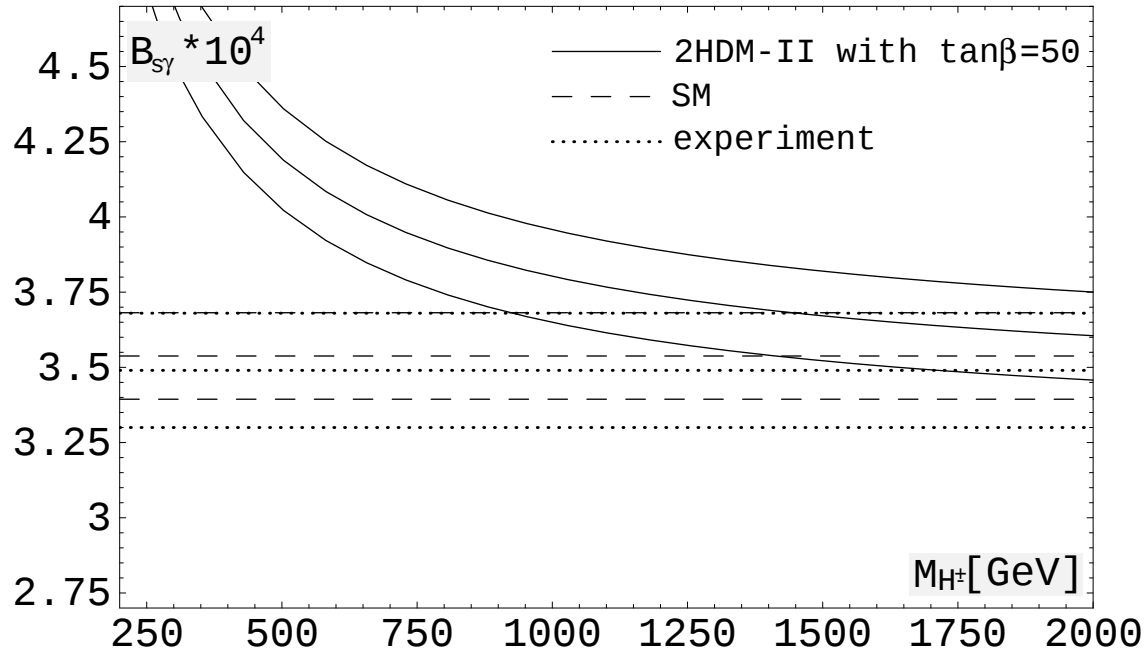


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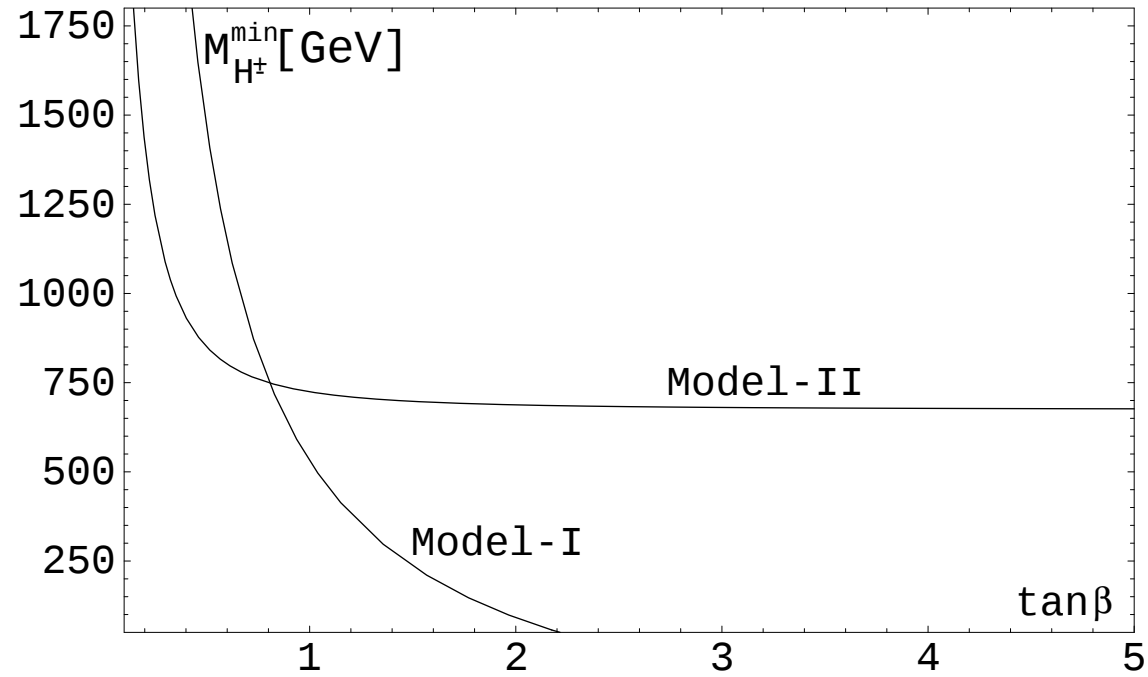
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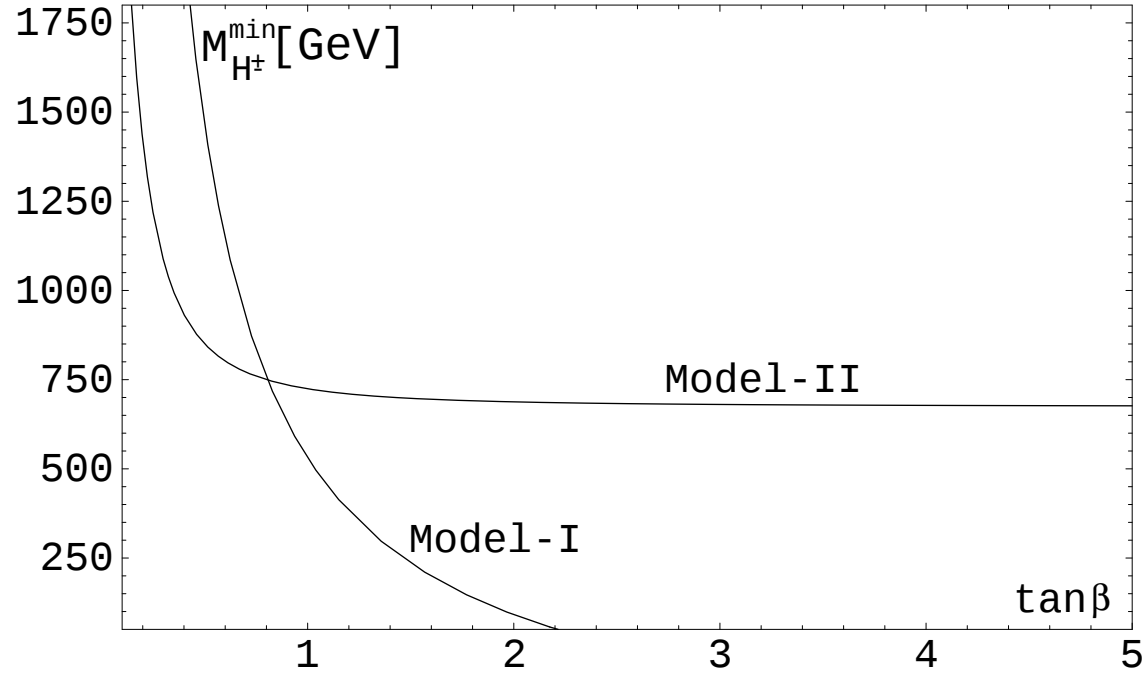


$$M_{H^\pm} > 675 \text{ GeV} \quad @ \text{ 95\% C.L.}$$

Exclusion bounds (95%C.L.) in the $\tan\beta$ – $M_{H^\pm}$ plane:



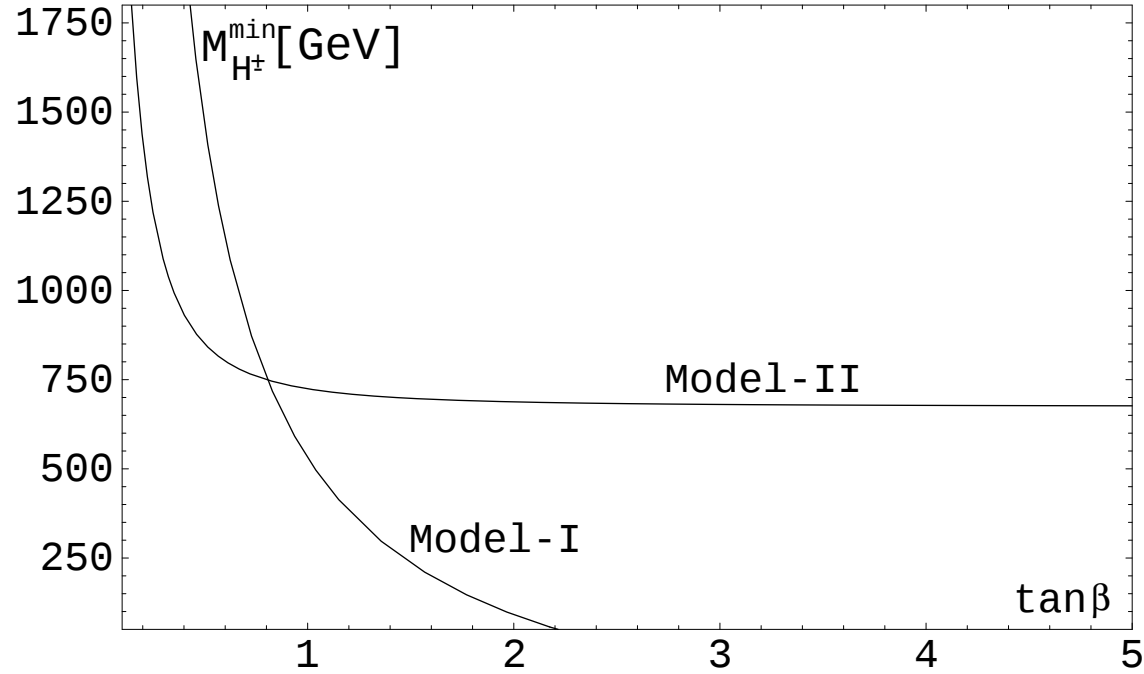
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In Model-I:

$$\delta_{H^\pm} A(b \rightarrow s\gamma) \sim \left[-F_1 \left(\frac{m_t^2}{M_{H^\pm}^2} \right) + F_2 \left(\frac{m_t^2}{M_{H^\pm}^2} \right) \right] \cot^2 \beta$$

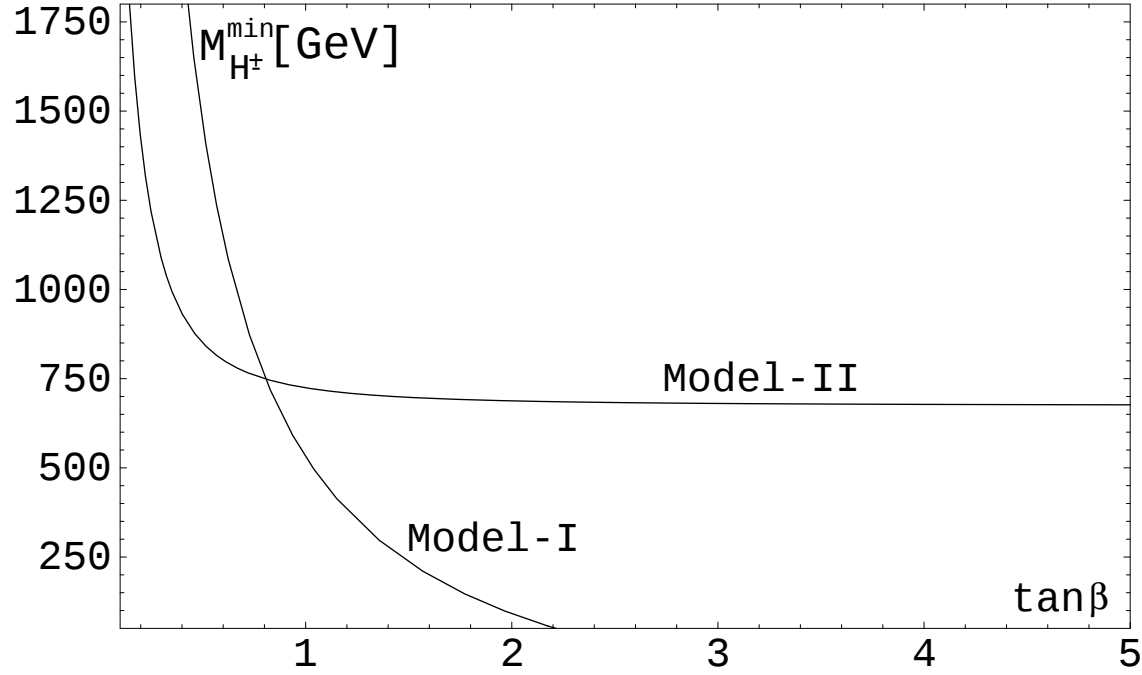
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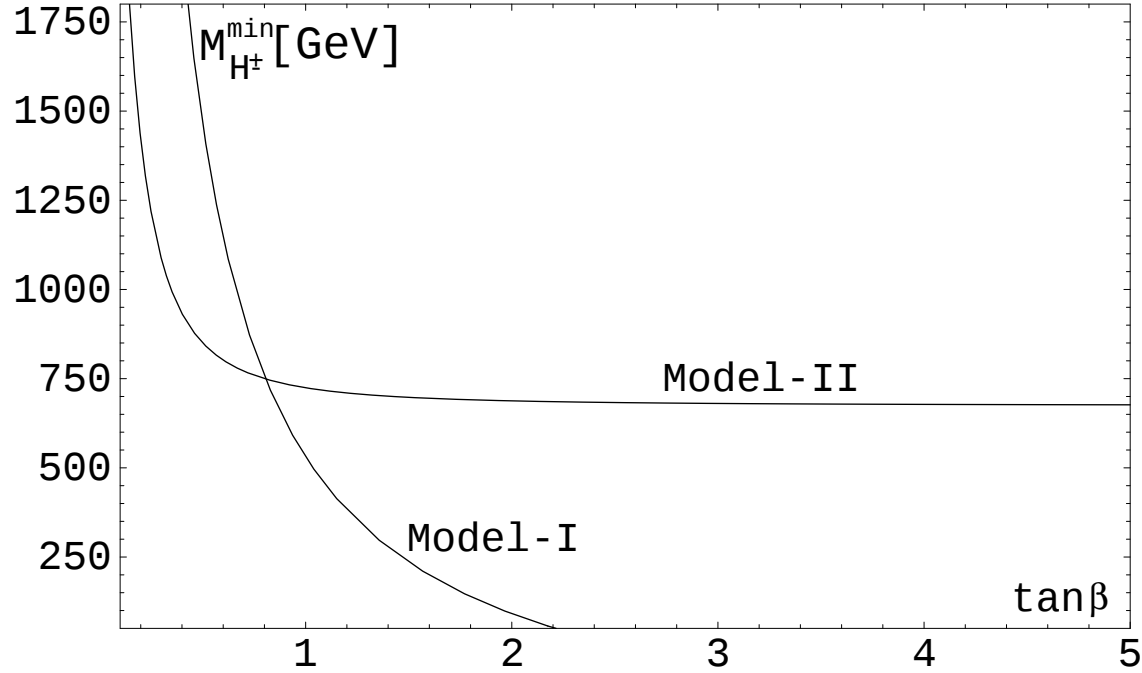


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In a large class of beyond-SM theories: $\mathcal{B}_{s\gamma} \times 10^4 = (3.54 \pm 0.14) - 8.51 \Delta_{\text{BSM}} C_7(\mu_0) - 2.16 \Delta_{\text{BSM}} C_8(\mu_0),$

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with $\Delta_{\text{BSM}} C_7(\mu_0)$ and $\Delta_{\text{BSM}} C_8(\mu_0)$ provided in hep-ph/9904413 up to the NLO in QCD.

Updated SM prediction:

$$\mathcal{B}_{s\gamma}^{\text{SM}} = (3.54 \pm 0.14) \times 10^{-4}$$

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$$\mathcal{B}_{s\gamma}^{\text{SM}} = (3.54 \pm 0.14) \times 10^{-4} \quad \left[\pm 4.0\% \simeq \pm \sqrt{(3.0\%)_{\text{h.o.}}^2 + (2.7\%)_{\text{param}}^2} \right].$$

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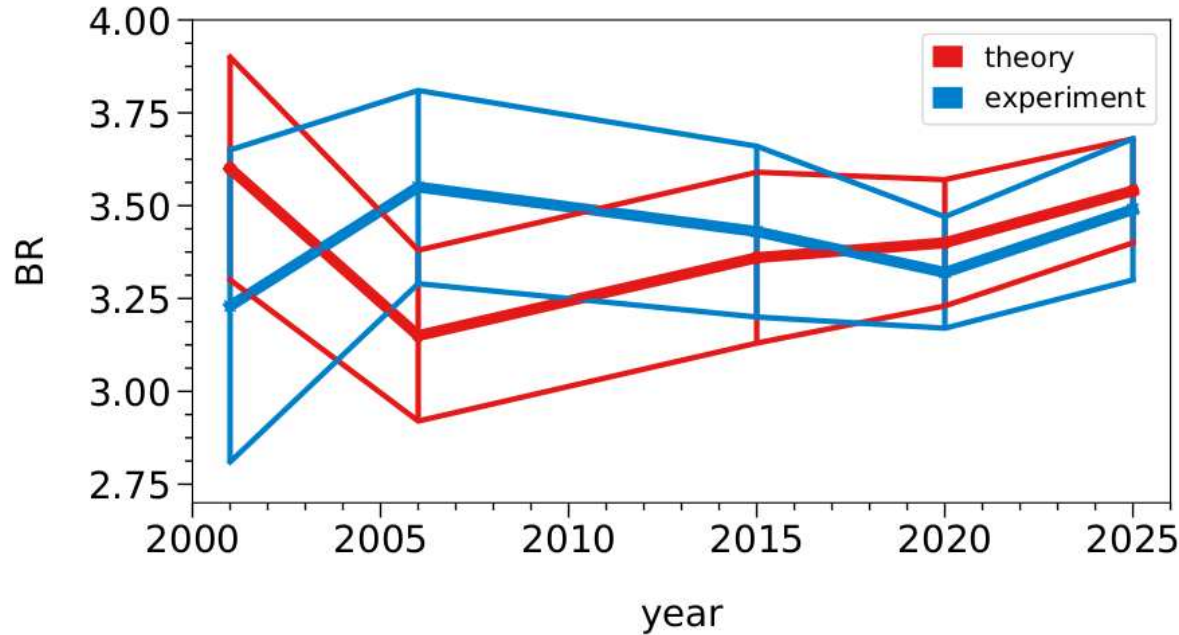
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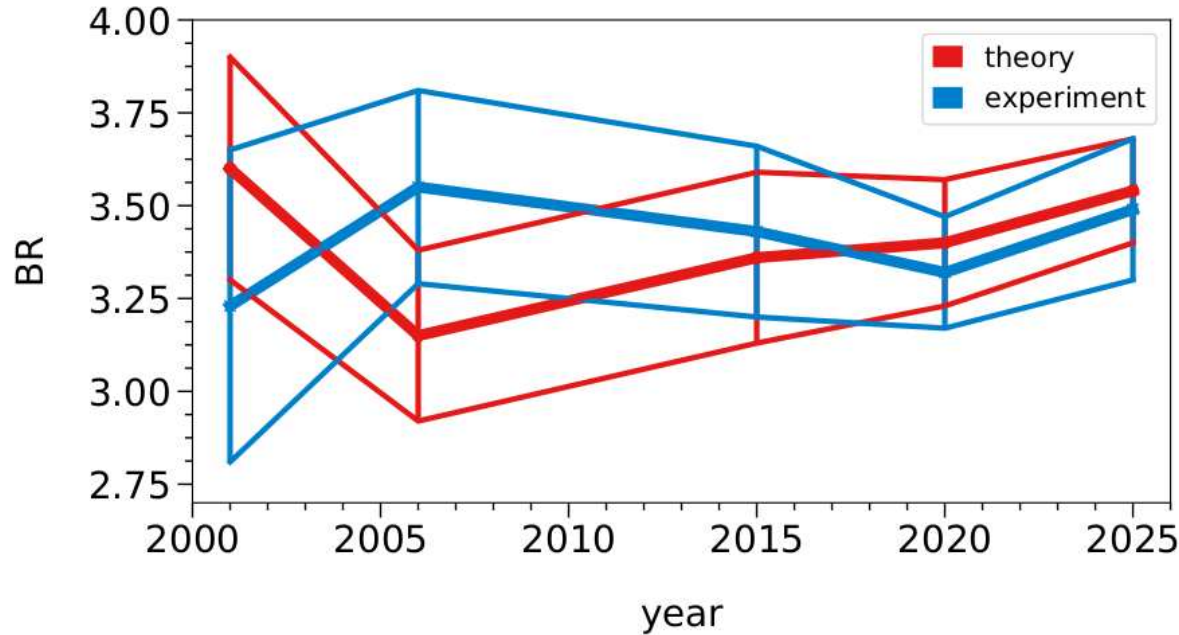
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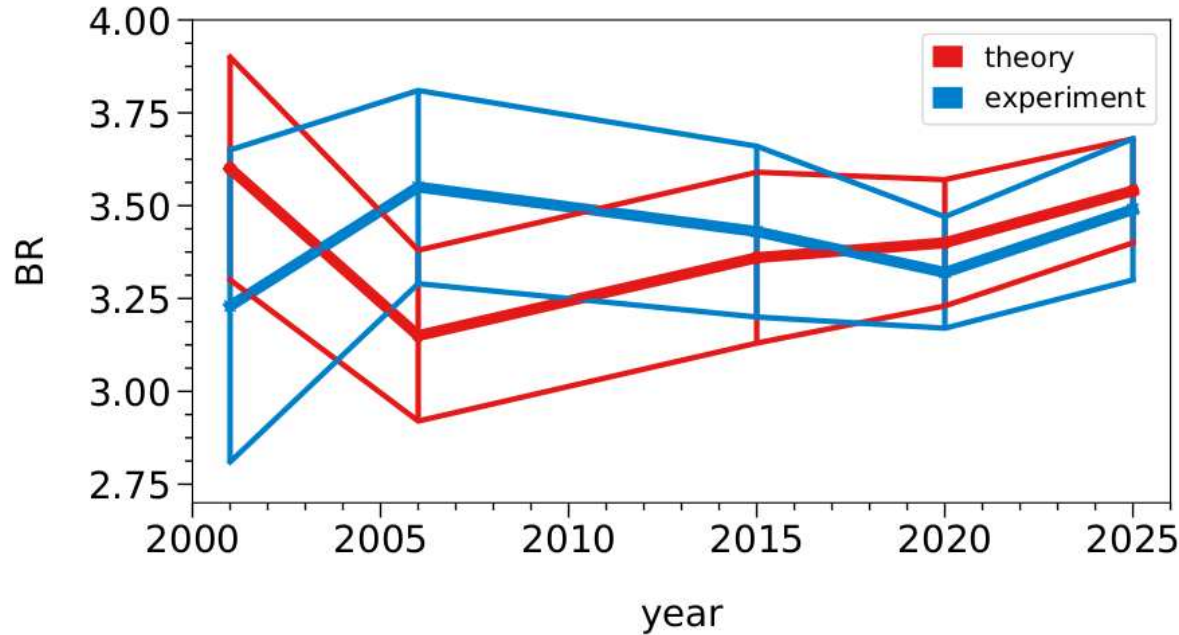


Improvements in the dominant charm-mass dependent NNLO QCD corrections:

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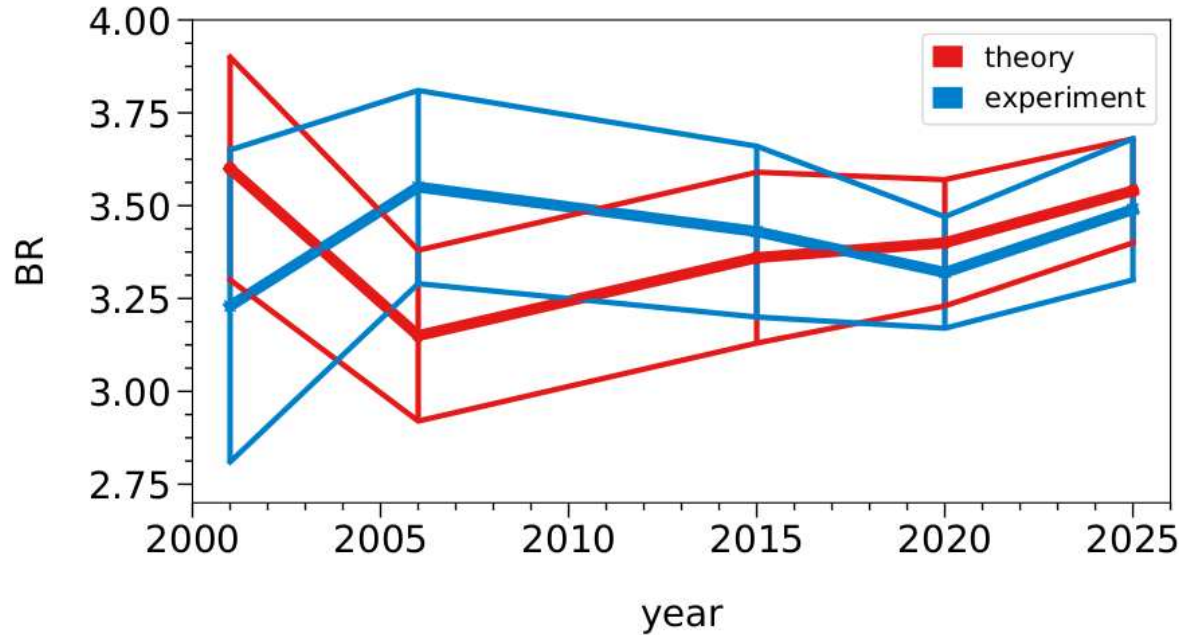
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2000: NLO only,

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Time evolution of selected SM predictions for $\mathcal{B}_{s\gamma}$ and experimental averages:



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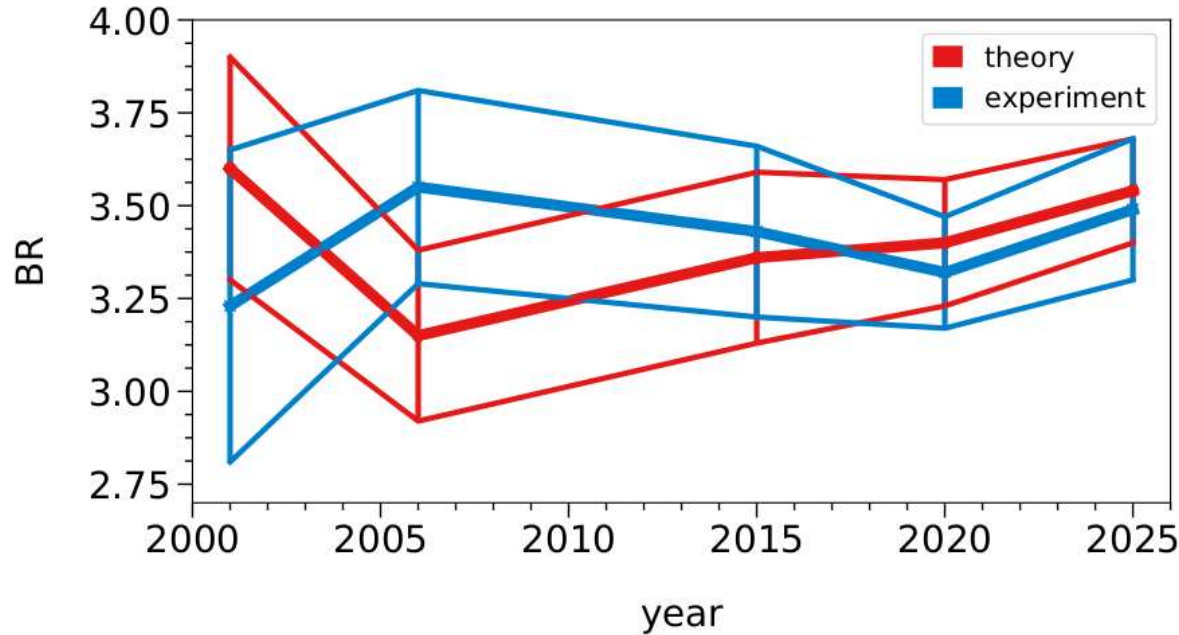
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2006: extrapolation of non-BLM NNLO from the $m_c \gg m_b$ limit,

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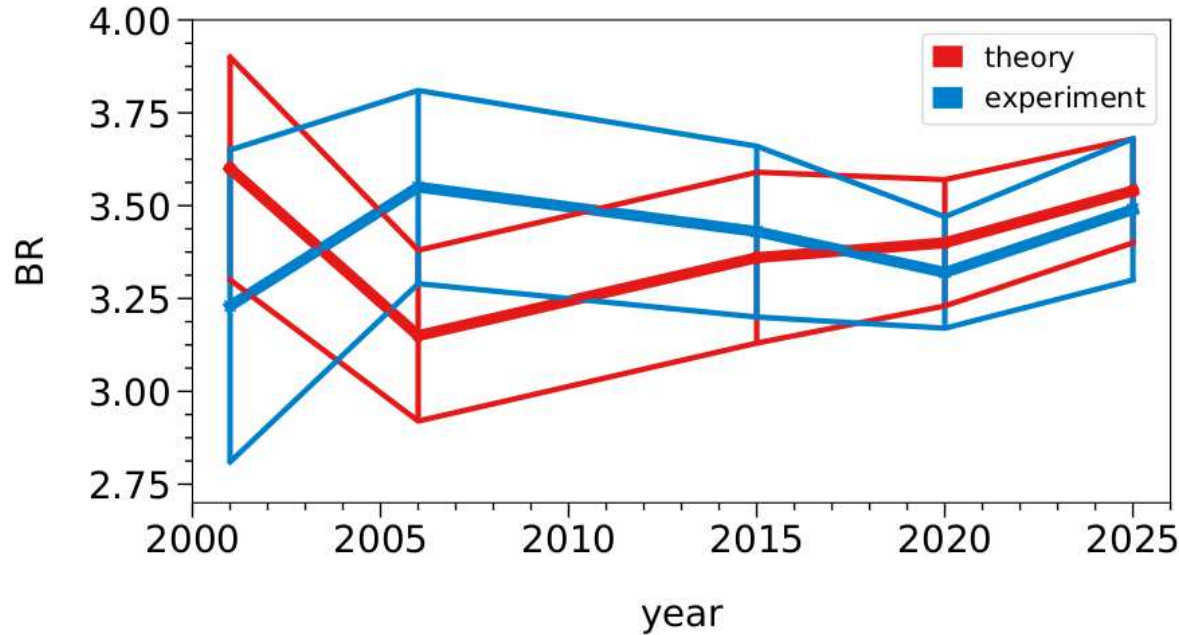
2006: extrapolation of non-BLM NNLO from the $m_c \gg m_b$ limit,

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Updated SM prediction:

$$\mathcal{B}_{s\gamma}^{\text{SM}} = (3.54 \pm 0.14) \times 10^{-4} \left[\pm 4.0\% \simeq \pm \sqrt{(3.0\%)_{\text{h.o.}}^2 + (2.7\%)_{\text{param}}^2} \right].$$

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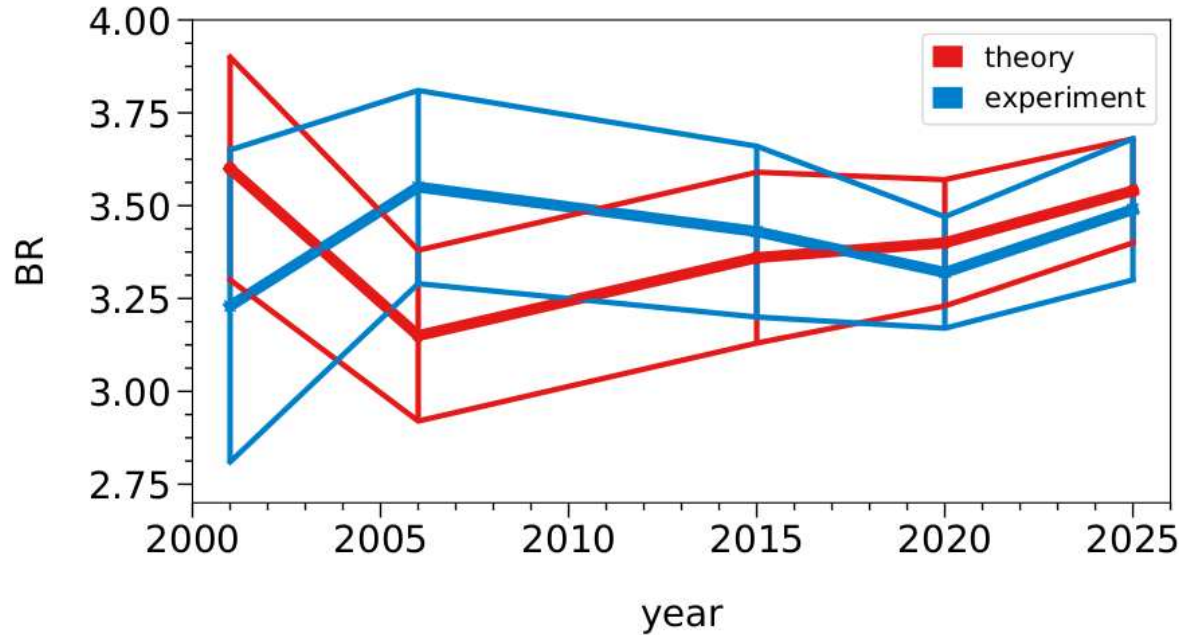
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The perturbative NLO corrections are now formally complete thanks to including formerly missing 4-body and 5-body contributions from arXiv:2510.nnnnn [K. M. Brune, T. Huber, L-T. Moos].

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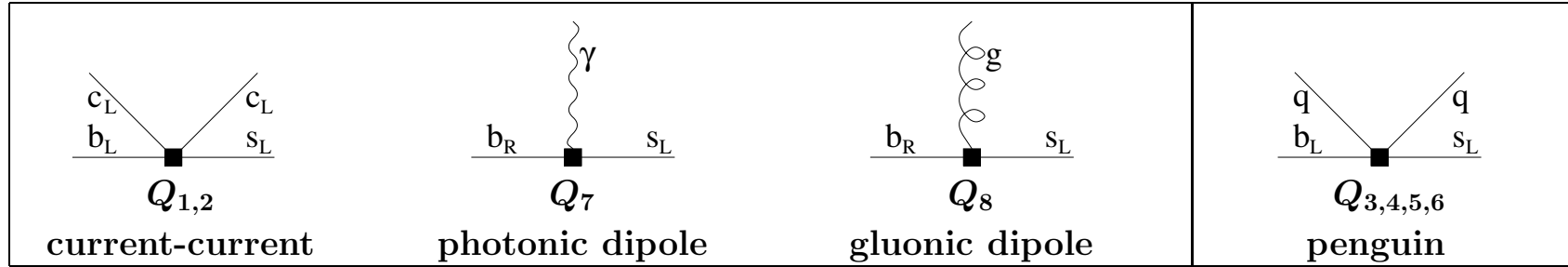
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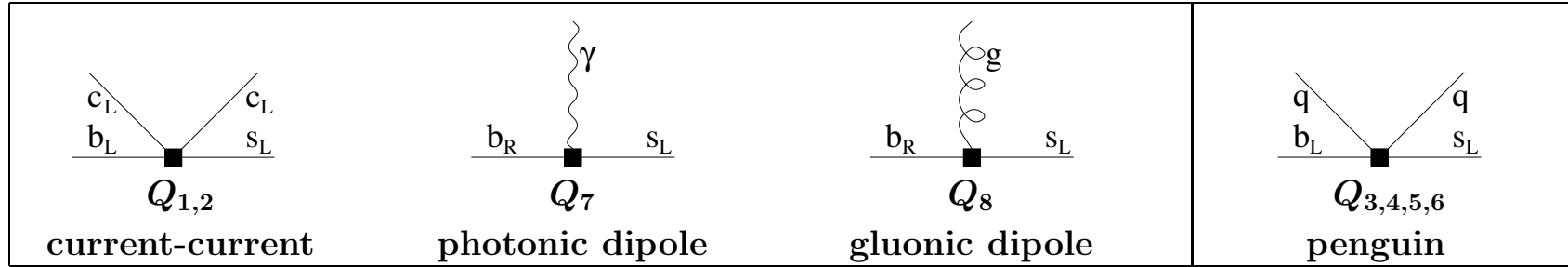
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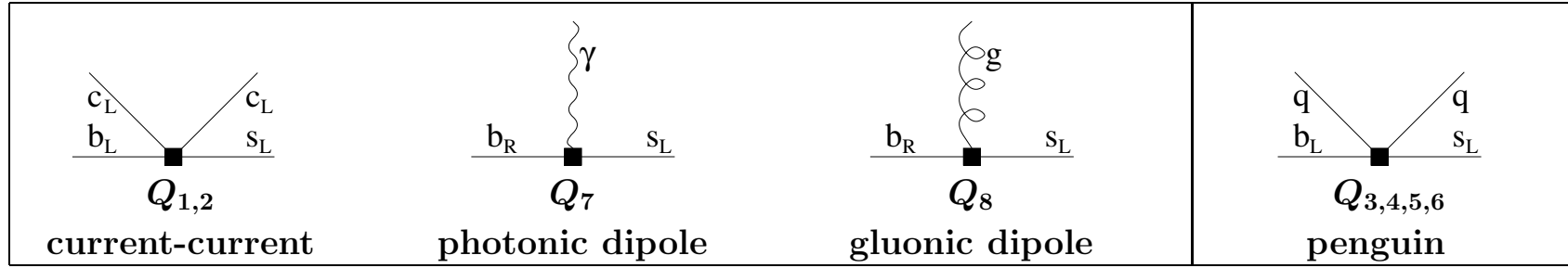
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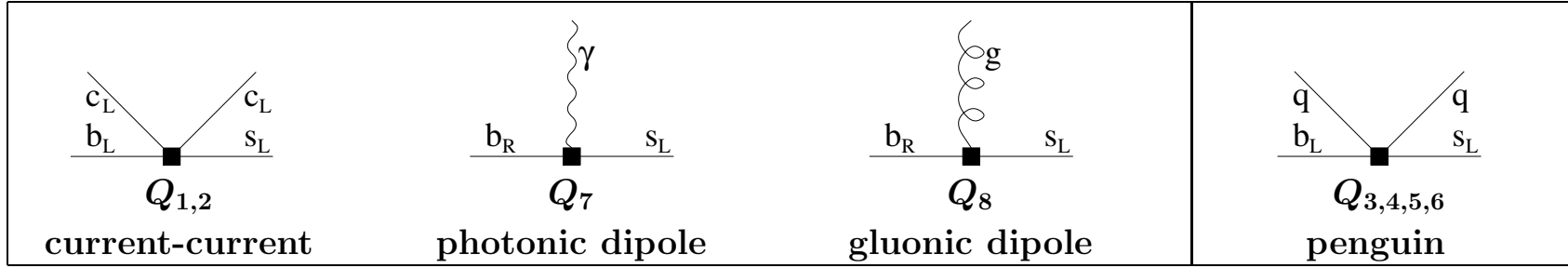
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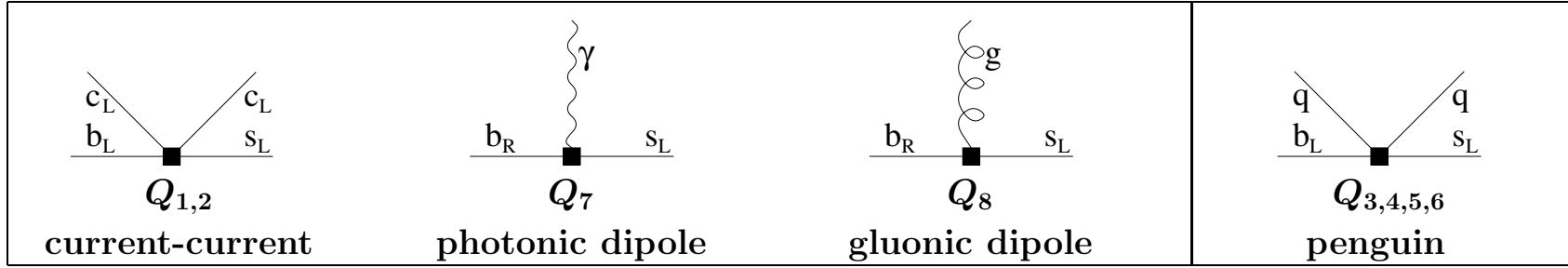
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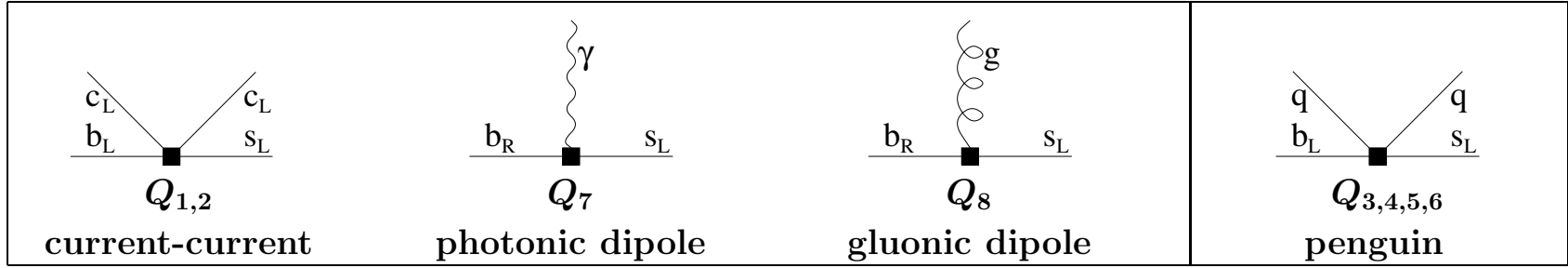
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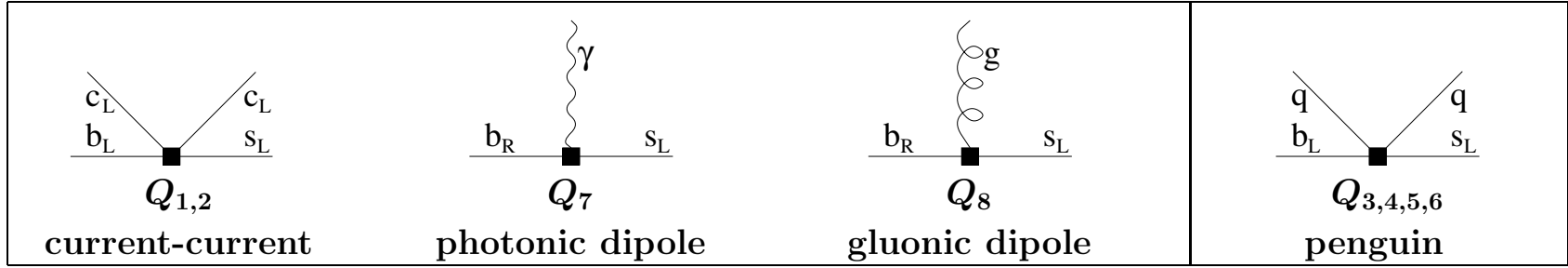
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NLO

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H. M. Asatrian, A. Hovhannisyan, V. Poghosyan, T. Ewerth, C. Greub and T. Hurth, hep-ph/0605009,

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$G_{78}^{(2)}$: H. M. Asatrian, T. Ewerth, A. Ferroglia, C. Greub and G. Ossola, arXiv:1005.5587.

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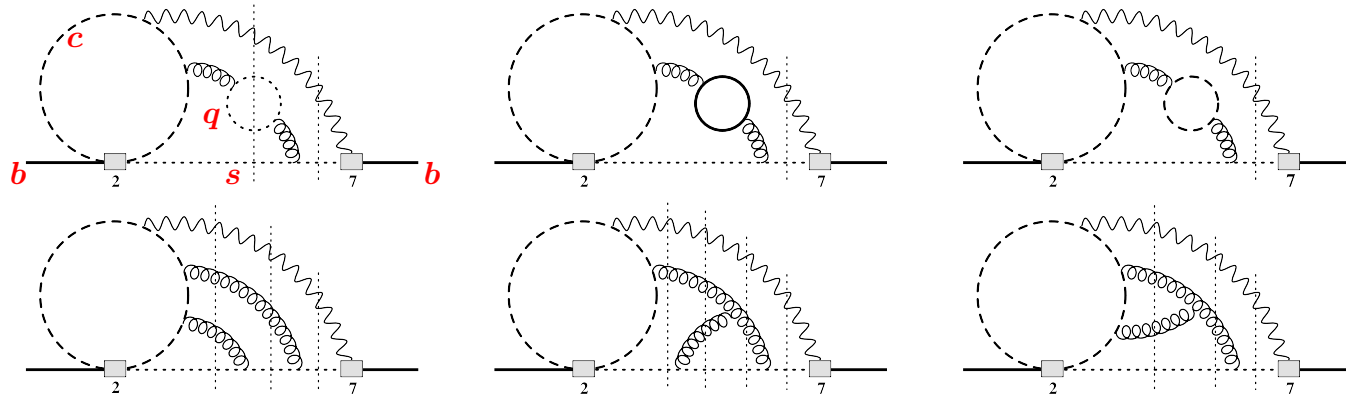
NNLO

$G_{77}^{(2)}$: K. Melnikov and A. Mitov, hep-ph/0505097,
I. Blokland, A. Czarnecki, MM, M. Ślusarczyk and F. Tkachov, hep-ph/0506055,
H. M. Asatryan, A. Hovhannisyan, V. Poghosyan, T. Ewerth, C. Greub and T. Hurth, hep-ph/0605009,
H. M. Asatryan, T. Ewerth, A. Ferroglia, P. Gambino and C. Greub, hep-ph/0607316,
H. M. Asatryan, T. Ewerth, H. Gabrielyan and C. Greub, hep-ph/0611123.

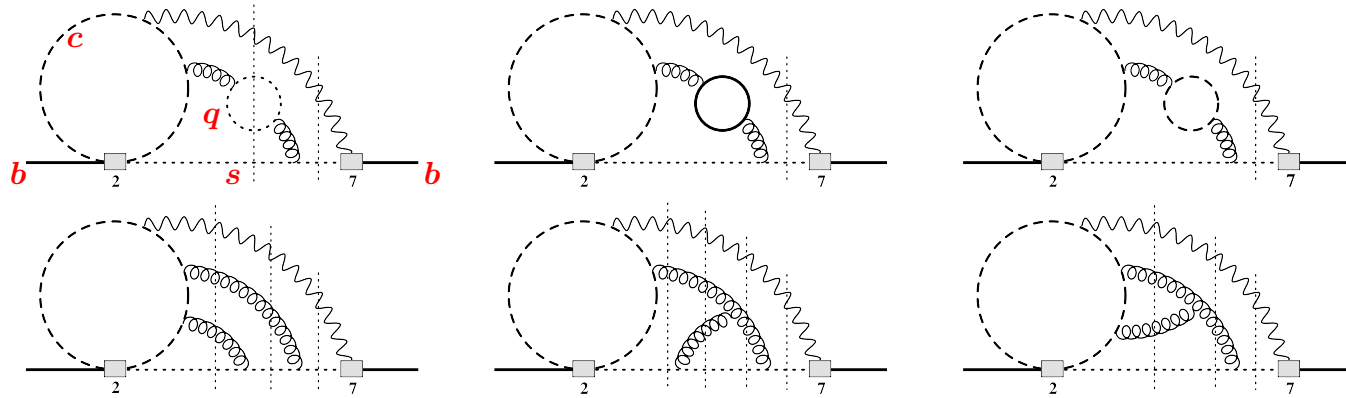
$G_{78}^{(2)}$: H. M. Asatryan, T. Ewerth, A. Ferroglia, C. Greub and G. Ossola, arXiv:1005.5587.

$G_{27}^{(2)}$: Z. Ligeti, M.E. Luke, A.V. Manohar and M.B. Wise, hep-ph/9903305, **large- β_0 , 4-body**,
K. Bieri, C. Greub and M. Steinhauser, hep-ph/0302051, **large- β_0 , 2-body**,
MM and M. Steinhauser, hep-ph/0609241, arXiv:1005.1173, **$m_c \gg m_b$** ,
R. Boughezal, M. Czakon and T. Schutzmeier, arXiv:0707.3090, **large- β_0 and massive loops, 2-body**,
M. Czakon, P. Fiedler, T. Huber, MM, T. Schutzmeier and M. Steinhauser, arXiv:1503.01791, **$m_c = 0$** ,
MM, A. Rehman and M. Steinhauser, arXiv:1702.07674, arXiv:2002.01548, **counterterms, large- β_0 and massive loops**,
C. Greub, H. M. Asatryan, F. Saturnino and C. Wiegand, arXiv:2303.01714, **physical m_c , partial 2-body**,
M. Fael, F. Lange, K. Schönwald and M. Steinhauser, arXiv:2309.14706, **physical m_c , full 2-body**,
M. Czaja, M. Czakon, T. Huber, MM, M. Niggetiedt, A. Rehman,
K. Schönwald and M. Steinhauser, arXiv:2309.14707, **physical m_c , full 2-body**,
C. Greub, H. M. Asatryan, H. H. Asatryan, L. Born and J. Eicher, arXiv:2407.17270, **physical m_c , full 2-body**,
M. Czaja, M. Czakon, T. Huber, MM, M. Niggetiedt, A. Rehman,
K. Schönwald and M. Steinhauser, arXiv:2510.nnnnn, **physical m_c , fully inclusive, renormalized result**.

Sample propagator diagrams with cuts contributing to G_{27} @ NNLO:

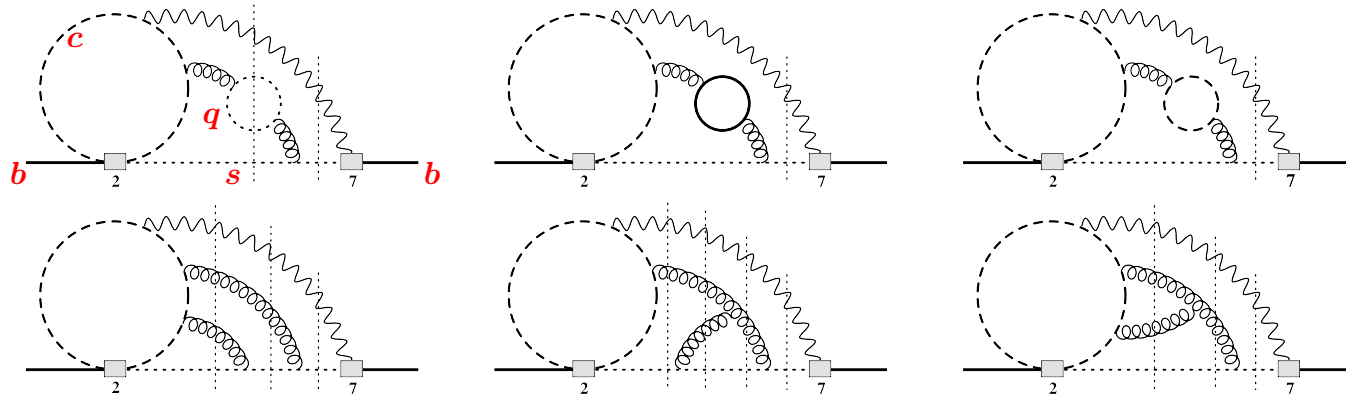


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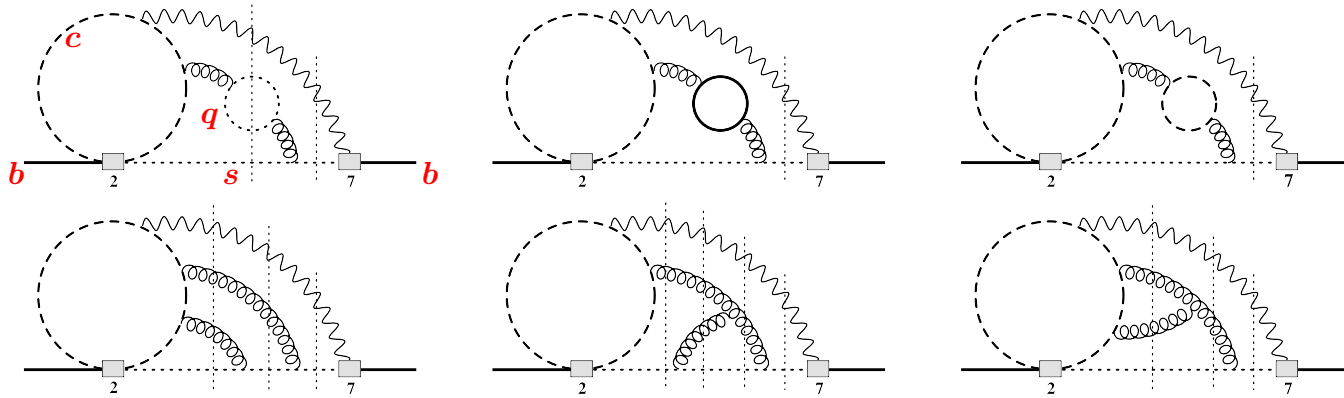
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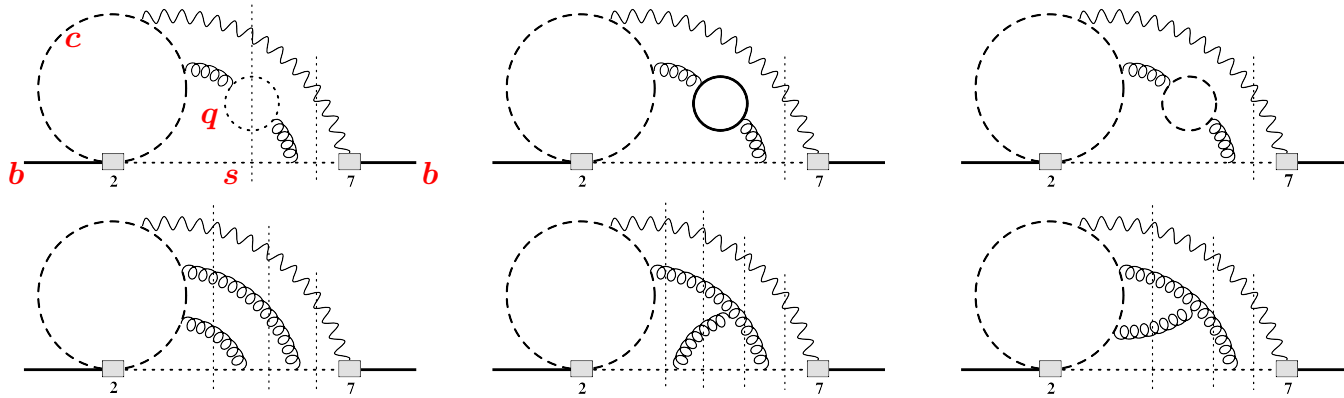


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3. Extending the set of master integrals M_k so that it closes under differentiation with respect to $z = m_c^2/m_b^2$. This way one obtains a system of differential equations

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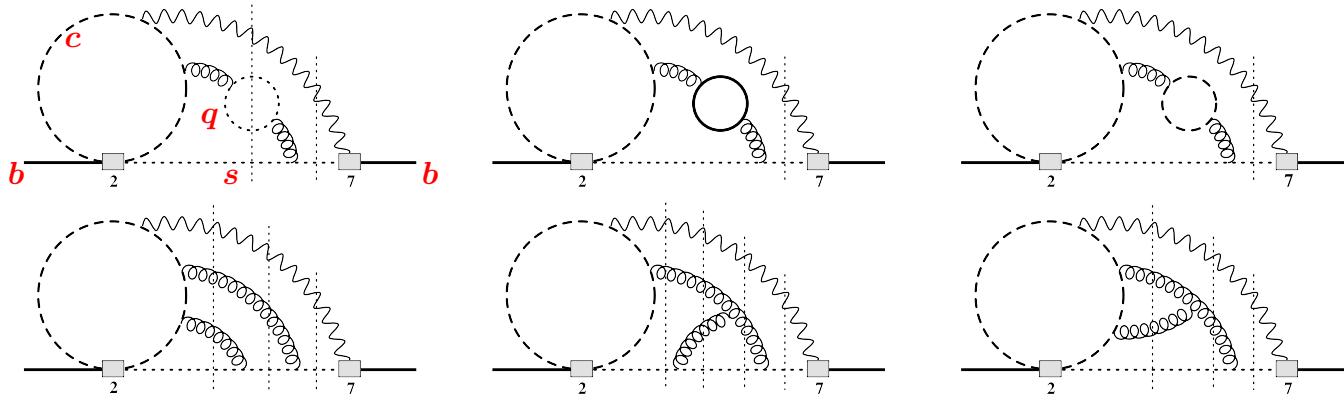
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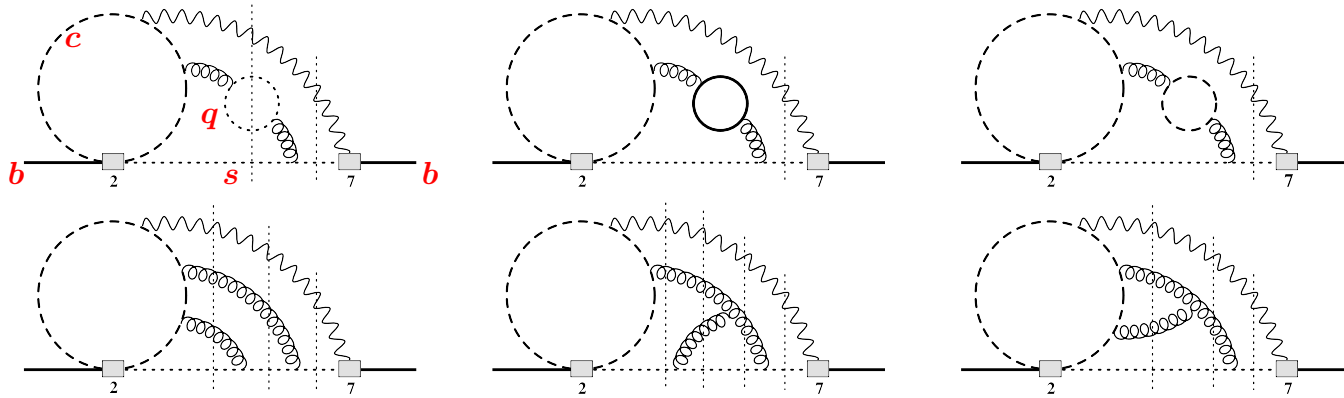
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Fully inclusive (2-, 3- and 4-body), renormalized results for $G_{17}^{(2)}$ and $G_{27}^{(2)}$.

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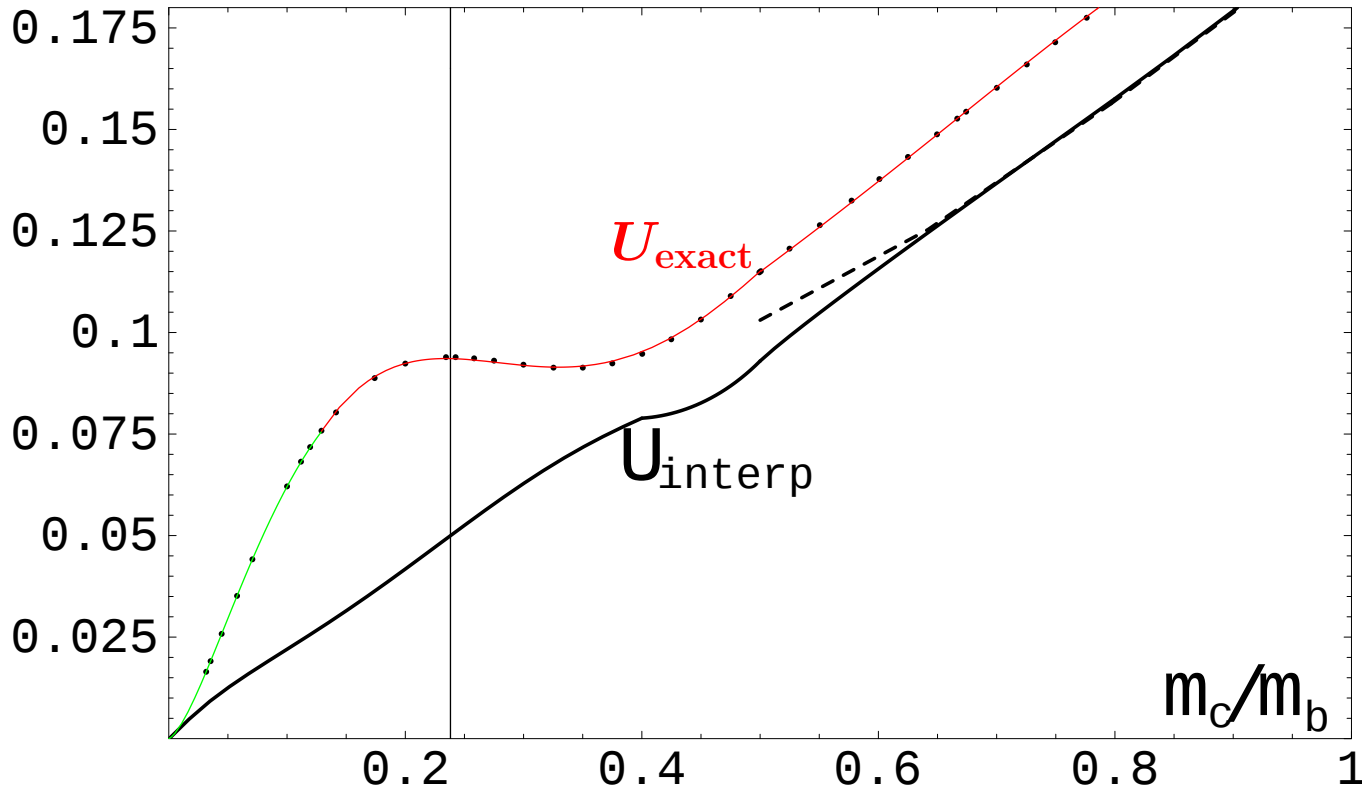
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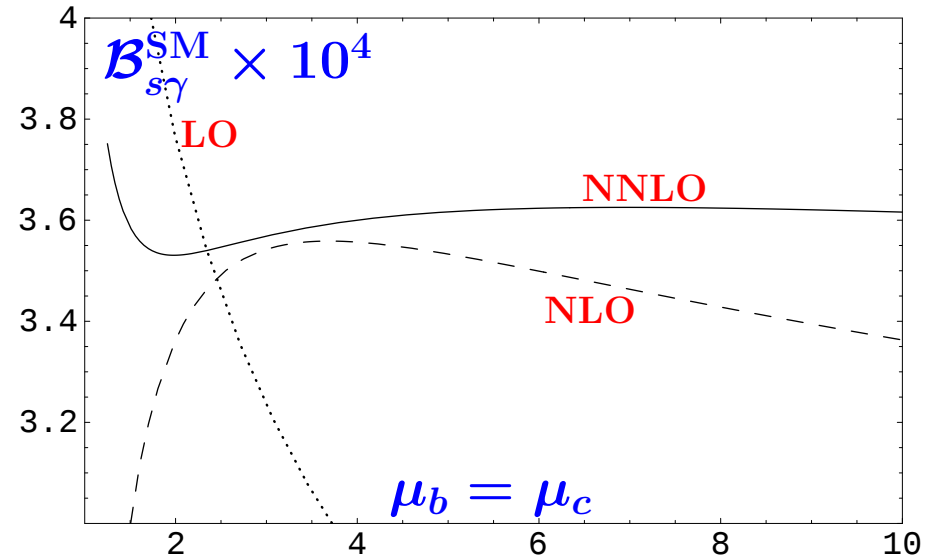
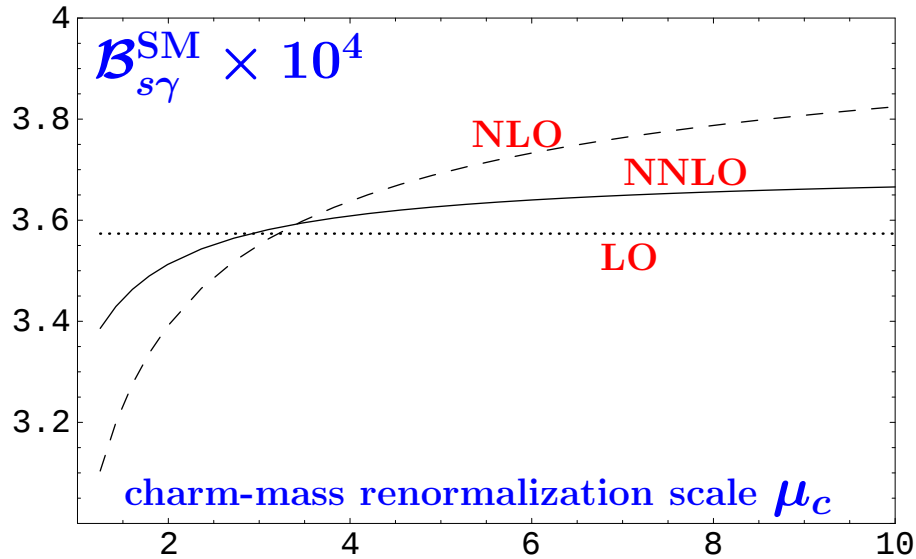
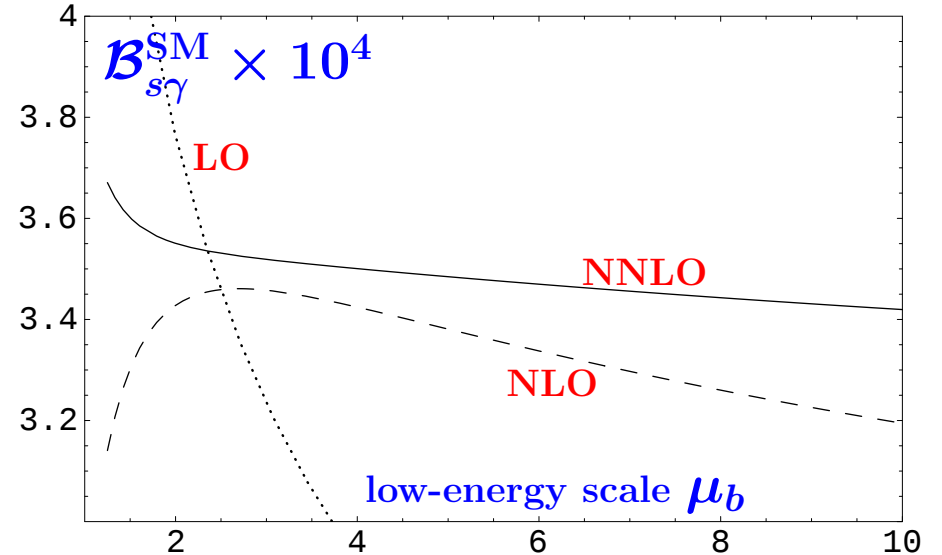
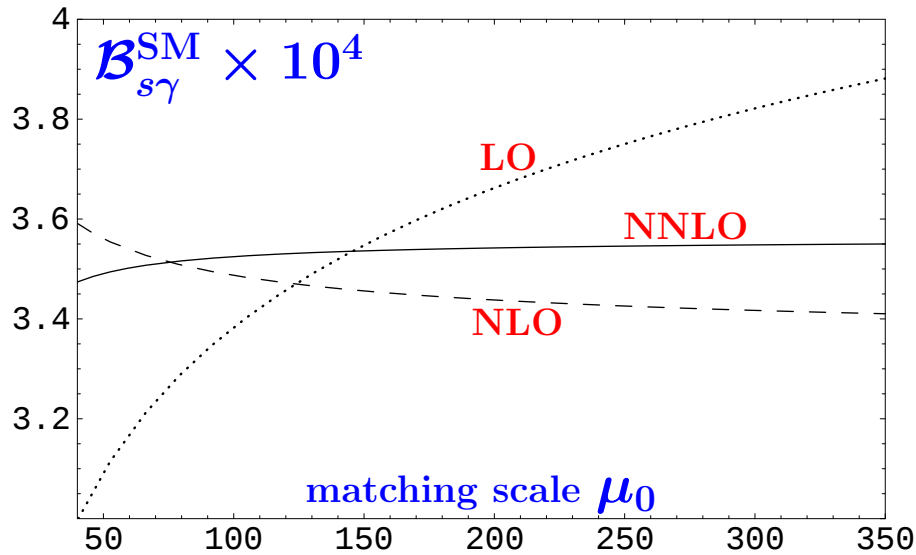
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Interpolated and exact results for $\delta = 1$ (no cut on E_γ):



Renormalization scale dependence of $\mathcal{B}_{s\gamma}^{\text{SM}}$ for $E_\gamma > 1.6$ GeV



“Central” values: $\mu_0 = 160$ GeV, $\mu_b = \mu_c = \frac{1}{2}m_{b,kin}(1 \text{ GeV}) \simeq 2.29$ GeV.

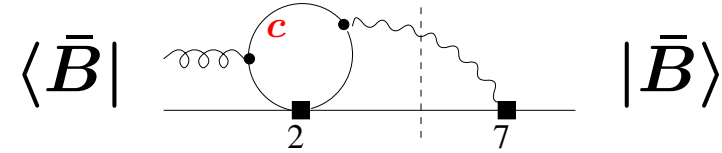
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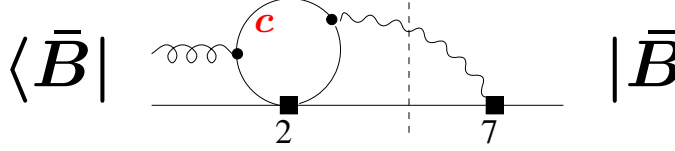
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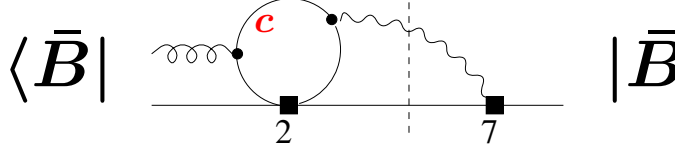
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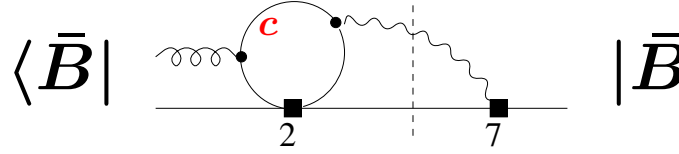
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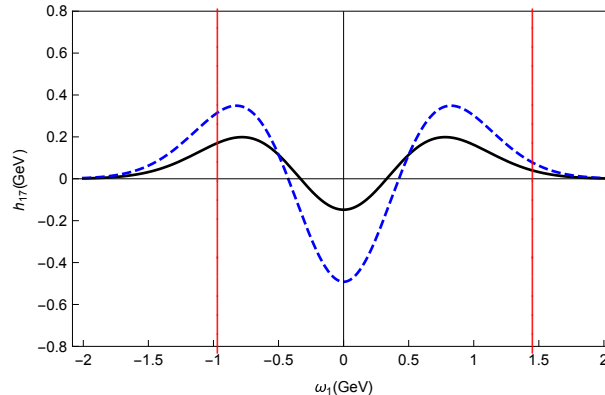
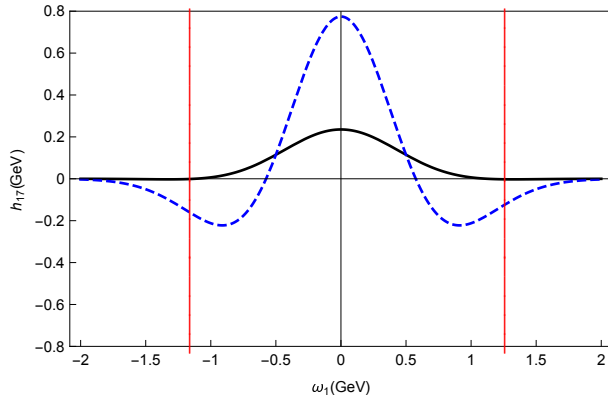
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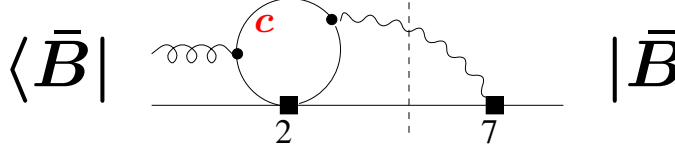
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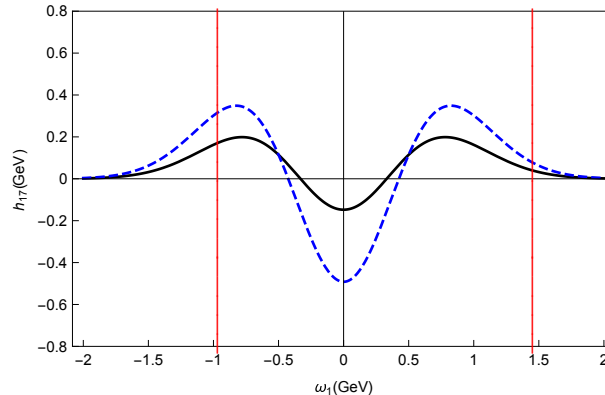
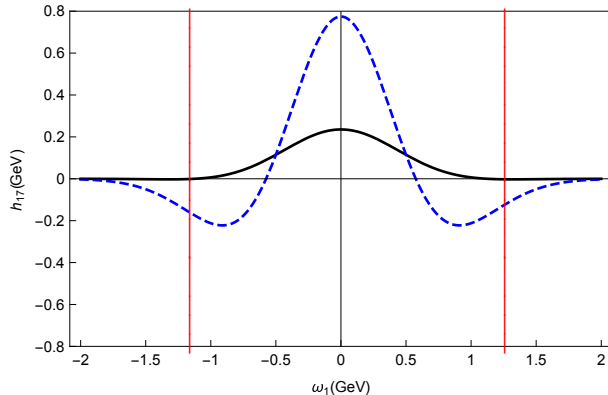
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G+P numerically:
 $\Lambda_{17} \in [-24, 5] \text{ MeV}$ for $m_c = 1.17 \text{ GeV}$.

In our code: $\kappa_V = 1.2 \pm 0.3$.
Warning: scheme for m_c !

Moment constraints vs. models of h_{17}

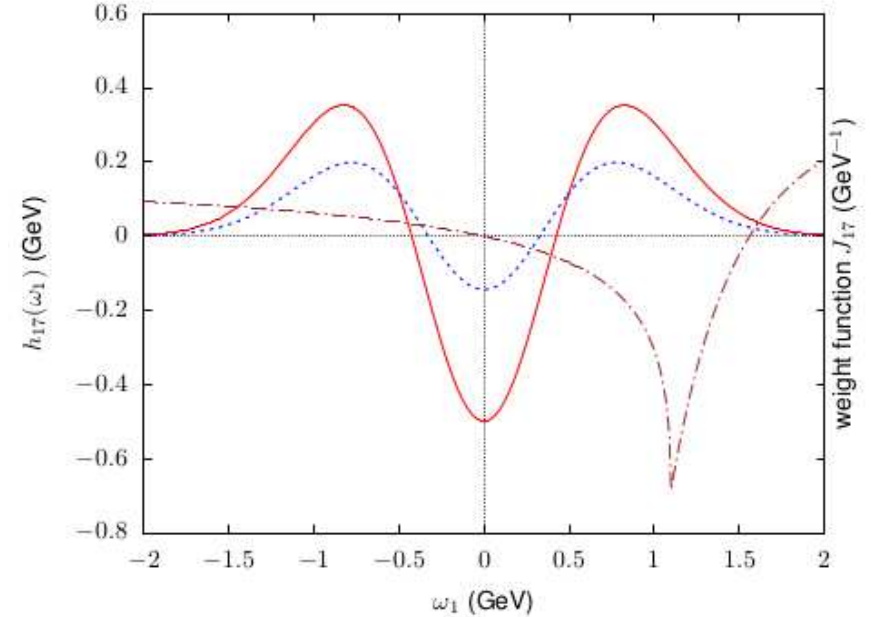
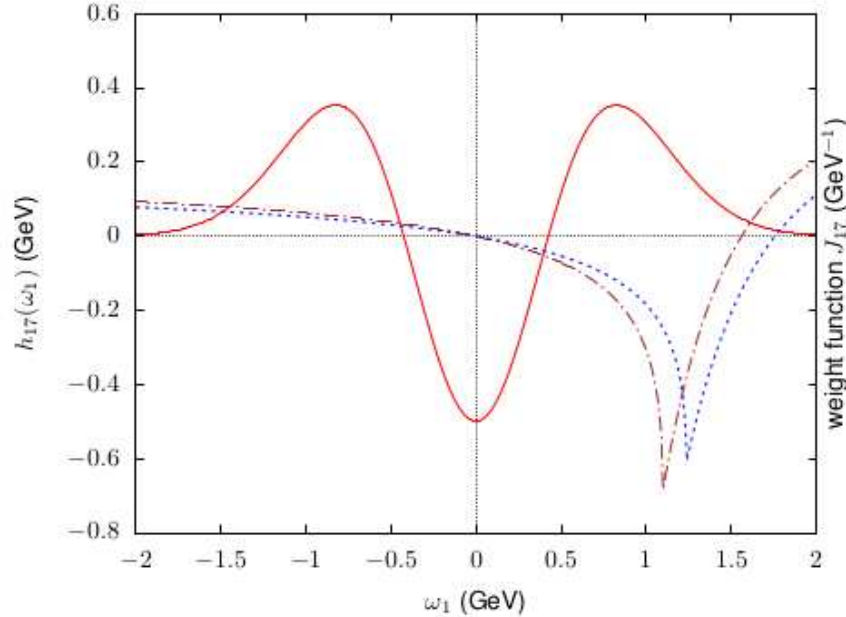
M. Benzke, S.J. Lee, M. Neubert, **G. Paz**, arXiv:1003.5012 – only the leading moment included.

A. Gunawardana, **G. Paz**, arXiv:1908.02812 – estimates of the subleading moments from LLSA included.

M. Benzke, T. Hurth, arXiv:2006.00624 – as above but with more generous modeling and partial $1/m_b^2$ corrections.

R. Bartocci, P. Böer, T. Hurth, arXiv:2411.16634 – RG evolution of $h_{17}(\omega_1, \mu)$.

Plots from arXiv:2006.00624:



Another recent contribution: T. Hurth and R. Szafron, arXiv:2301.01739 – clarifying the SCET treatment of resolved photons in the Q_8 - Q_8 interference.

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$$\mathcal{B}_{s\gamma}^{\text{SM}} = (3.54 \pm 0.14) \times 10^{-4} \quad (\pm 4.0\%).$$

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- Non-perturbative outlook:

Complete the $\mathcal{O}\left(\frac{1}{m_b^2}\right)$ and $\mathcal{O}(\alpha_s)$ resolved-photon corrections in the $Q_{1,2}$ - Q_7 interference.

Use arXiv:2301.01739 to update the Q_8 - Q_8 interference.

Use arXiv:2211.07663 [B. Dehnadi, I. Novikov, F. J. Tackmann] and the **SIMBA** analysis in arXiv:2007.04320 to improve HFLAV/PDG averaging and extrapolation to $E_0 = 1.6$ GeV.

BACKUP SLIDES

The “hard” contribution to $\bar{B} \rightarrow X_s \gamma$

J. Chay, H. Georgi, B. Grinstein PLB 247 (1990) 399.
A.F. Falk, M. Luke, M. Savage, PRD 49 (1994) 3367.

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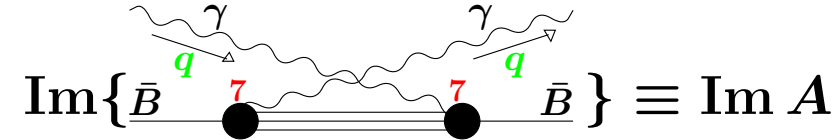
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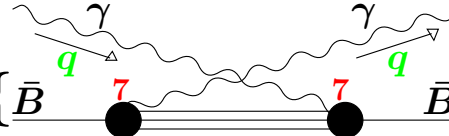
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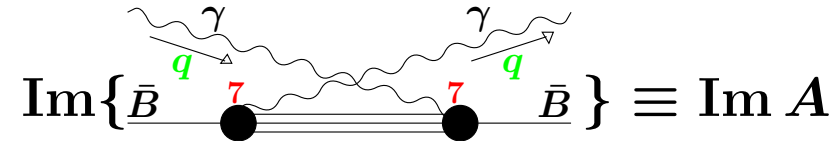
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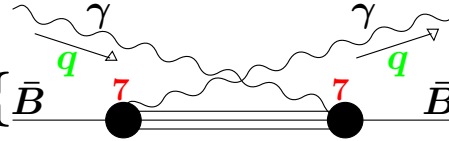
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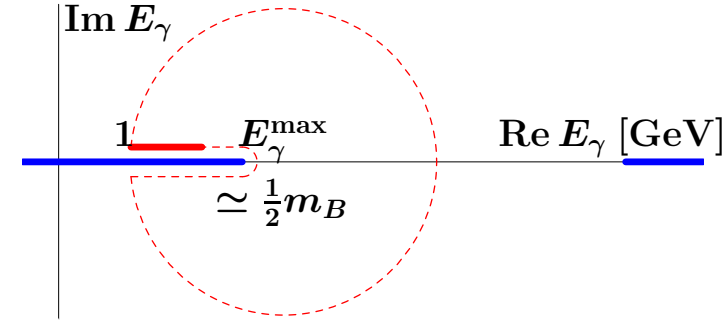
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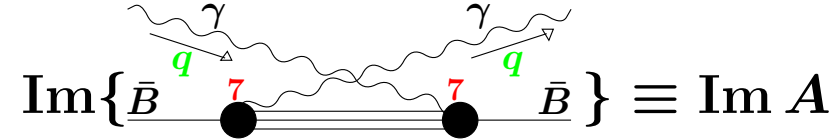


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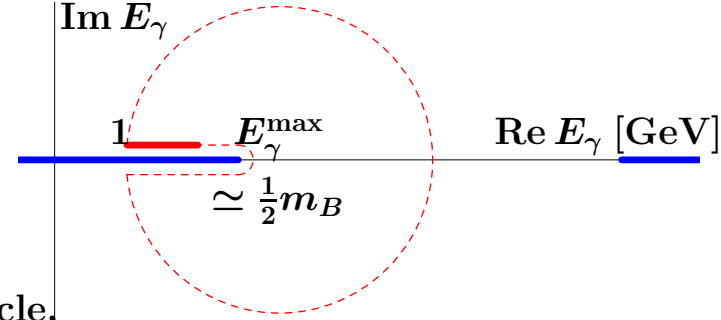
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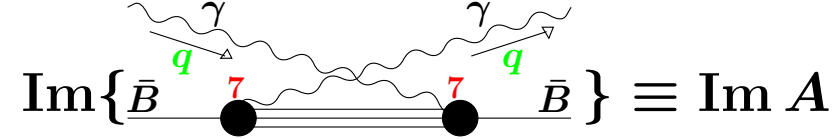
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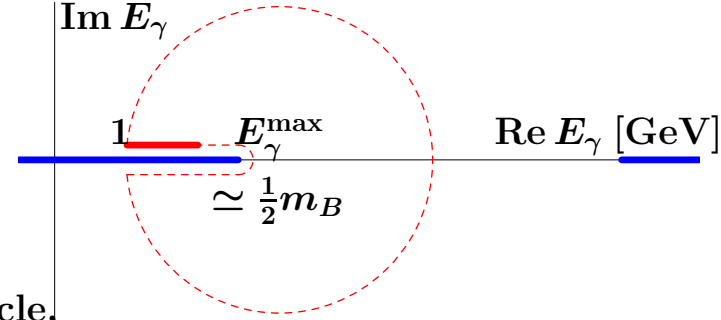
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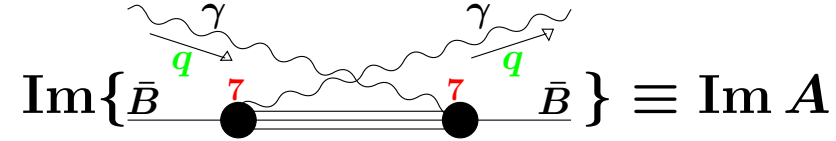
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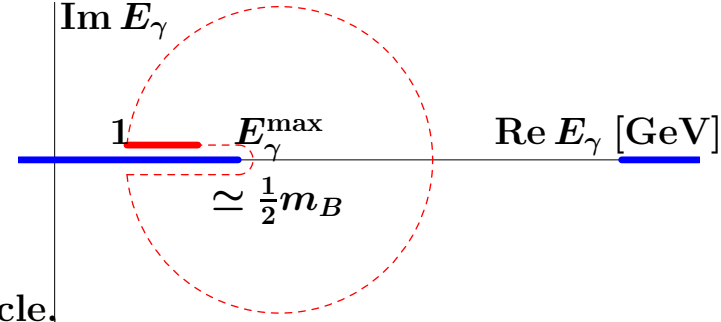
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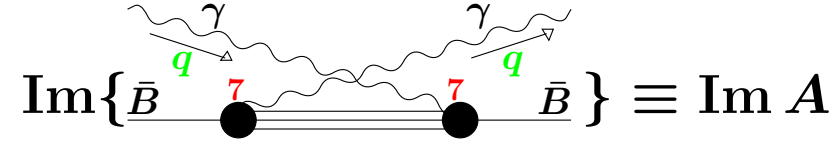
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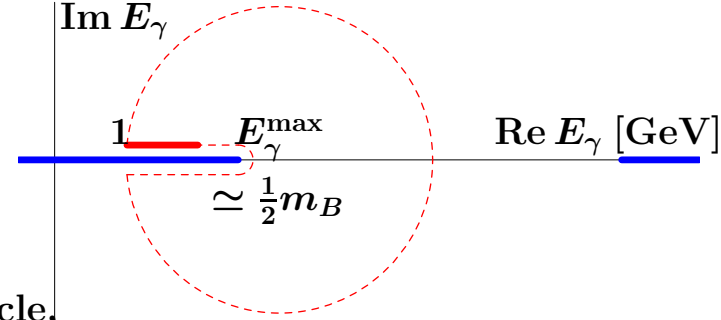
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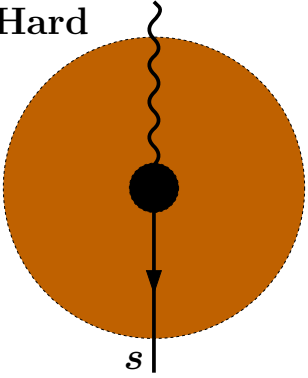
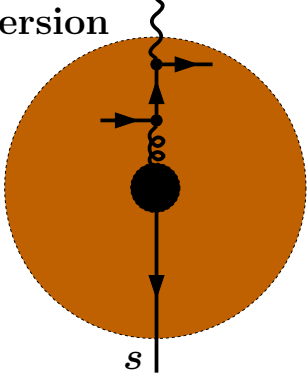
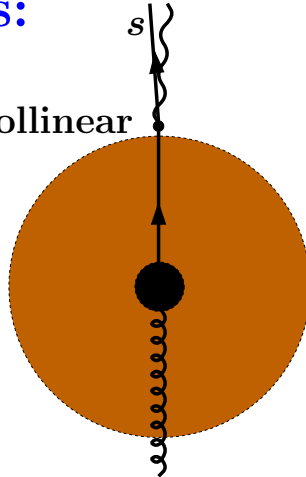
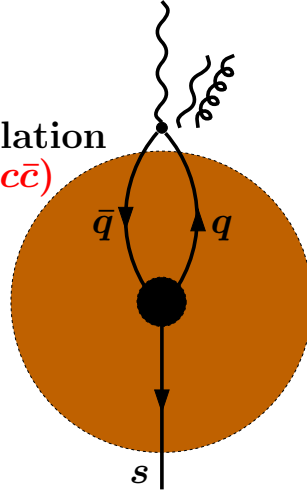
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The HQET heavy-quark field: $b_v(x) = \frac{1}{2}(1 + \not{v})b(x) \exp(im_b v \cdot x)$ with $v = p/m_B$.

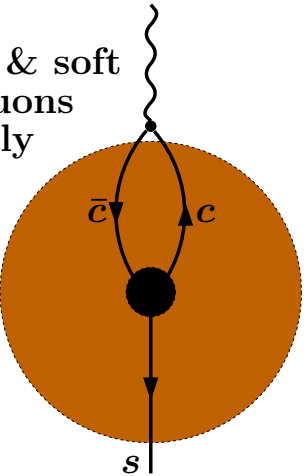
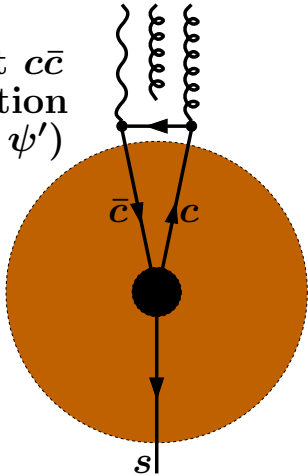
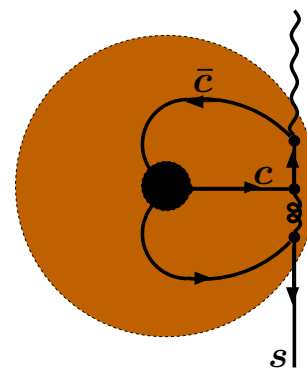
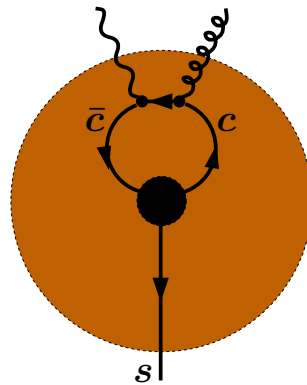
Energetic photon production in charmless decays of the $\bar{B}$ -meson

($E_\gamma \gtrsim \frac{m_b}{3} \simeq 1.6 \text{ GeV}$)

A. Without long-distance charm loops:

1. Hard

 Dominant, well-controlled.
2. Conversion

 $\mathcal{O}(\alpha_s \Lambda/m_b)$,
 $\leftrightarrow$ isospin asymmetry.
3. Collinear

 Perturbatively $< 1\%$,
 $\leftrightarrow$ fragmentation functions.
4. Annihilation
 ($q\bar{q} \neq c\bar{c}$)

 Exp. $\pi^0, \eta, \eta', \omega$ subtracted,
 perturbatively $\sim 0.1\%$.

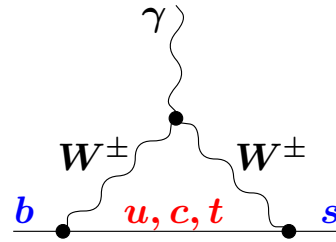
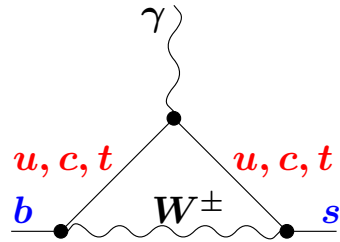
B. With long-distance charm loops:

5. $c\bar{c}$ & soft gluons only

 HQET & SCET.
6. Boosted light $c\bar{c}$ state annihilation
 (e.g. $\eta_c, J/\psi, \psi'$)

 Exp. J/ψ subtracted ($< 1\%$).
7. Annihilation of $c\bar{c}$ in a heavy $(\bar{c}s)(\bar{q}c)$ state



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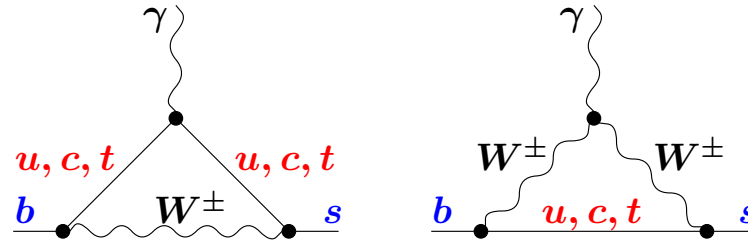
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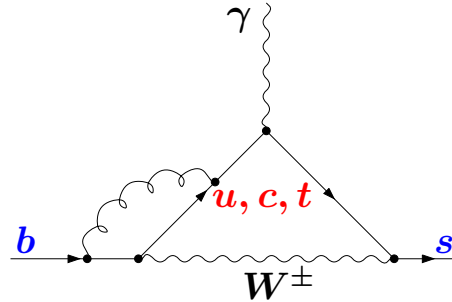
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NLO:

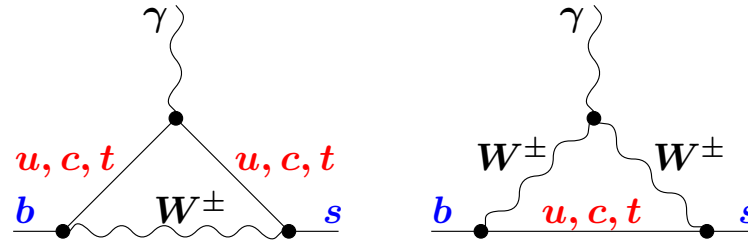
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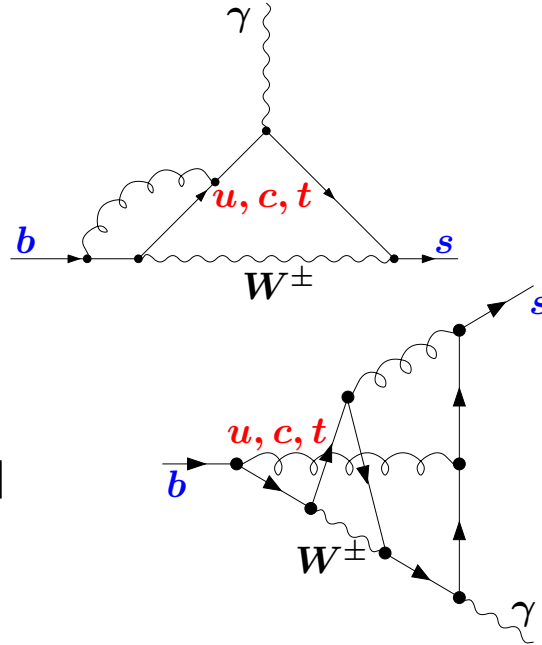
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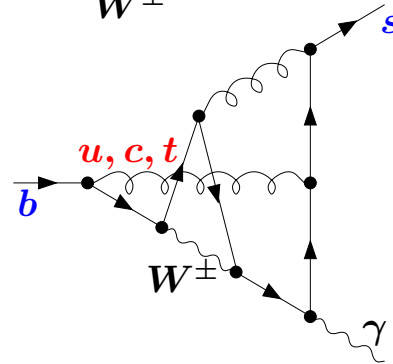
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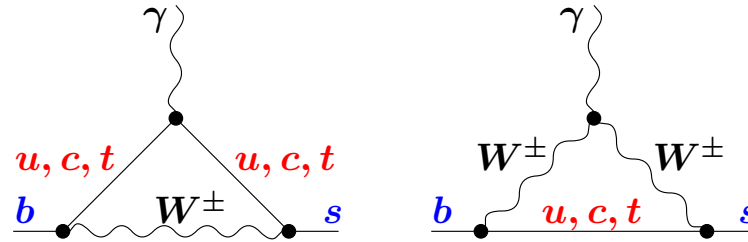
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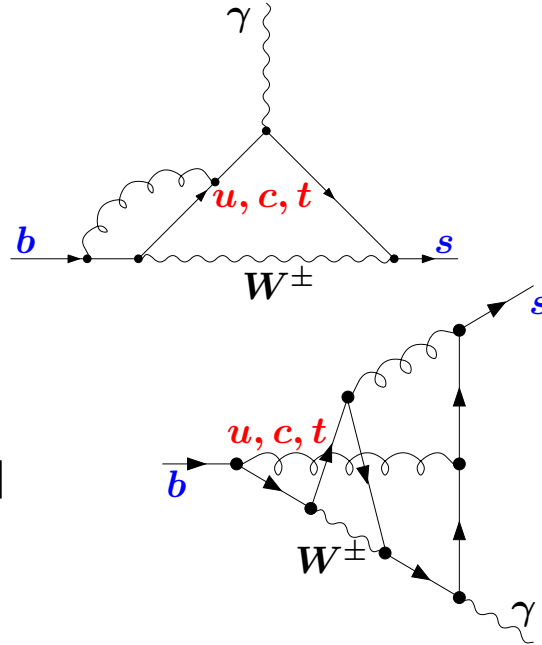
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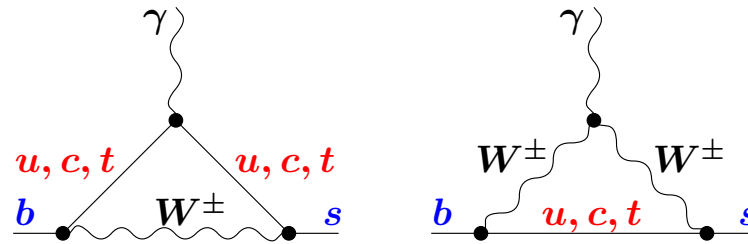
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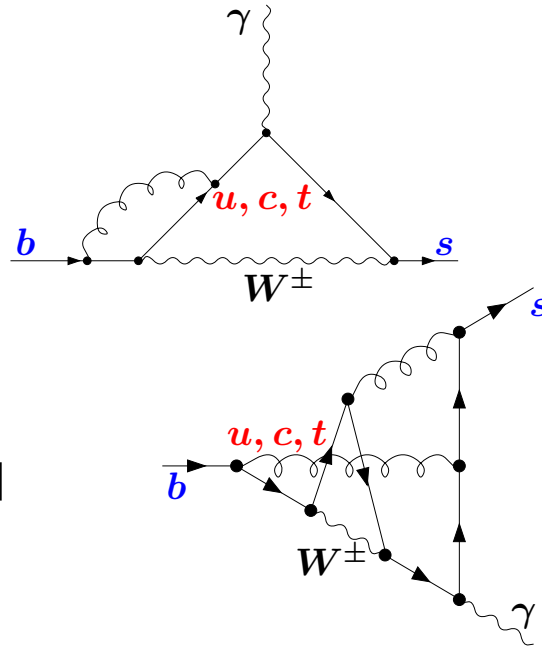
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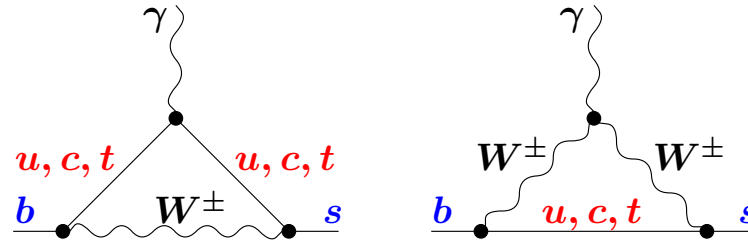
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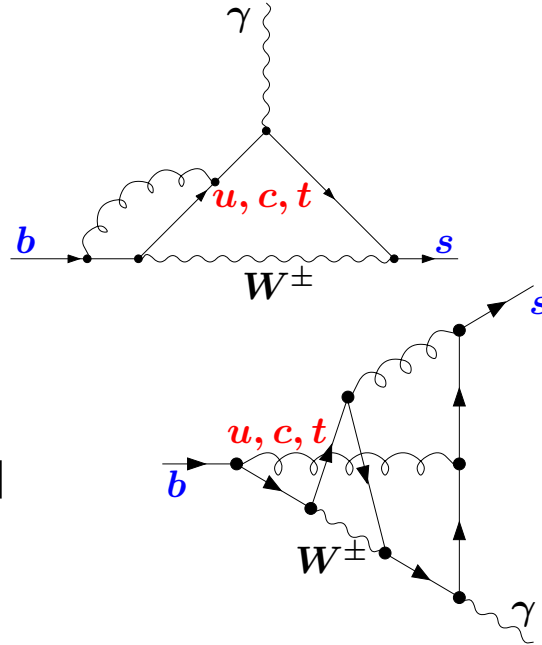
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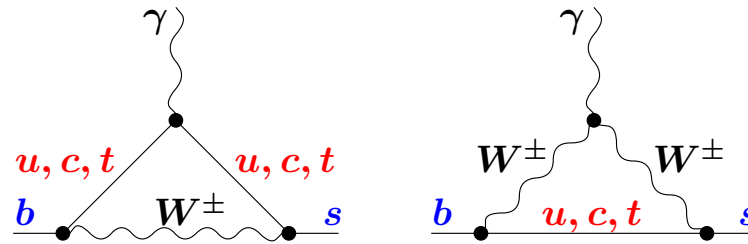
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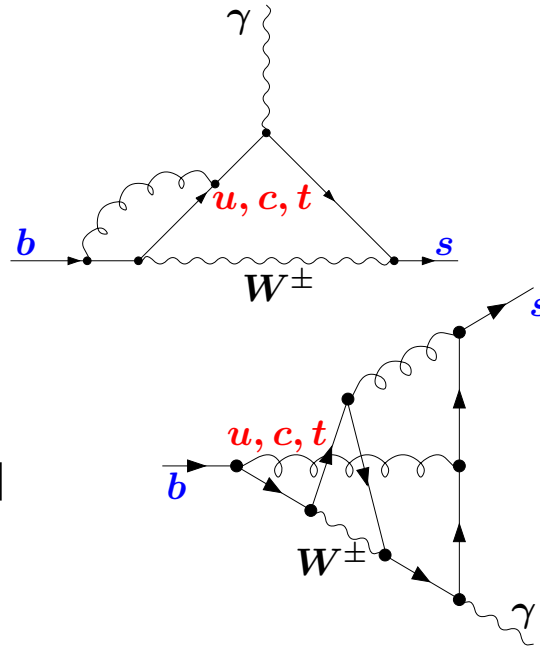
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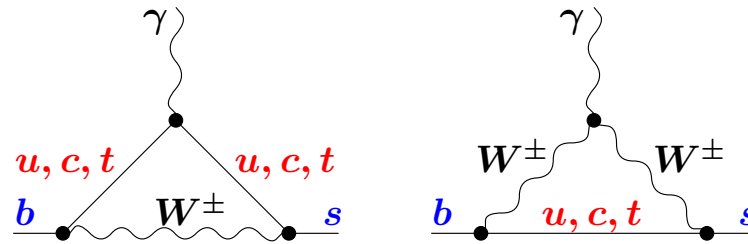
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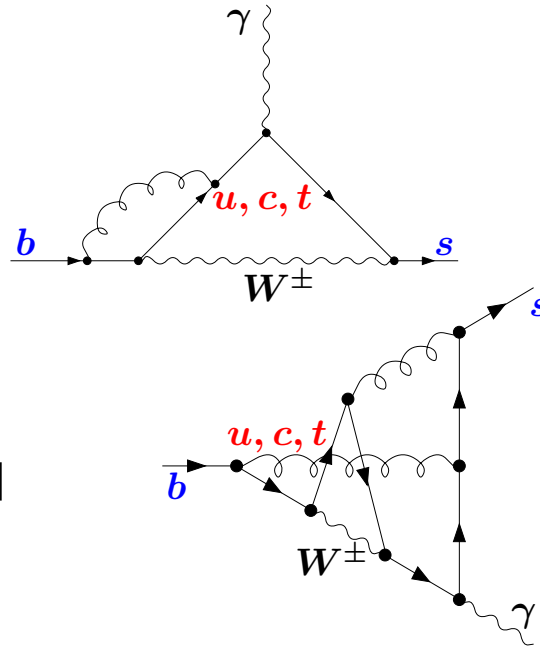
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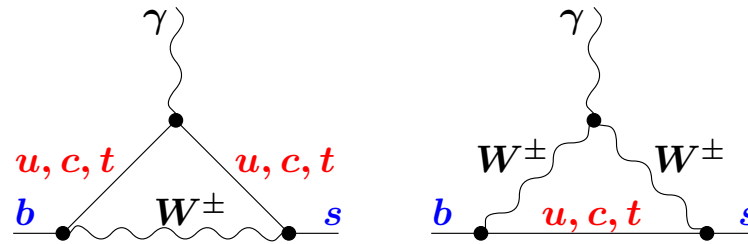
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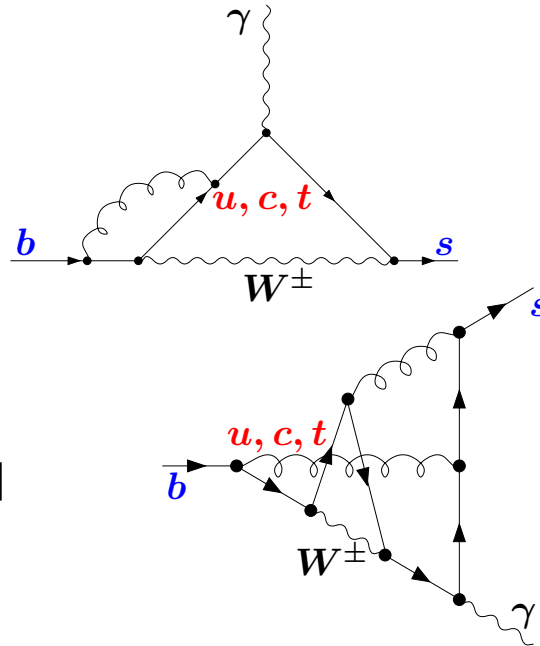
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- At the 3-loop level, the difference $m_t - M_W$ is taken into account with the help of expansions in y^n and $(1 - y^2)^n$ up to $n = 8$, where $y = M_W/m_t$.

Resummation of large logarithms $\left(\alpha_s \ln \frac{\mu_0^2}{\mu_b^2}\right)^n$ in the $b \rightarrow s\gamma$ amplitude.

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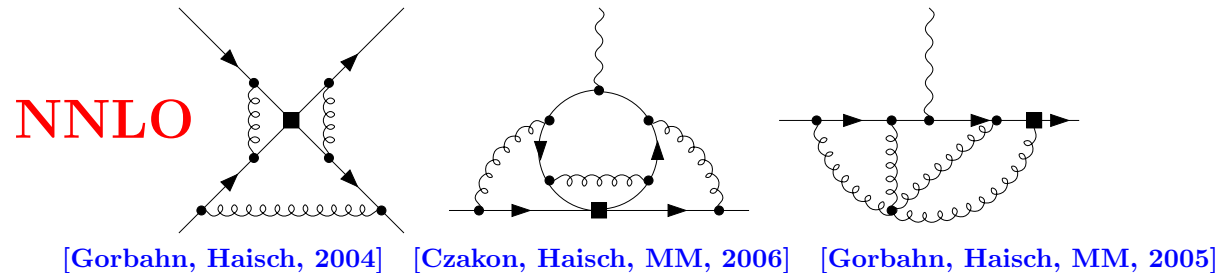
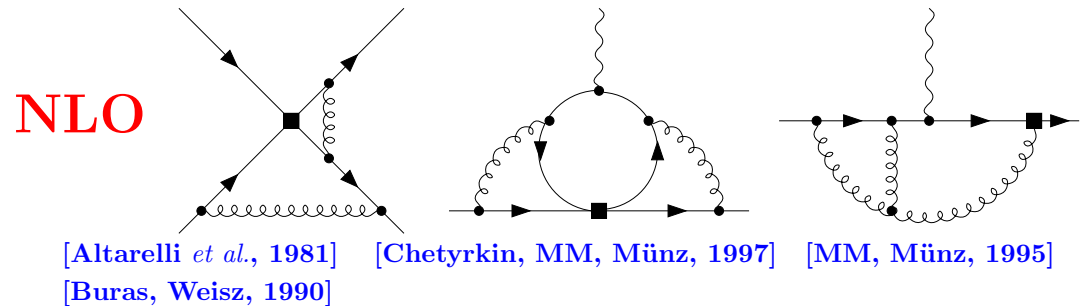
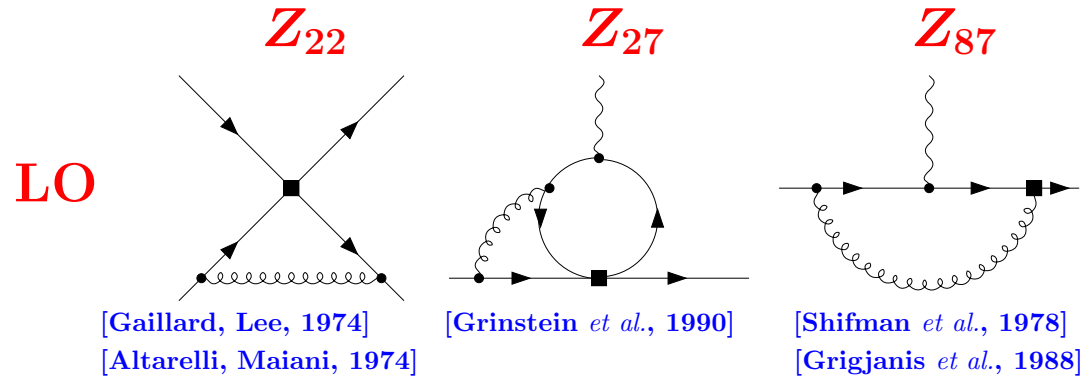
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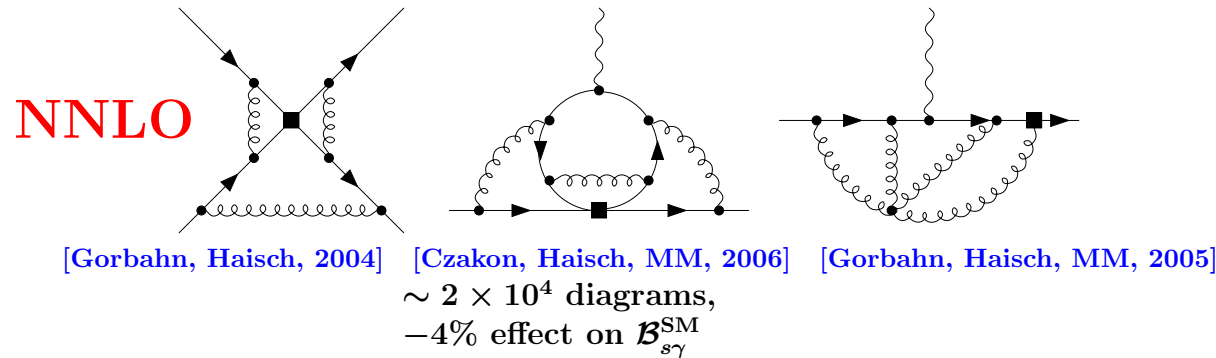
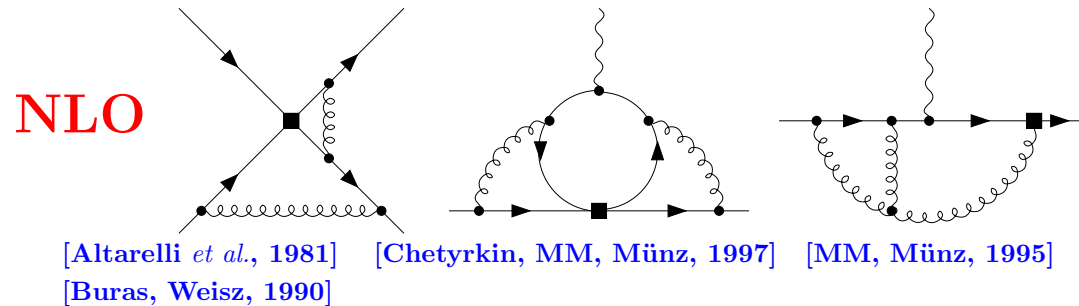
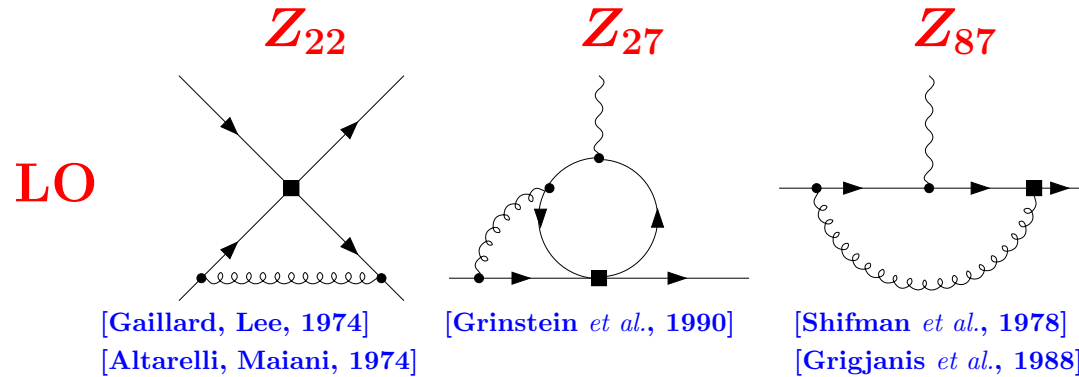
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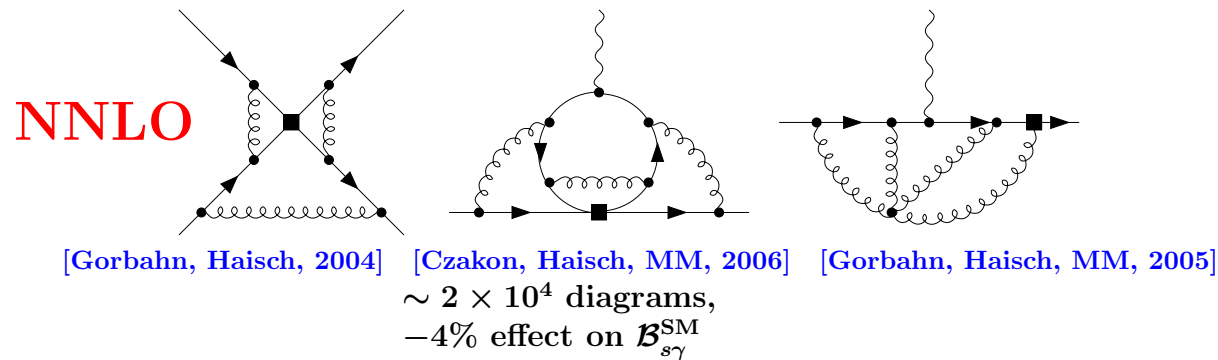
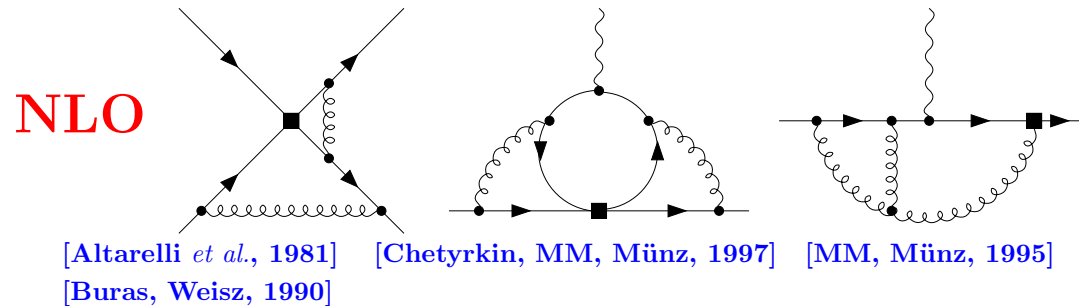
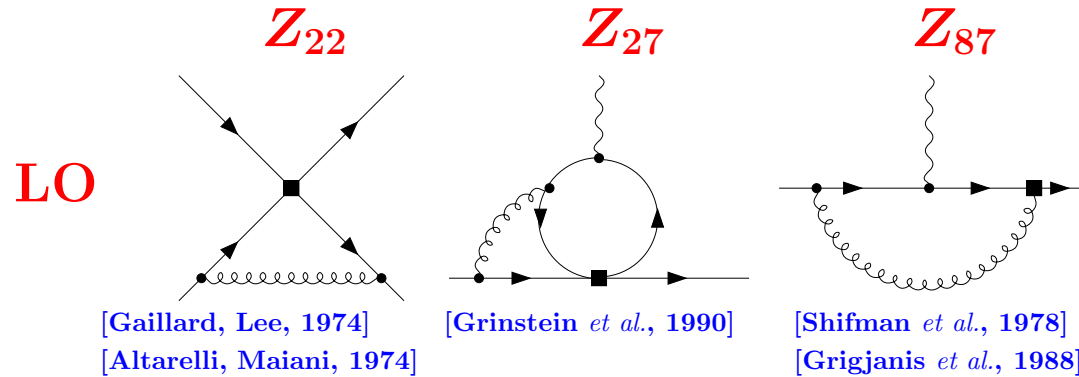
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Since 2006, all the Wilson coefficients $C_1(\mu_b), \dots, C_8(\mu_b)$ are known at the NNLO.