

Freezing-in Cannibal DM during early matter domination

Based on: [arXiv:2506.09155](https://arxiv.org/abs/2506.09155)

Esau Cervantes

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Bernal and **Kuldeep Deka**

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Scalars 2025: Higgs bosons and cosmology



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Cannibal Dark Matter

SELF-INTERACTING DARK MATTER

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AND

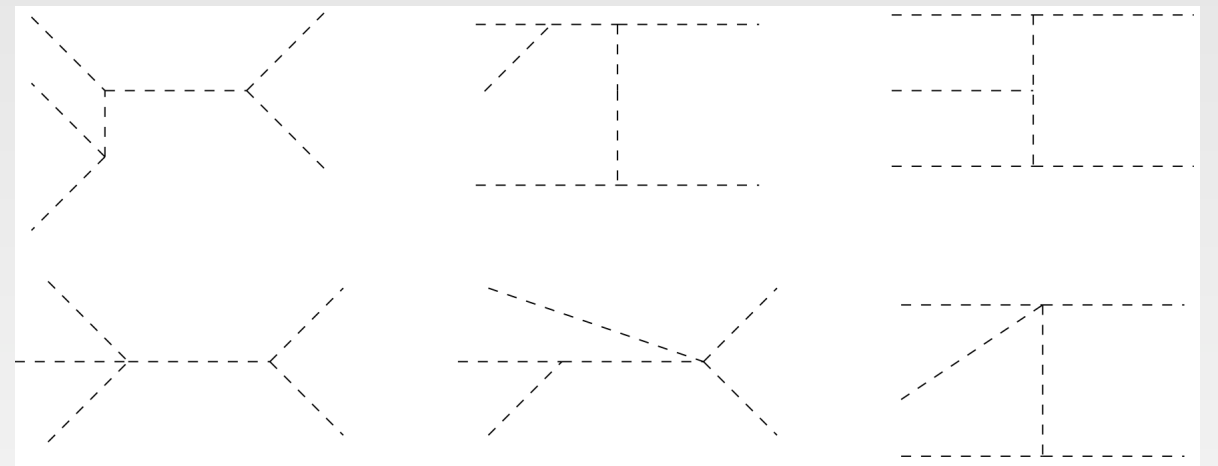
LAWRENCE J. HALL

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Simple realisation with a scalar

field: $\frac{g}{3!}\phi^3 + \frac{\lambda}{4!}\phi^4$



If DM is non-relativistic, $\Gamma_{3 \rightarrow 2} > \Gamma_{2 \rightarrow 3}$. The DM fluid **exchanges** particle number for kinetic energy!



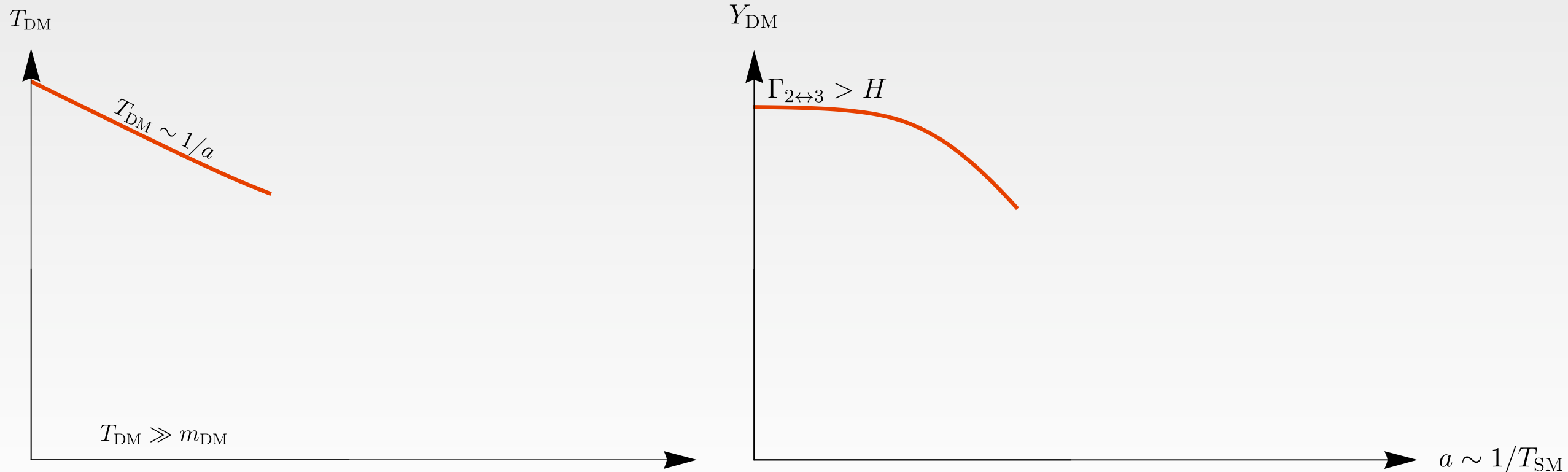
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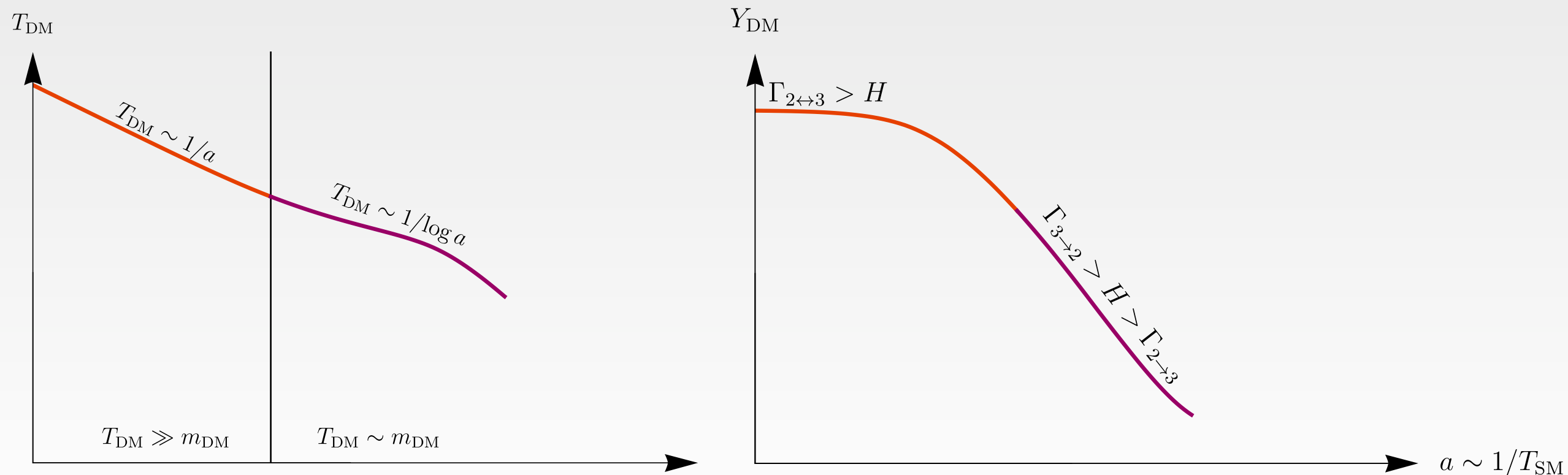
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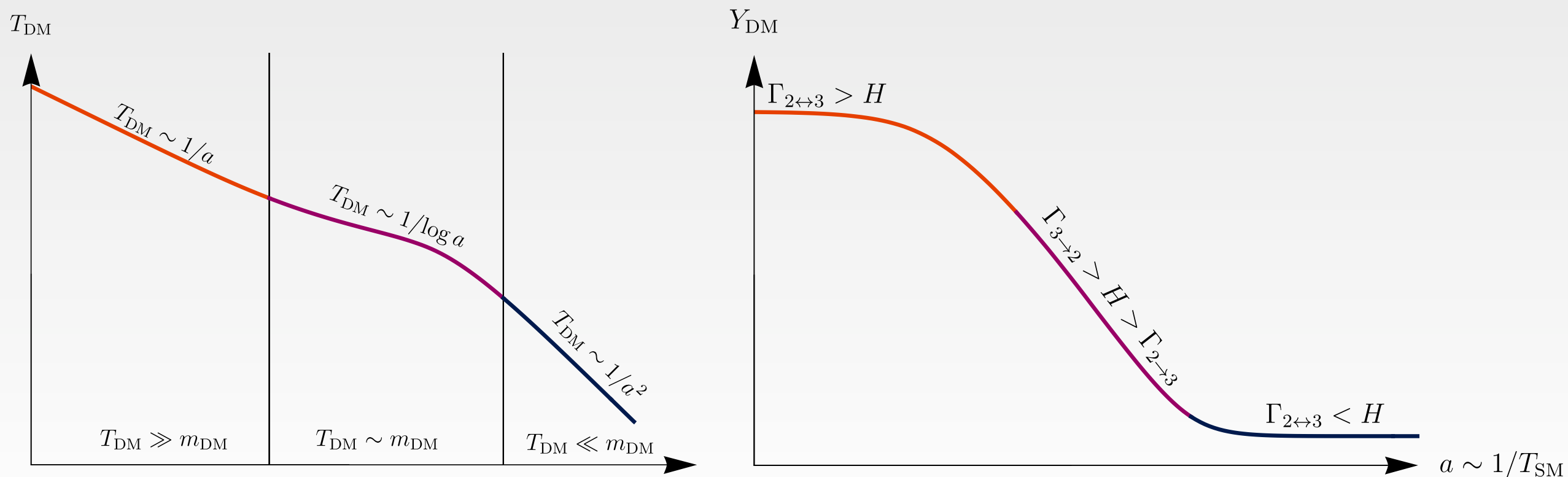
- DM is initially *relativistic*;
- as the DM fluid cools down, the dark sector *exchanges* number of particles for kinetic energy;



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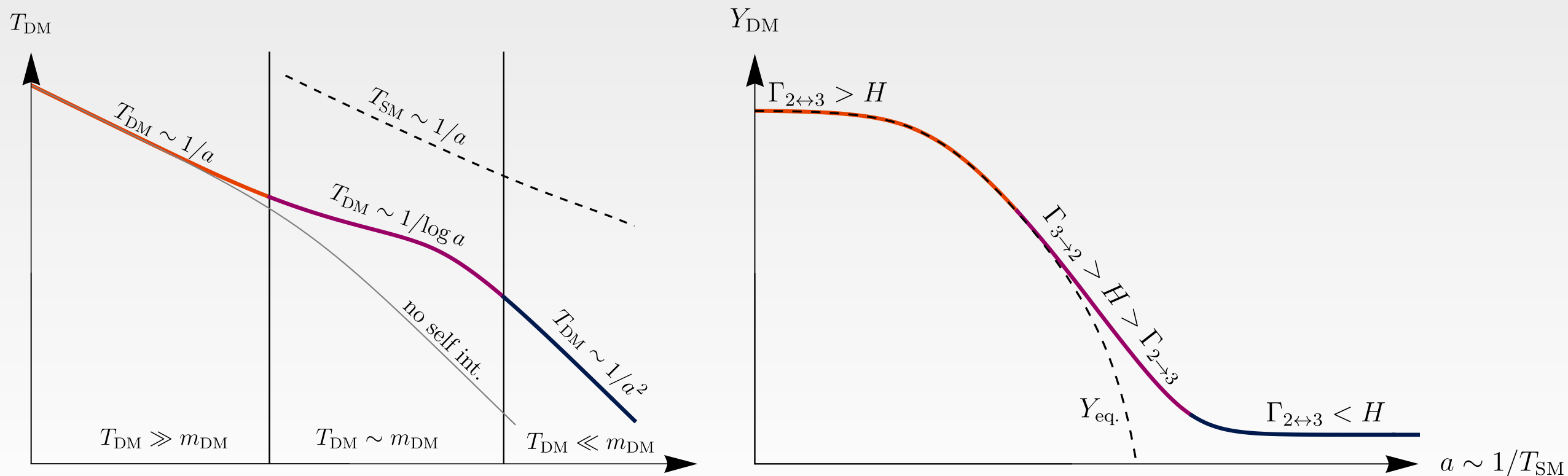
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- as the DM fluid cools down, the dark sector *exchanges* number of particles for kinetic energy;
- all interactions decouple and the system behaves as a non-relativistic gas.



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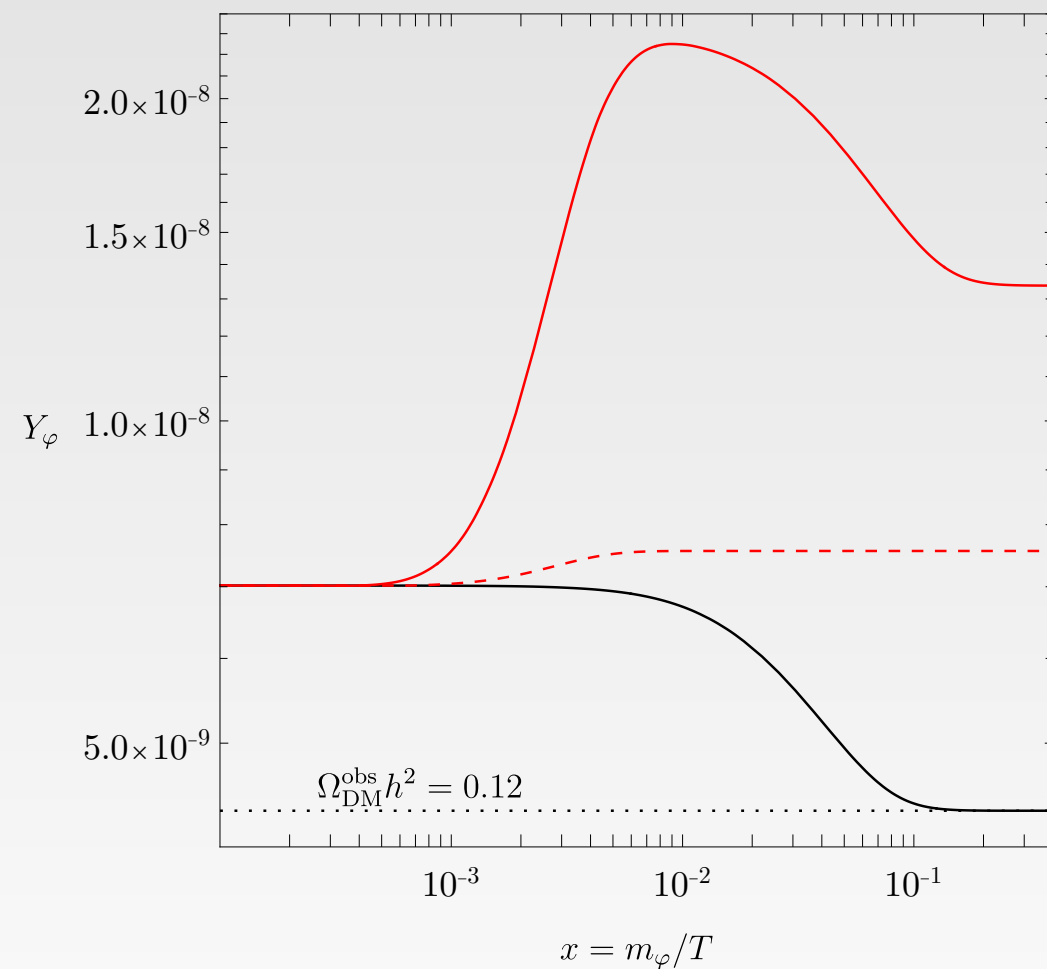
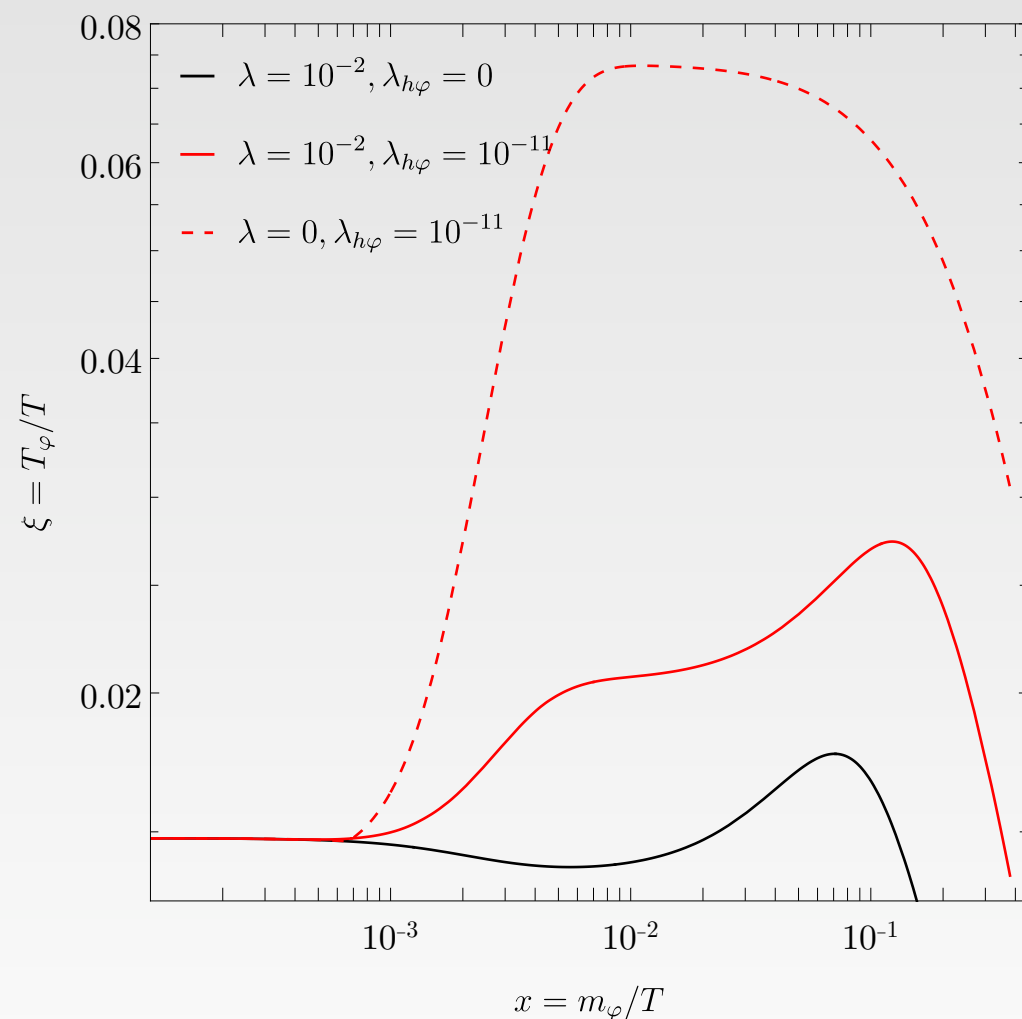
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Cannibals produced via freeze-in

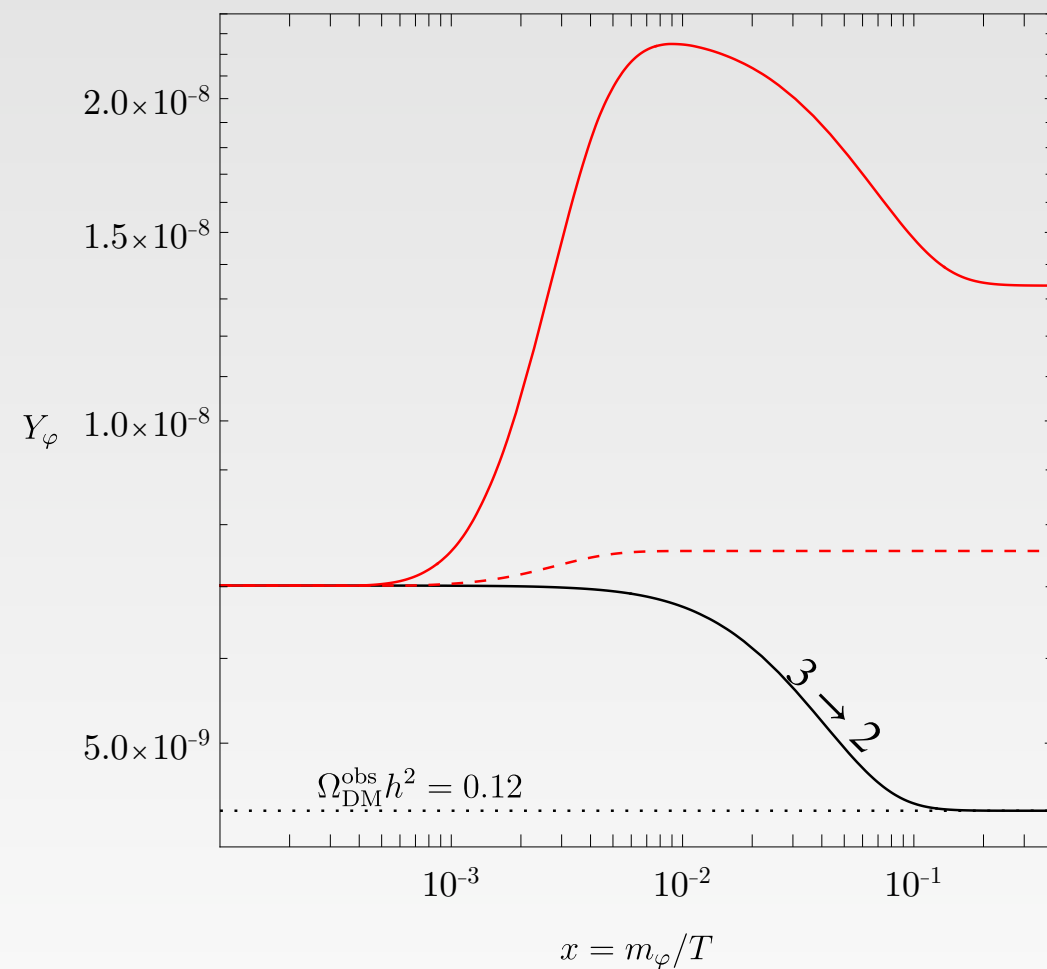
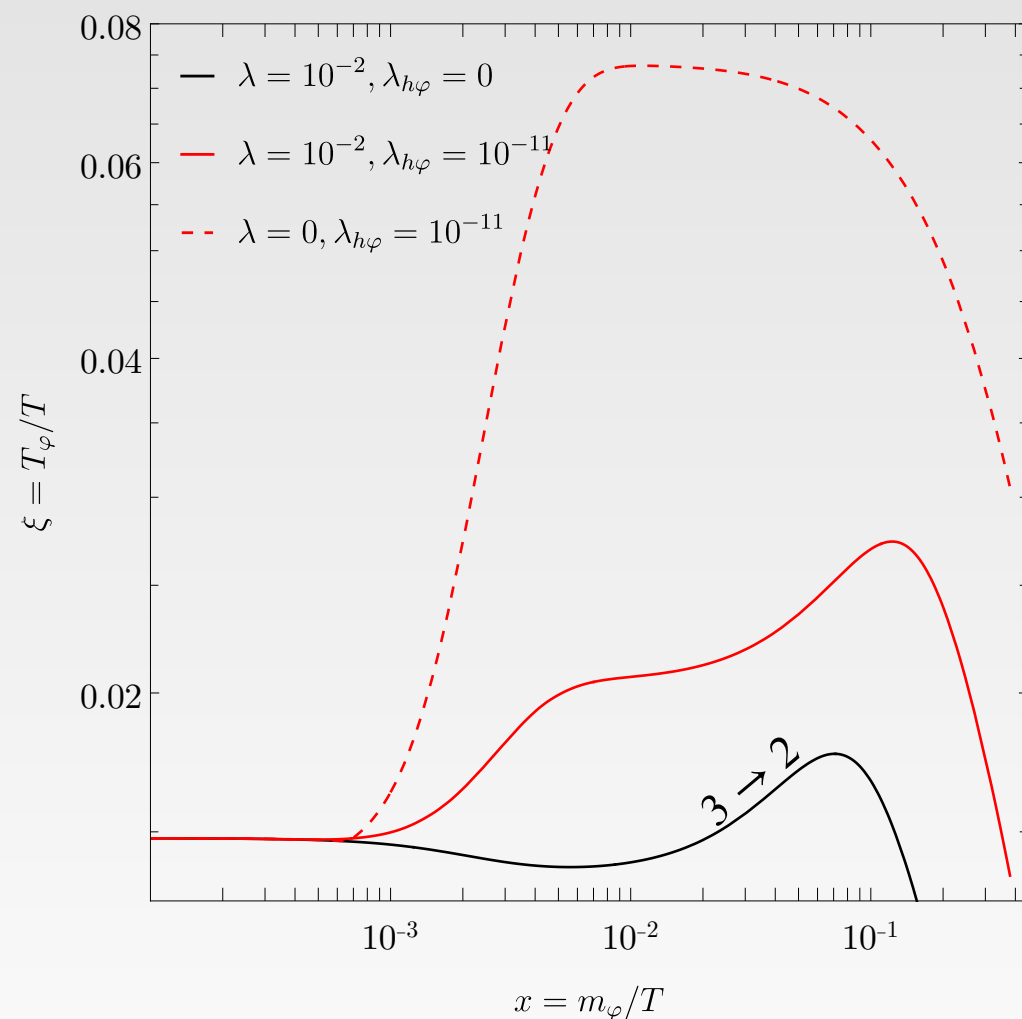
Consider $\mathcal{L} \supset -\lambda_{h\varphi}\varphi^2 H^\dagger H$, $\lambda_{h\varphi} \ll 1$, $\lambda_\varphi \geq 10^{-4}$ and initially cold DM; $T_{DM}/T_{SM} = 10^{-2}$:



See EC, A. Hryczuk 24

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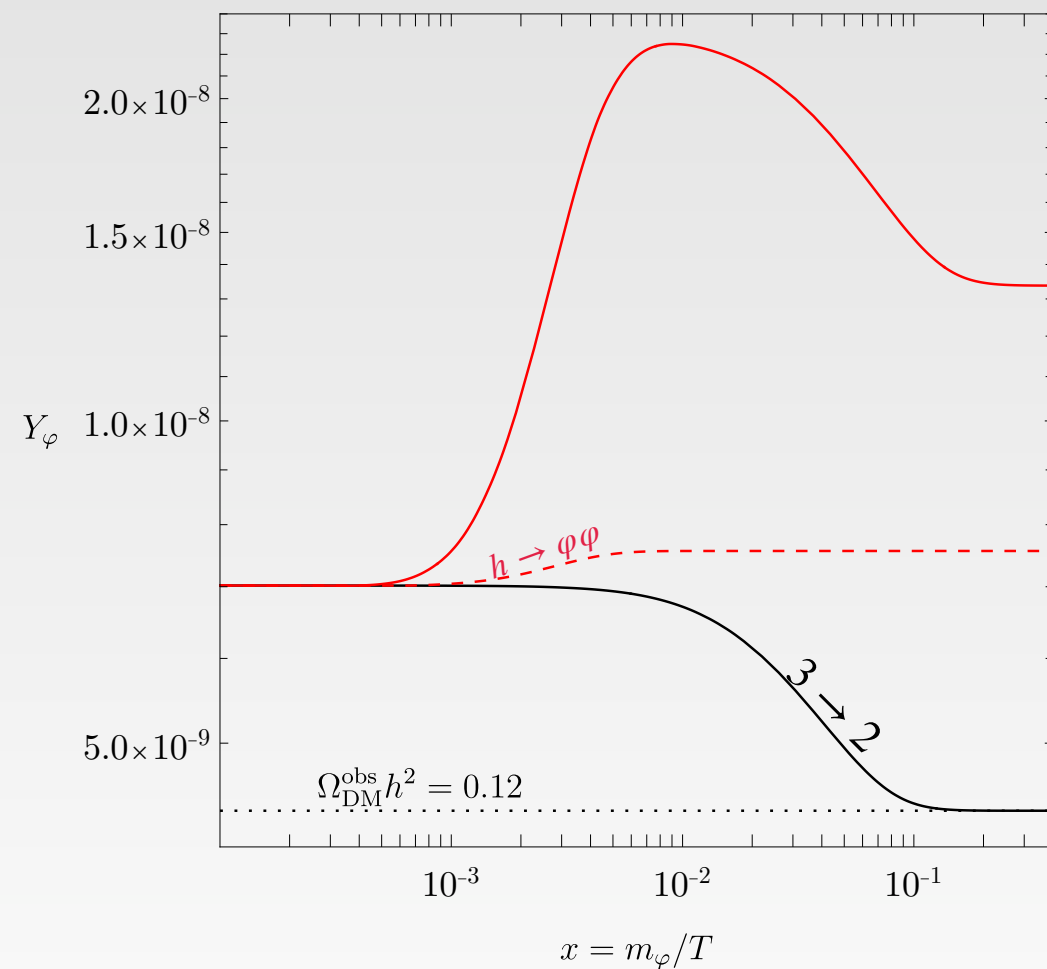
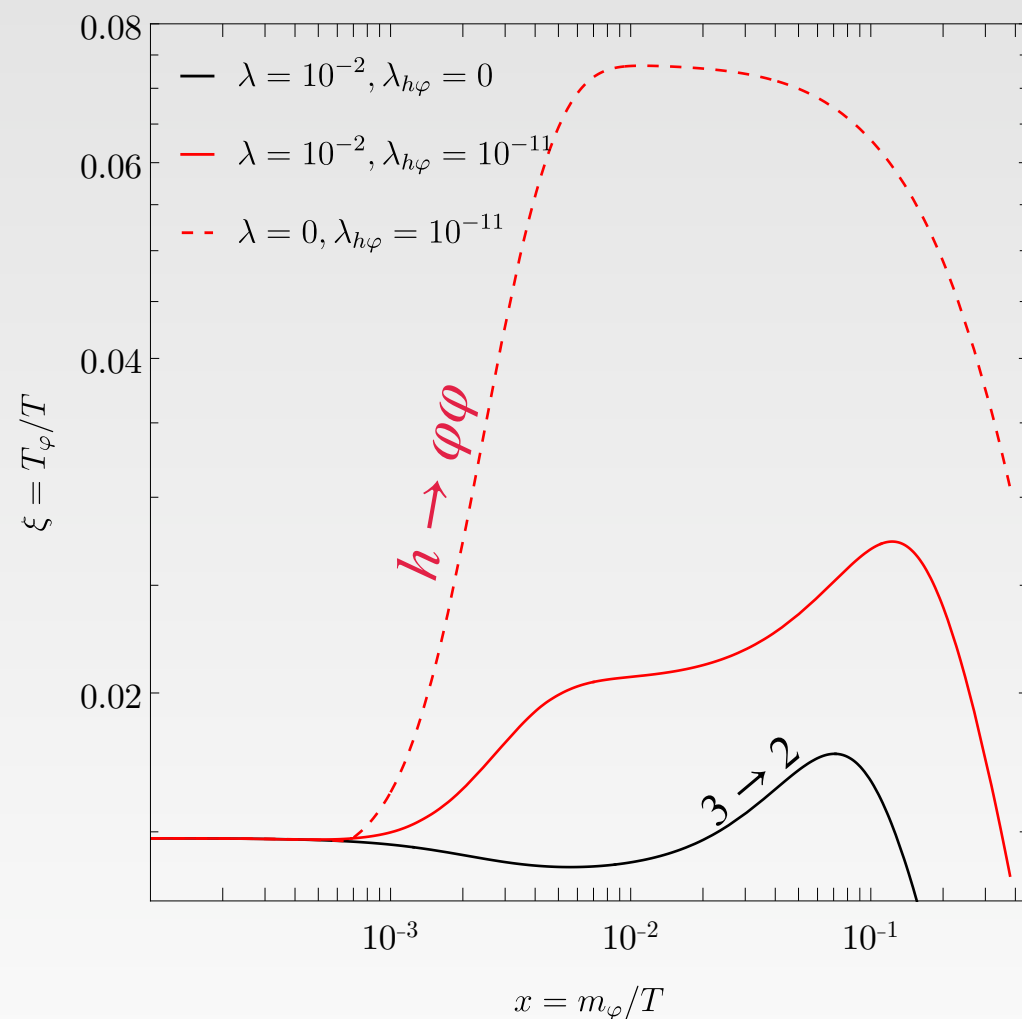
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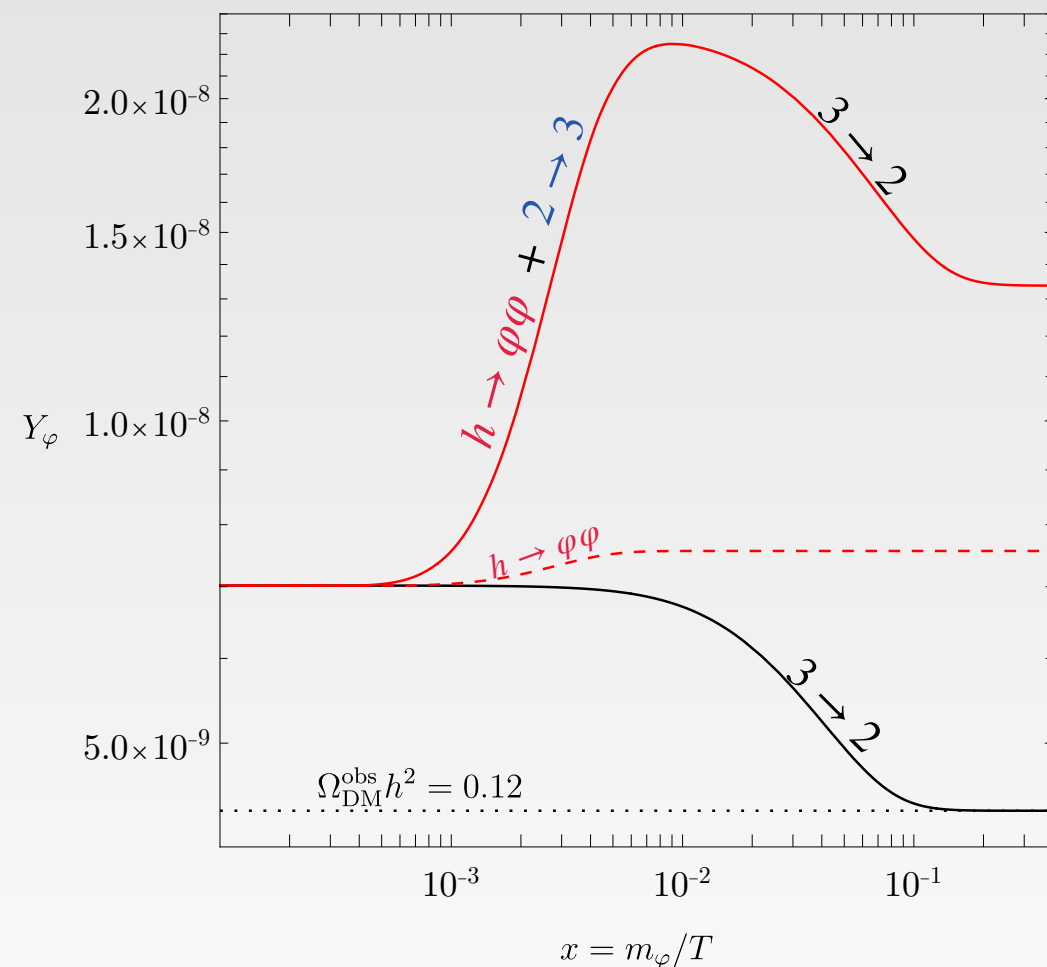
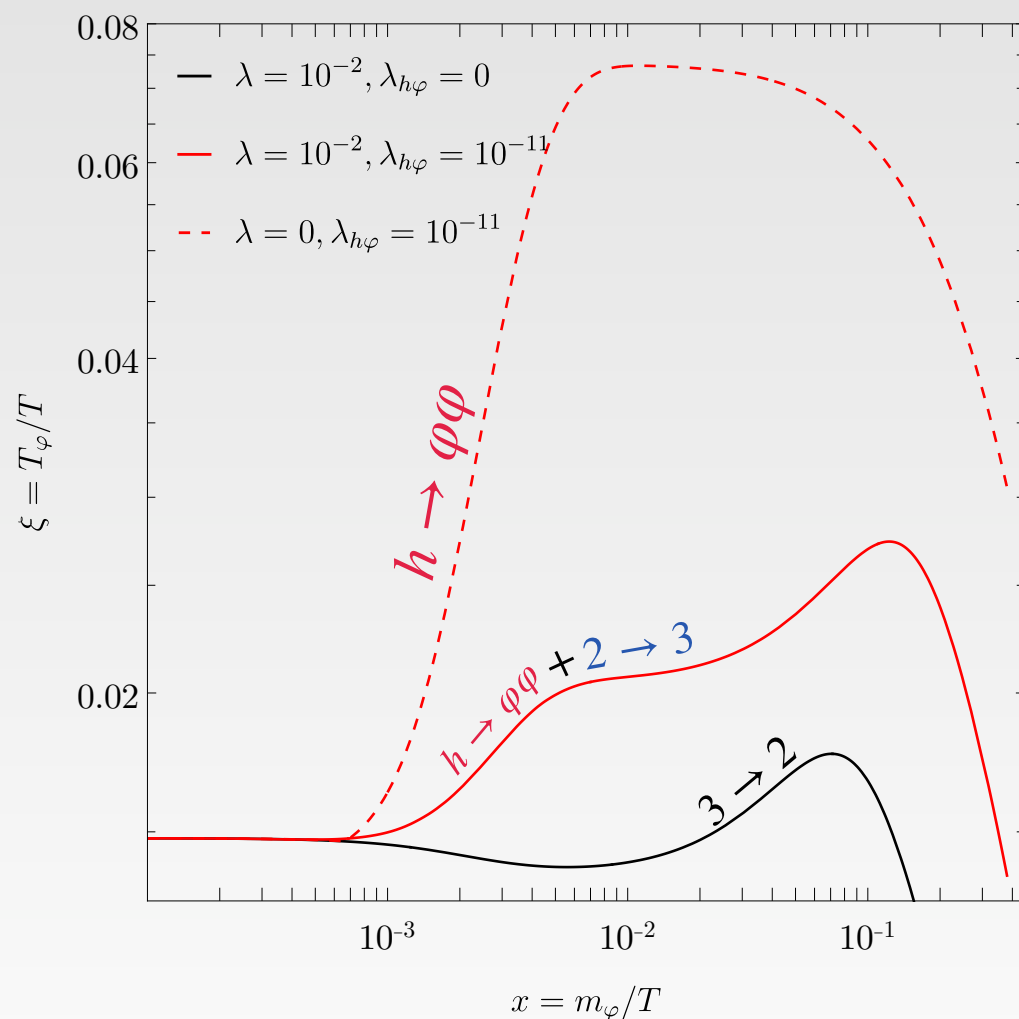
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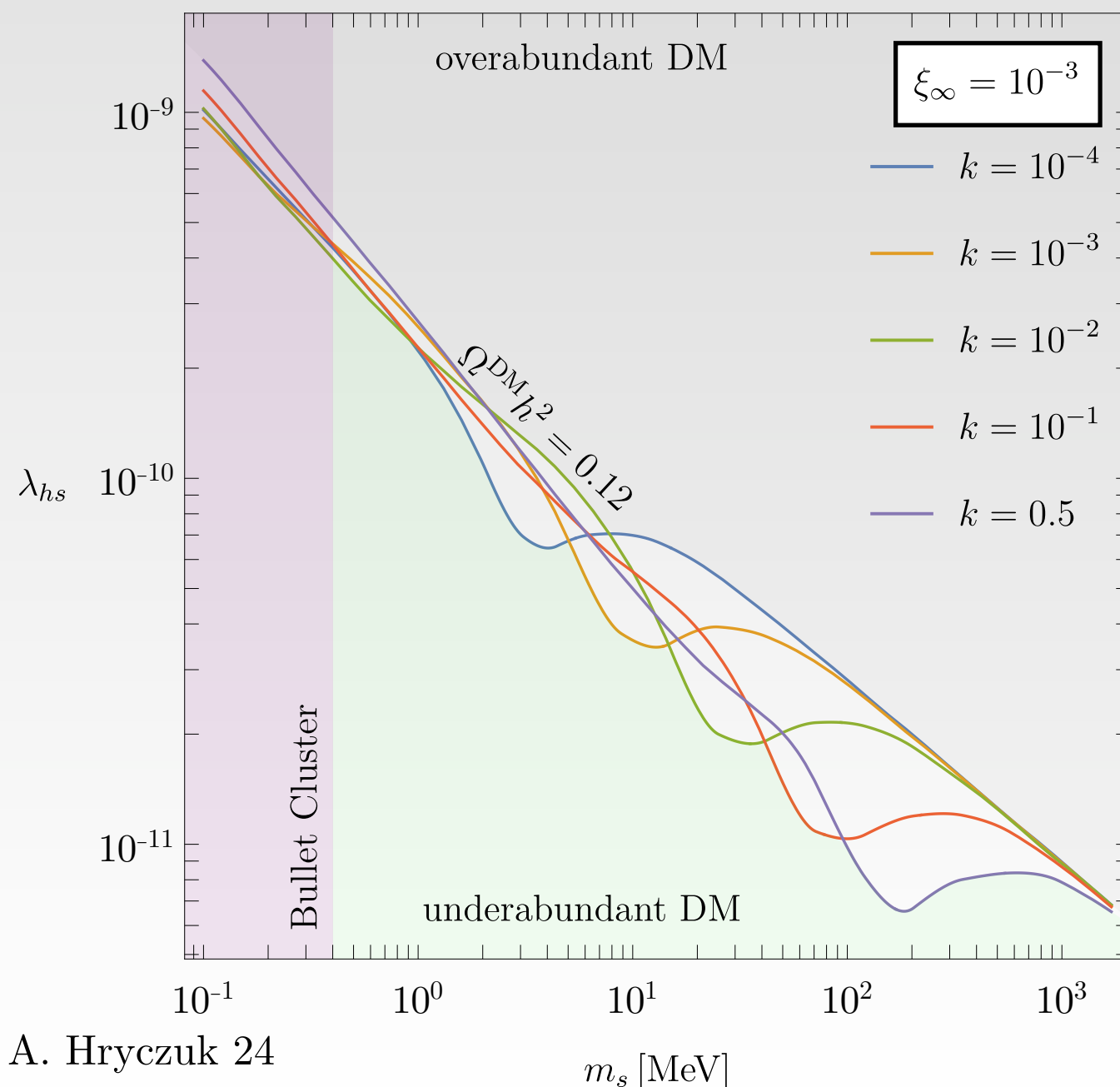
and also Bernal et.al 15 (SIMP \mathbb{Z}_2 DM)

Cannibals produced via freeze-in

Toy model: $\mathcal{L} \supset -\frac{1}{3!}g_s(S^3 + (S^*)^3) - \frac{\lambda_s}{4}|S|^4 - \lambda_{hs}|S|^2|H|^2$

DM self interactions (cannibal)

Portal



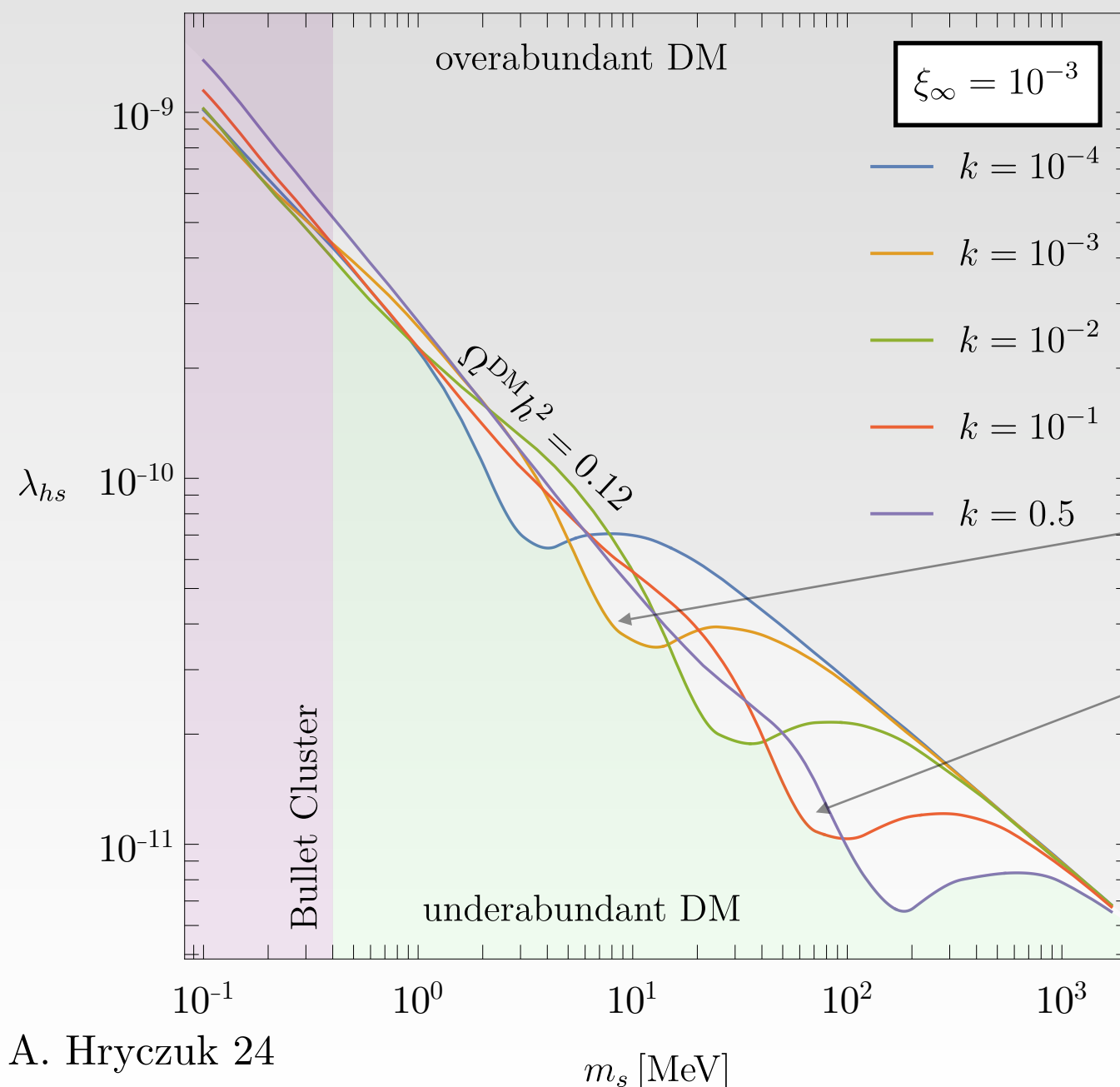
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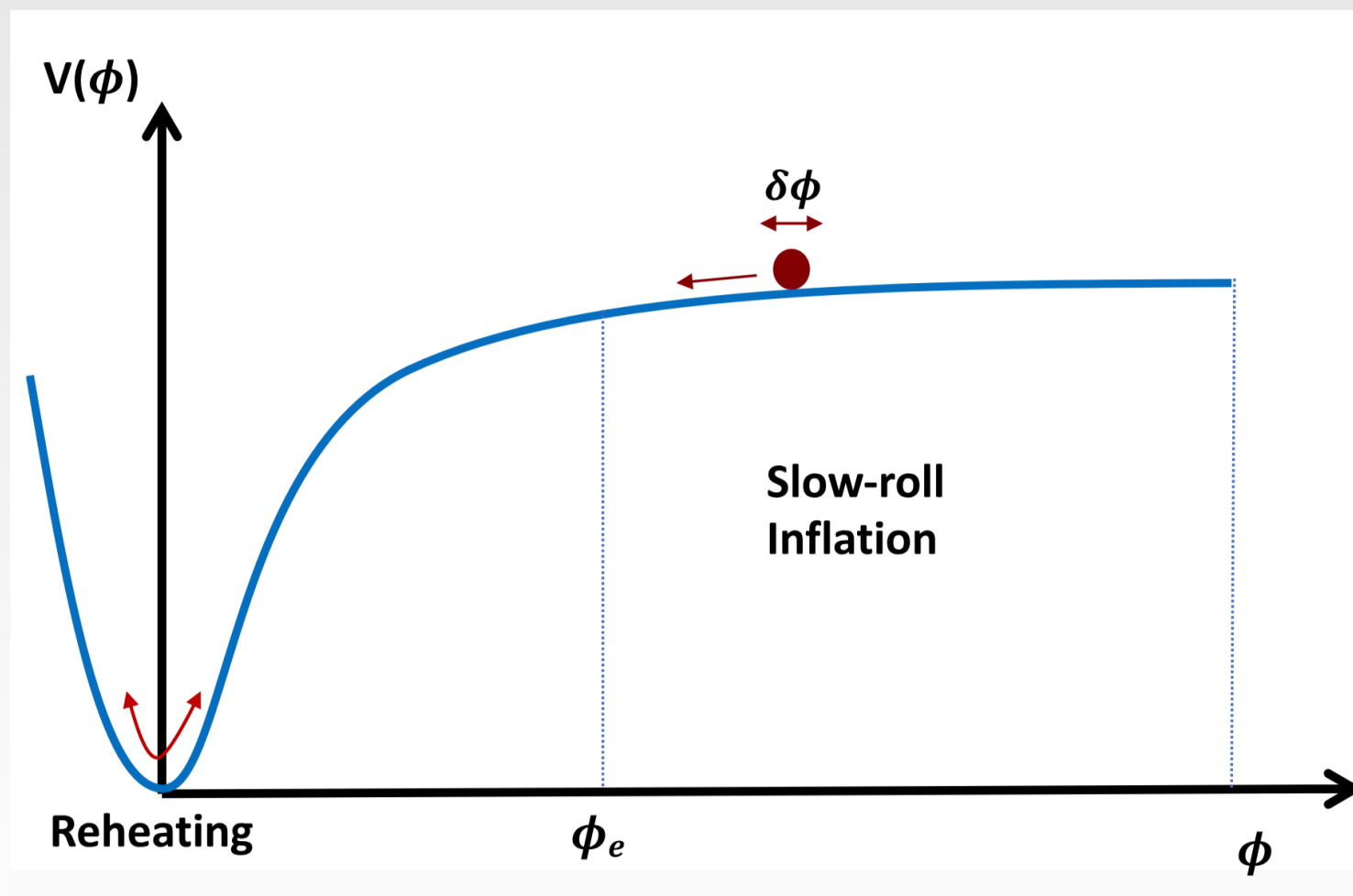


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2 → 3 at play

Inflaton decay and reheating

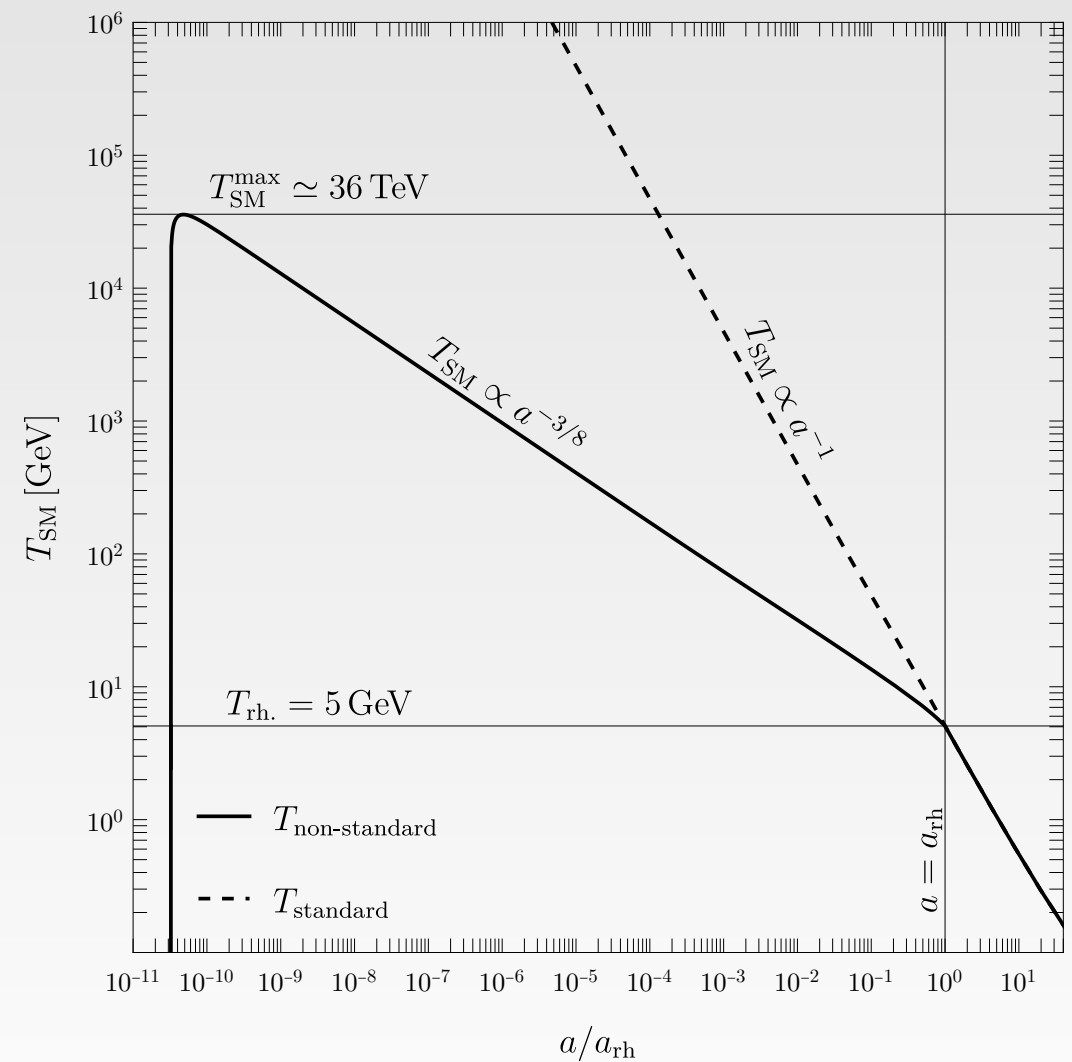
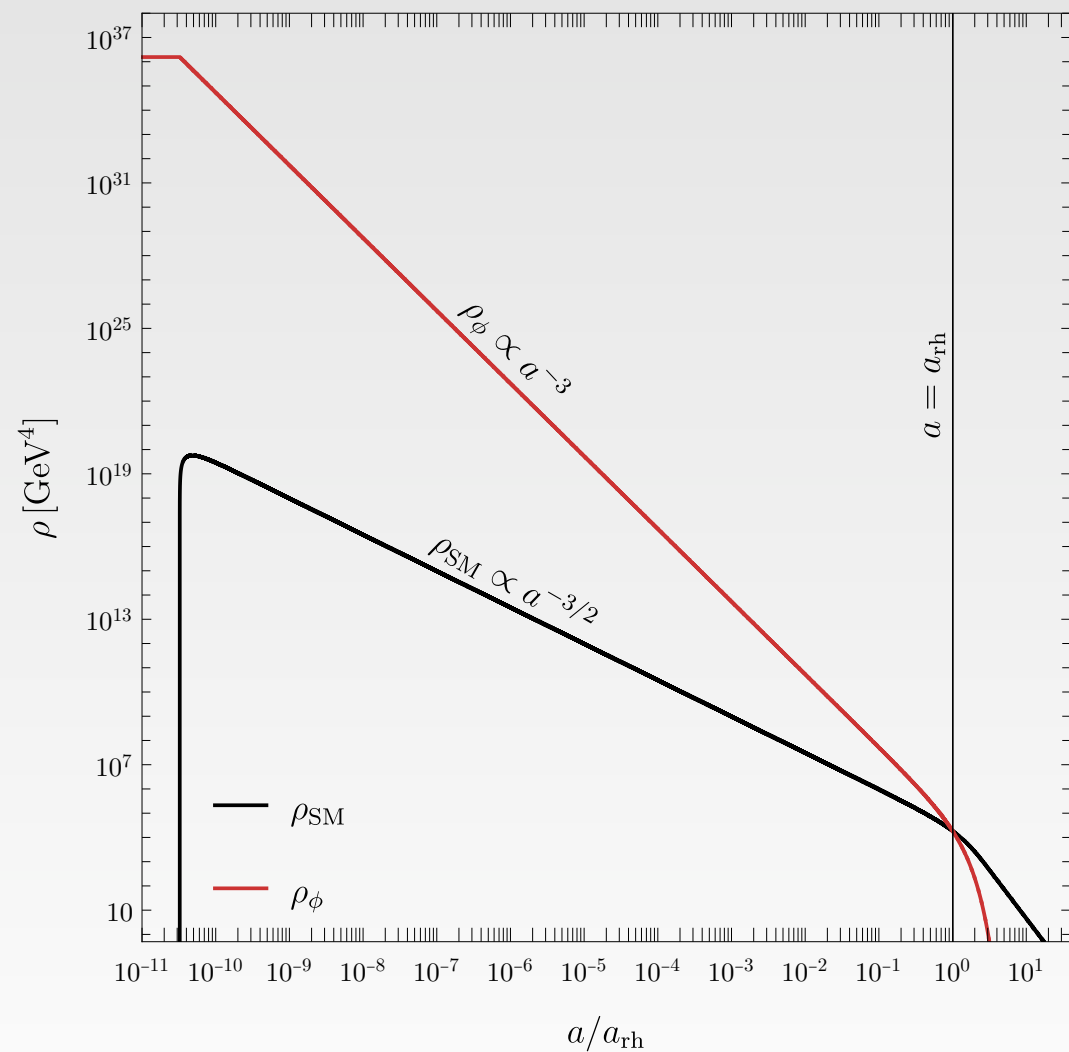
Transition between *matter* domination and *radiation* can be due to a scalar (**inflaton**) field ϕ that rolls ($a \propto e^{Ht}$) in the potential and subsequently oscillates in the minimum decaying into SM states.



$$\frac{d\rho_\phi}{dt} + 3H\rho_\phi = -\Gamma\rho_\phi,$$
$$\frac{d\rho_R}{dt} + 4H\rho_R = +\Gamma\rho_\phi,$$

Inflaton decay and reheating

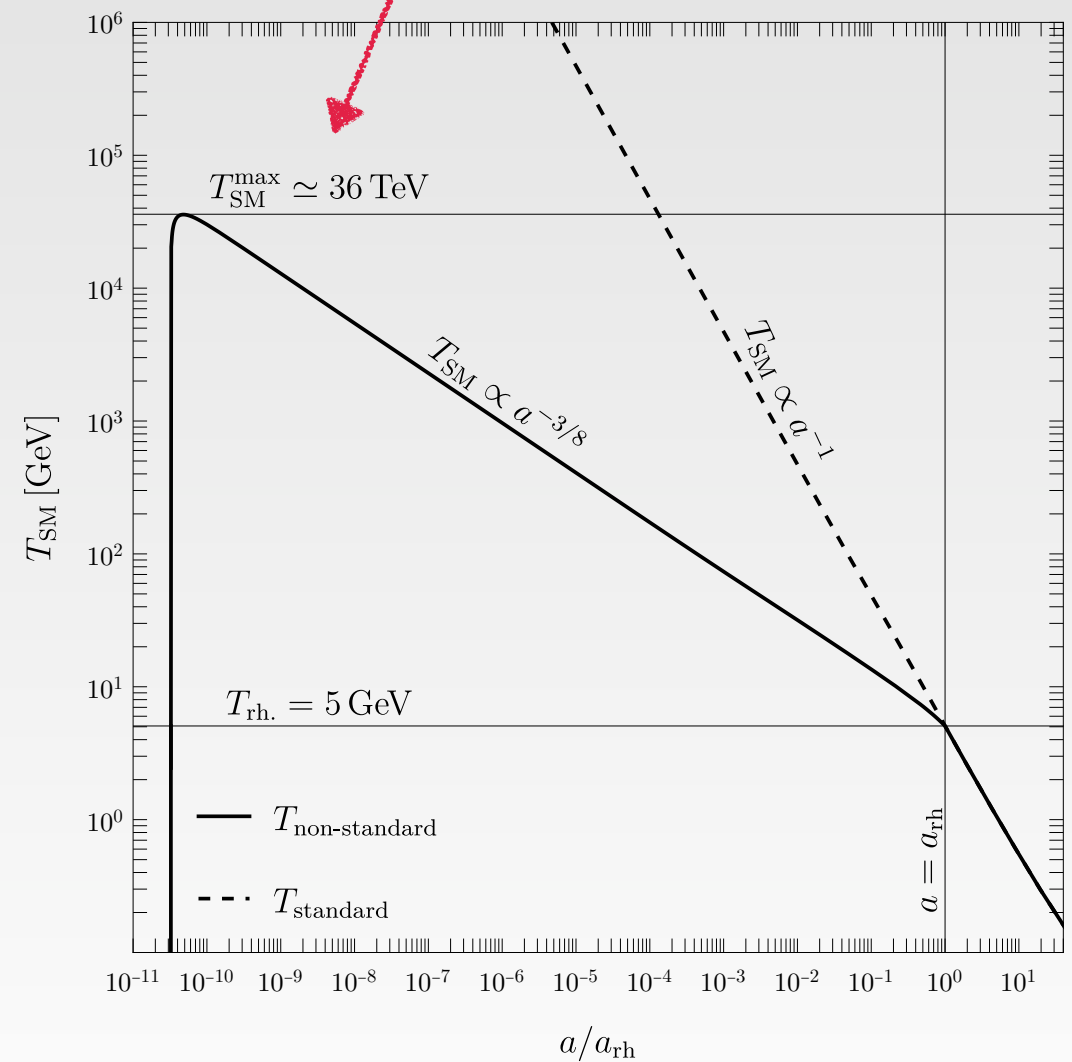
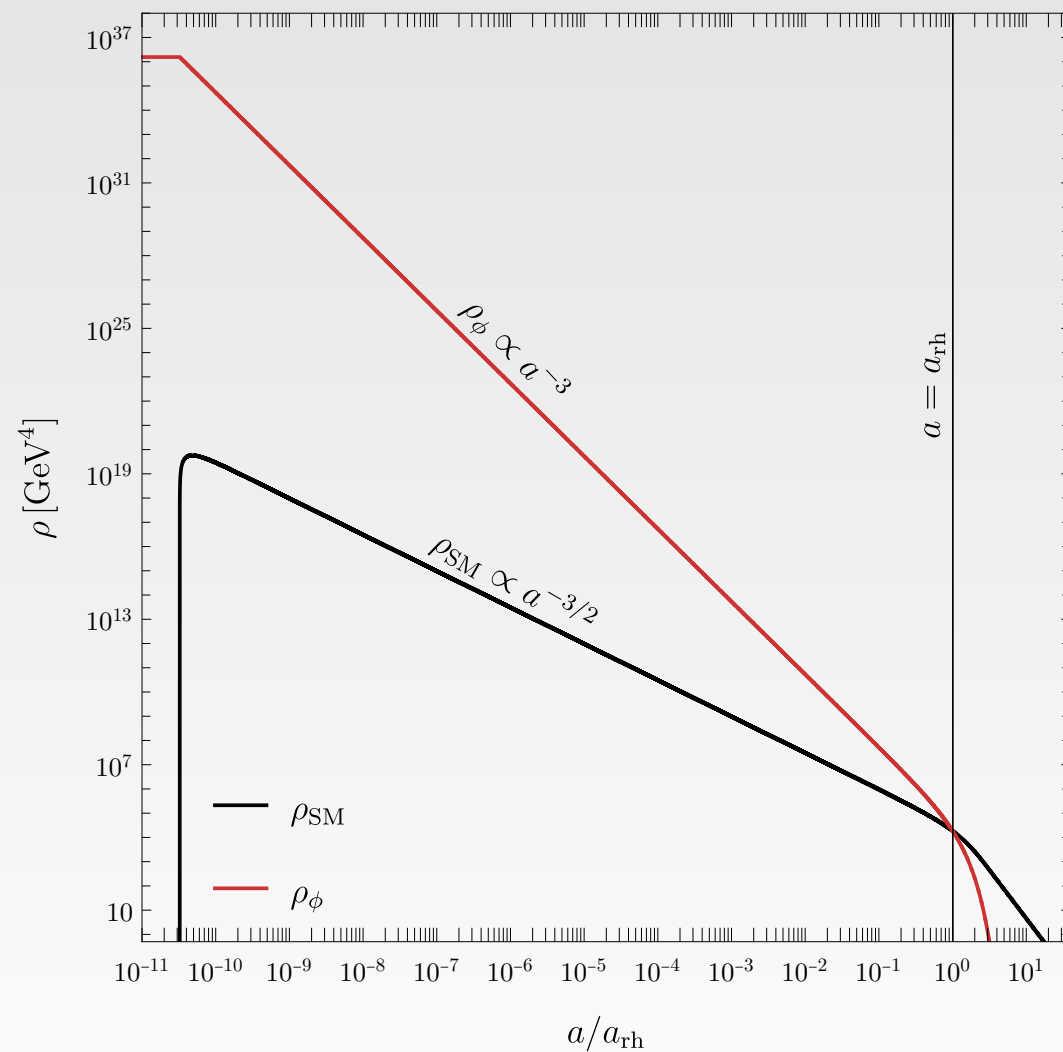
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To produce *all* SM species

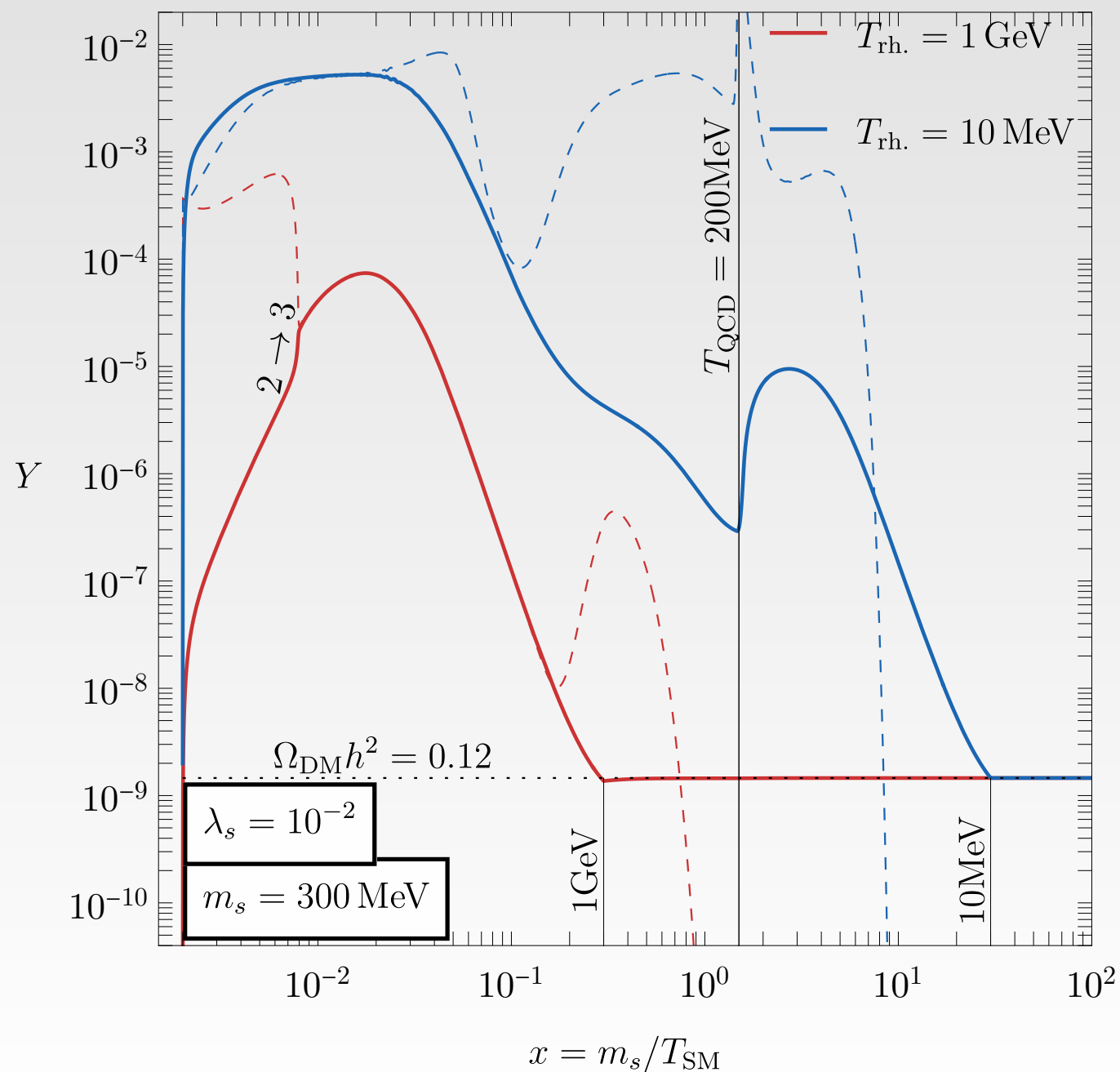


Production during reheating

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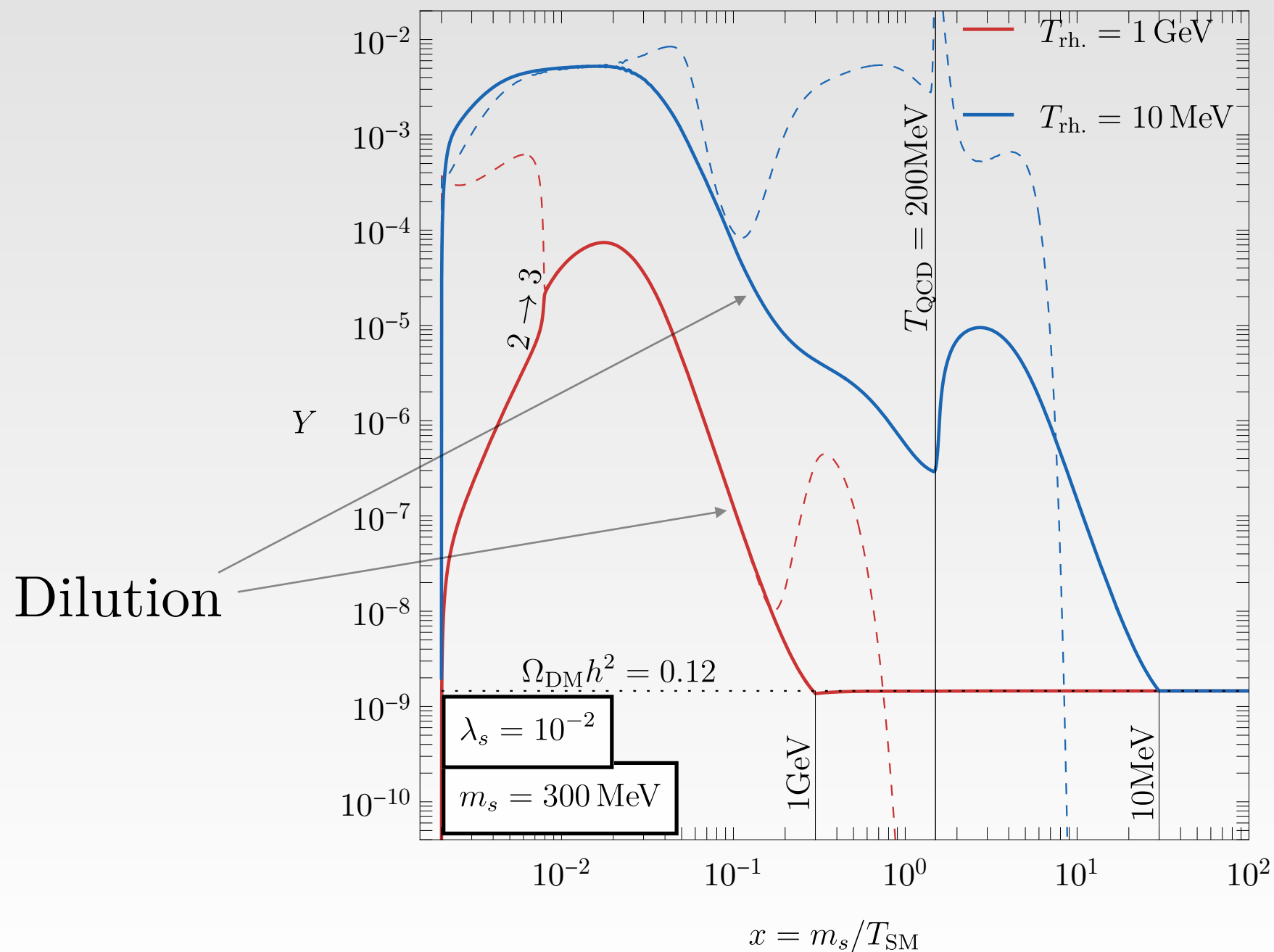


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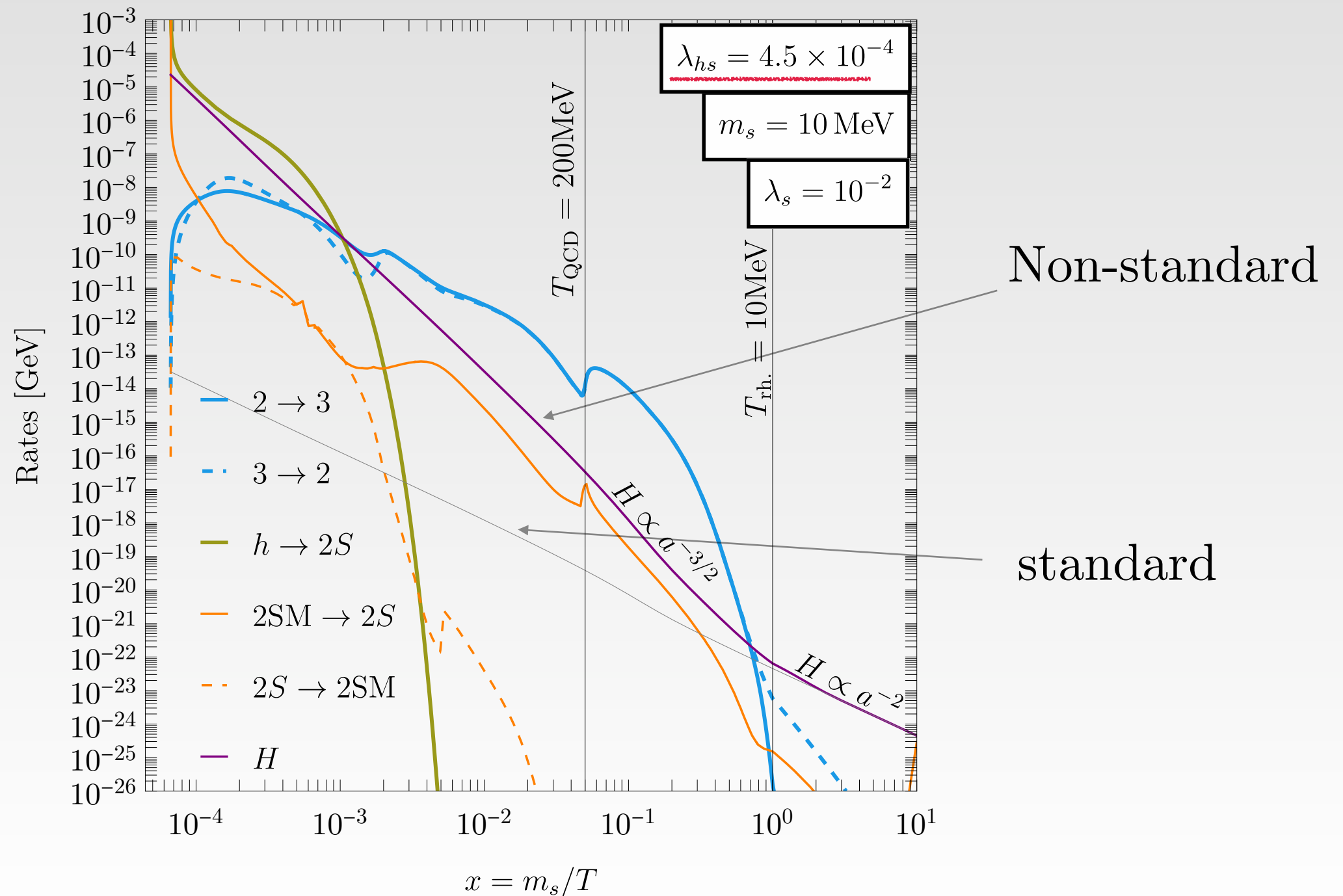
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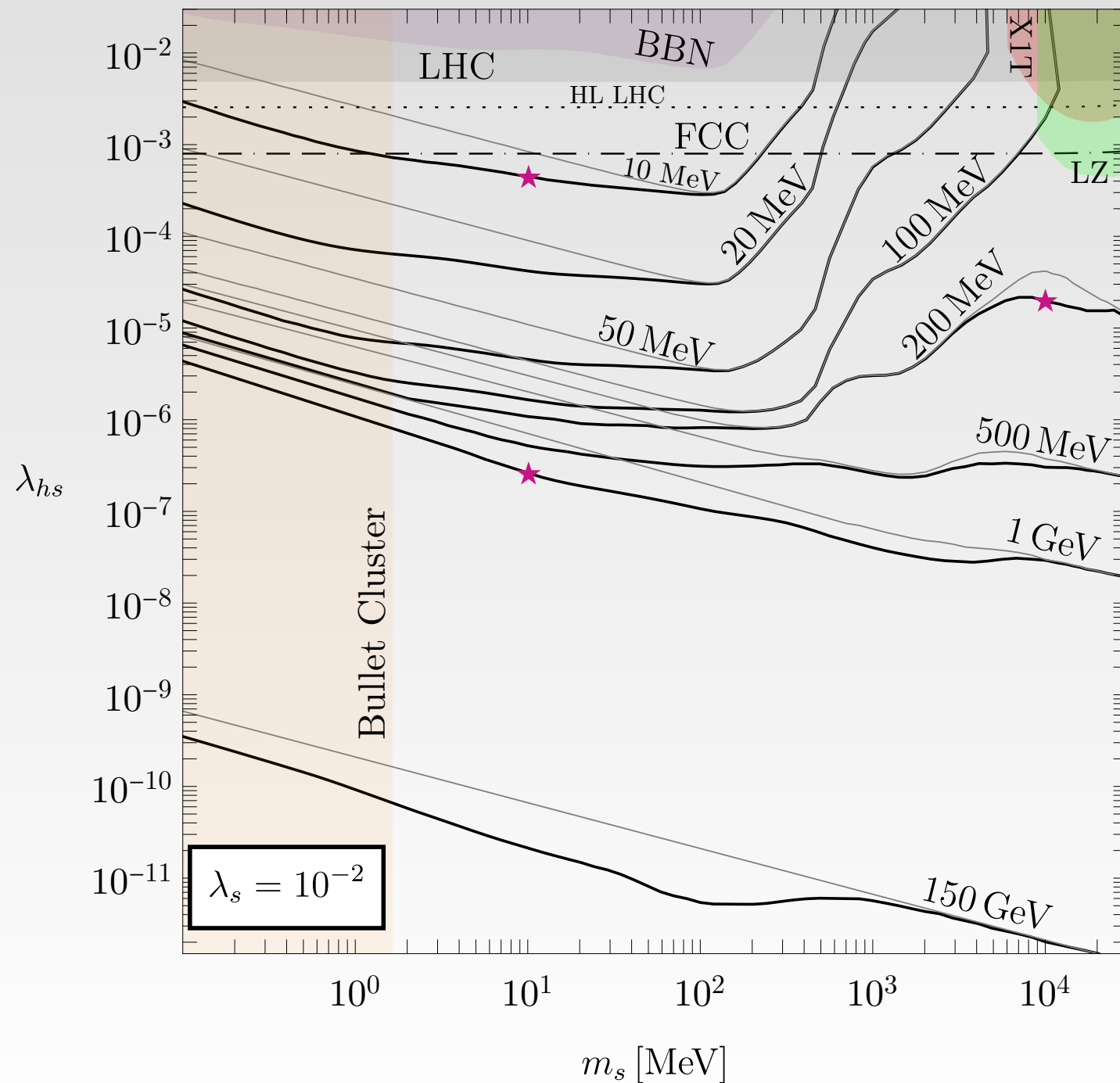


Production during reheating

Production rate from SM has to catch up with $H \propto T^4$, and ρ_{DM} dilutes during reheating.

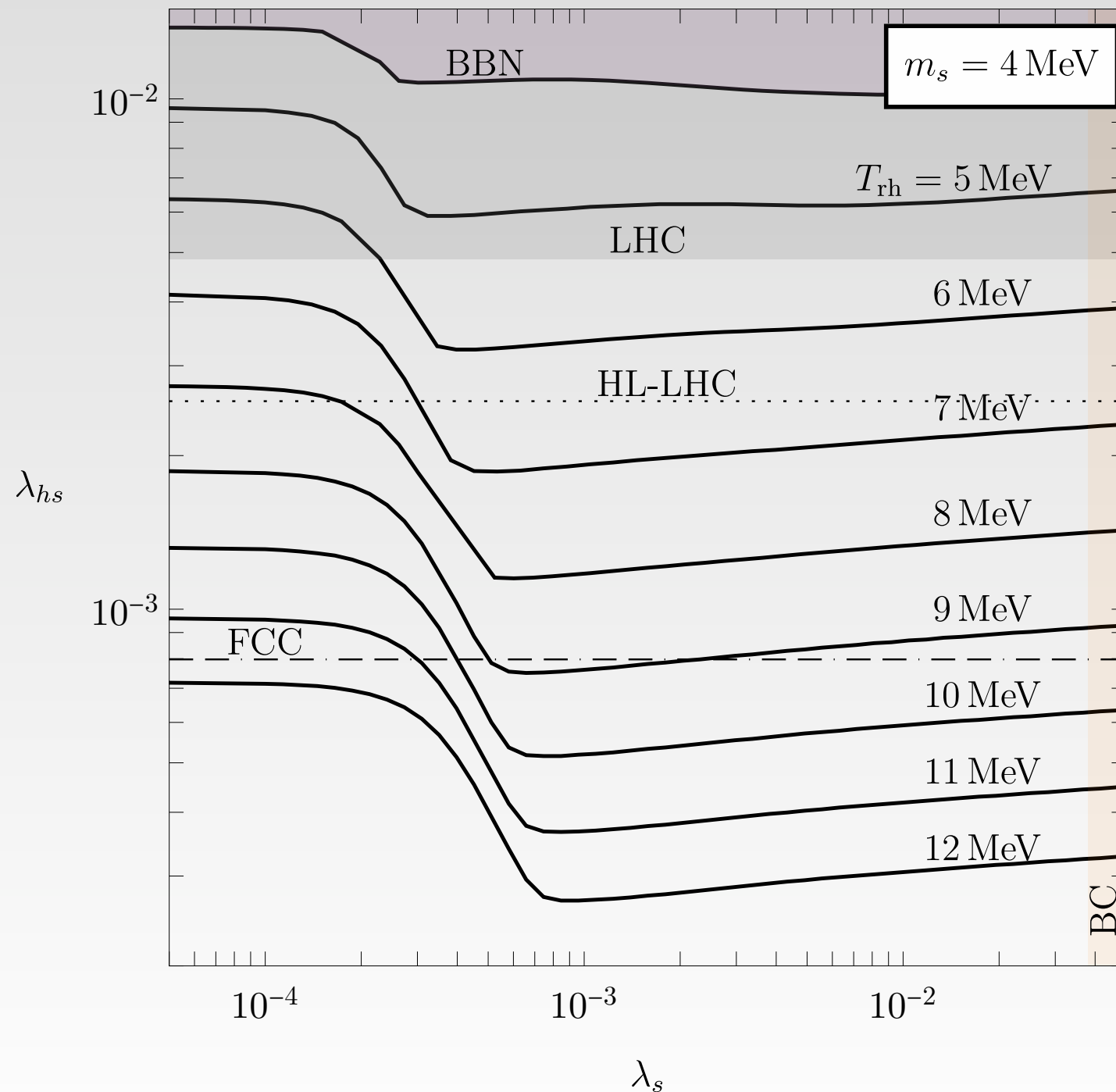


Impact on collider phenomenology



- Low T_{rh} leads to detectability;
- The case of instantaneous reheating is studied in Lebedev, Morais, Oliveira, Pasechnik 24.

Self interactions with low T_{rh}



- $T_{rh} = 6$ MeV is either **excluded** or **detectable** depending on λ_s ;
- $T_{rh} = 11$ MeV is either **out of reach** or **detectable** depending on λ_s ;
- The peculiar behaviour of the curves is due to the $2 \rightarrow 3$ reaction overproducing DM.

Summary

- SIDM produced via the freeze-in mechanism has a unique evolution in the Early Universe;
- Temperature can have a **non-trivial** impact in such scenarios and **need to be studied** carefully;
- Non-standard cosmologies might be able to test SIDM.

Summary

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- Temperature can have a **non-trivial** impact in such scenarios and **need to be studied** carefully;
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Dziękuję bardzo!

Coupled Boltzmann equations

From the fBE we can obtain a ‘temperature’ Boltzmann equation:

We define $T' := \frac{g_{dm}}{3n} \int \frac{d^3p}{(2\pi)^3} \frac{p^2}{E} f(p);$

we integrate $g(2\pi)^{-3} \int d^3p \frac{p^2}{E} (\partial_t - H\vec{p} \cdot \vec{\nabla}_p) f = g(2\pi)^{-3} \int d^3p \frac{p^2}{E} C[f] =: C_2;$

to obtain $\frac{dT'}{da} = -\frac{2T'}{a} + \frac{1}{3a} \left\langle \frac{p^4}{E^3} \right\rangle + \frac{a^2}{3HN} C_2 - \frac{a^2 T'}{HN} C_0;$

along with the usual nBE: $\frac{dN}{da} = \frac{a^2}{H} g \int \frac{d^3p}{(2\pi)^3} C[f] =: \frac{a^2}{H} C_0, N = na^3;$

we close the system by assuming $f(E, T') = \frac{n}{n_{\text{eq}}} \exp \left[-\frac{E}{T'} \right].$