

Gravitational Waves from a First-Order Phase Transition of the Inflaton

Jörn Kersten



연세대학교
YONSEI UNIVERSITY



UNIVERSITY OF BERGEN

Based on JK, Seong Chan Park, Yeji Park, Juhoon Son, Liliana Velasco Sevilla,
JCAP **04** (2025) 053 [arXiv:2412.17278]

A Window to the Very Early Universe?

- Pulsar Timing Arrays

→ Evidence for **Stochastic GW Background**

NANOGrav, ApJL 951 (2023)

EPTA, InPTA, A&A 678 (2023)

Parkes PTA, ApJL 951 (2023)

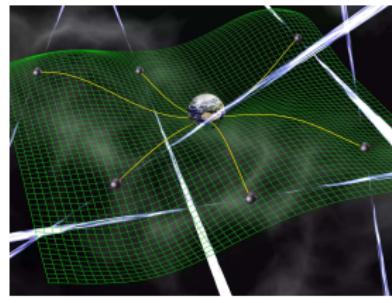
CPTA, RAA 23 (2023)

- Lower frequency than LIGO/Virgo events

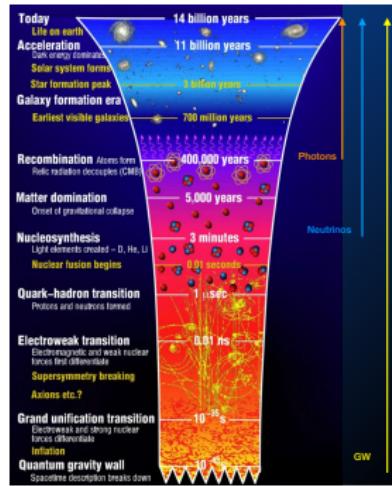
→ Mergers of **supermassive** black holes?

- More interesting: particle physics origin

- First-Order Phase Transition (FOPT)
- Cosmic strings or domain walls
- Inflation



D. Champion/MPI for Radio Astronomy



Higgs Inflation

Bezrukov & Shaposhnikov, PLB 659 (2008)

$$\mathcal{L}_J = \sqrt{-g_J} \left[\frac{M_P^2}{2} R_J + \xi \phi^\dagger \phi R_J + g_J^{\mu\nu} (D_\mu \phi)^\dagger (D_\nu \phi) - V_J(\phi) + \dots \right]$$

- J: Jordan frame
- R: Ricci scalar
- V_J : SM Higgs potential
- ξ : non-minimal coupling to gravity
 - Consistent with all symmetries
 - Required for renormalization in curved spacetime
 - Not necessarily small

Higgs Inflation

Bezrukov & Shaposhnikov, PLB 659 (2008)

$$\mathcal{L}_J = \sqrt{-g_J} \left[\frac{M_P^2}{2} R_J + \xi \phi^\dagger \phi R_J + g_J^{\mu\nu} (D_\mu \phi)^\dagger (D_\nu \phi) - V_J(\phi) + \dots \right]$$

- J: Jordan frame
- R: Ricci scalar
- V_J : SM Higgs potential
- ξ : non-minimal coupling to gravity
 - Consistent with all symmetries
 - Required for renormalization in curved spacetime
 - Not necessarily small
- Def.: conformal factor $\Omega^2(\phi) \equiv 1 + 2\xi \frac{\phi^\dagger \phi}{M_P^2}$

$$\mathcal{L}_J = \sqrt{-g_J} \left[\frac{M_P^2}{2} \Omega^2(\phi) R_J + g_J^{\mu\nu} (D_\mu \phi)^\dagger (D_\nu \phi) - V_J(\phi) + \dots \right]$$

Higgs Inflation

- Unitary gauge: $\phi = \begin{pmatrix} 0 \\ (\varphi + v)/\sqrt{2} \end{pmatrix}$
- Weyl transformation to **Einstein frame**: $g_J \rightarrow g_E \equiv \Omega^2(\varphi) g_J$
- Canonical normalization: $\chi = \int_0^\varphi d\varphi \sqrt{\frac{3}{2} \frac{(\Omega^2, \varphi)^2}{\Omega^4} + \frac{1}{\Omega^2}}$

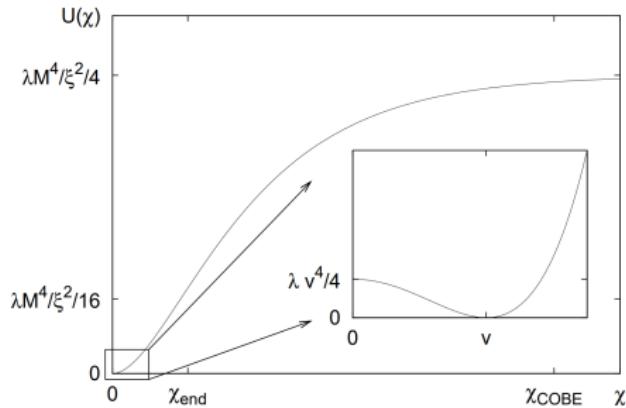
Higgs Inflation

- Unitary gauge: $\phi = \begin{pmatrix} 0 \\ (\varphi + v)/\sqrt{2} \end{pmatrix}$
- Weyl transformation to **Einstein frame**: $g_J \rightarrow g_E \equiv \Omega^2(\varphi) g_J$
- Canonical normalization: $\chi = \int_0^\varphi d\varphi \sqrt{\frac{3}{2} \frac{(\Omega^2,\varphi)^2}{\Omega^4} + \frac{1}{\Omega^2}}$

$$\mathcal{L}_E = \sqrt{-g_E} \left[\frac{M_P^2}{2} R_E + \frac{1}{2} g_E^{\mu\nu} (\partial_\mu \chi)(\partial_\nu \chi) - V_E(\chi) + \dots \right]$$

- Minimal coupling to gravity but modified potential
- Potential flat for large field values \rightsquigarrow slow-roll

Higgs Inflation



Bezrukov & Shaposhnikov, PLB 659 (2008)

- Minimal coupling to gravity but modified potential
- Potential flat for large field values \leadsto slow-roll

Higgs Inflation in a Dark Sector

$$\mathcal{L}_J = \sqrt{-g_J} \left[\frac{M_P^2}{2} \Omega^2(\phi) R_J + g_J^{\mu\nu} (D_\mu \phi)^\dagger (D_\nu \phi) - \frac{1}{4} g_J^{\mu\rho} g_J^{\nu\sigma} X_{\mu\nu} X_{\rho\sigma} - V_J(\phi) \right]$$

$$V_J(\phi) = -\mu^2 \phi^\dagger \phi + \lambda (\phi^\dagger \phi)^2 = \lambda \left(\phi^\dagger \phi - \frac{v^2}{2} \right)^2 + \text{const.}$$

- $U(1)_X$ gauge symmetry with gauge coupling g
- SM singlet scalar ϕ with $U(1)_X$ charge
 - Spontaneously breaks $U(1)_X \rightsquigarrow \text{FOPT} \rightsquigarrow \text{GW}$
 - Non-minimal coupling to gravity \rightsquigarrow Higgs-inflation-like inflation

Higgs Inflation in a Dark Sector

$$\mathcal{L}_J = \sqrt{-g_J} \left[\frac{M_P^2}{2} \Omega^2(\phi) R_J + g_J^{\mu\nu} (D_\mu \phi)^\dagger (D_\nu \phi) - \frac{1}{4} g_J^{\mu\rho} g_J^{\nu\sigma} X_{\mu\nu} X_{\rho\sigma} - V_J(\phi) \right]$$

$$V_J(\phi) = -\mu^2 \phi^\dagger \phi + \lambda (\phi^\dagger \phi)^2 = \lambda \left(\phi^\dagger \phi - \frac{v^2}{2} \right)^2 + \text{const.}$$

- $U(1)_X$ gauge symmetry with gauge coupling g
- SM singlet scalar ϕ with $U(1)_X$ charge
 - Spontaneously breaks $U(1)_X \rightsquigarrow \text{FOPT} \rightsquigarrow \text{GW}$
 - Non-minimal coupling to gravity \rightsquigarrow Higgs-inflation-like inflation
- Optional: fermions with mass from Yukawa coupling
- Weak coupling to SM \rightsquigarrow reheating

$$\mathcal{L}_\psi = \sqrt{-g_J} \sum_{i=1}^{n_\psi} \left[\bar{\psi}_{Li} i \not{D} \psi_{Li} + \bar{\psi}_{Ri} i \not{D} \psi_{Ri} - \frac{y}{2} (\bar{\psi}_{Li} \phi \psi_{Li}^c + \bar{\psi}_{Ri} \phi \psi_{Ri}^c + \text{h.c.}) \right]$$

Gravitational Waves from Phase Transition

Scalar potential with **1-loop** and **finite-temperature** corrections

- ↪ Potential **barrier** around $T \sim v$
- ↪ FOPT possible

Gravitational Waves from Phase Transition

Scalar potential with **1-loop** and **finite-temperature** corrections

- ↪ Potential **barrier** around $T \sim v$
- ↪ FOPT possible

GW spectrum determined by

- Nucleation temperature T_n
- $\alpha \leftrightarrow$ strength of PT
- $\beta \leftrightarrow$ duration

Calculated with help from CosmoTransitions

Wainwright, Comput. Phys. Commun. 183 (2012)

Effective Scalar Potential at Finite Temperature

$$V(h_c, T) = V_{\text{tree}}(h_c) + V_{\text{1-loop}}(h_c) + V_{\text{th}}(h_c, T)$$

$$V_{\text{tree}}(h_c) = -\frac{\mu^2}{2} h_c^2 + \frac{\lambda}{4} h_c^4$$

$$V_{\text{1-loop}}(h_c) = \sum_{i=h,\chi,g,f} \frac{n_i}{64\pi^2} m_i^4(h_c) \left[\ln \frac{|m_i^2(h_c)|}{v^2} - C_i \right]$$

$$V_{\text{th}}(h_c, T) = \sum_{i=h,\chi,g} \frac{n_i}{2\pi^2} T^4 \operatorname{Re} J_b\left(\frac{m_i^2(h_c)}{T^2}\right) + \frac{n_f}{2\pi^2} T^4 J_f\left(\frac{m_f^2(h_c)}{T^2}\right)$$

$$m_h^2(h_c) = -\mu^2 + 3\lambda h_c^2$$

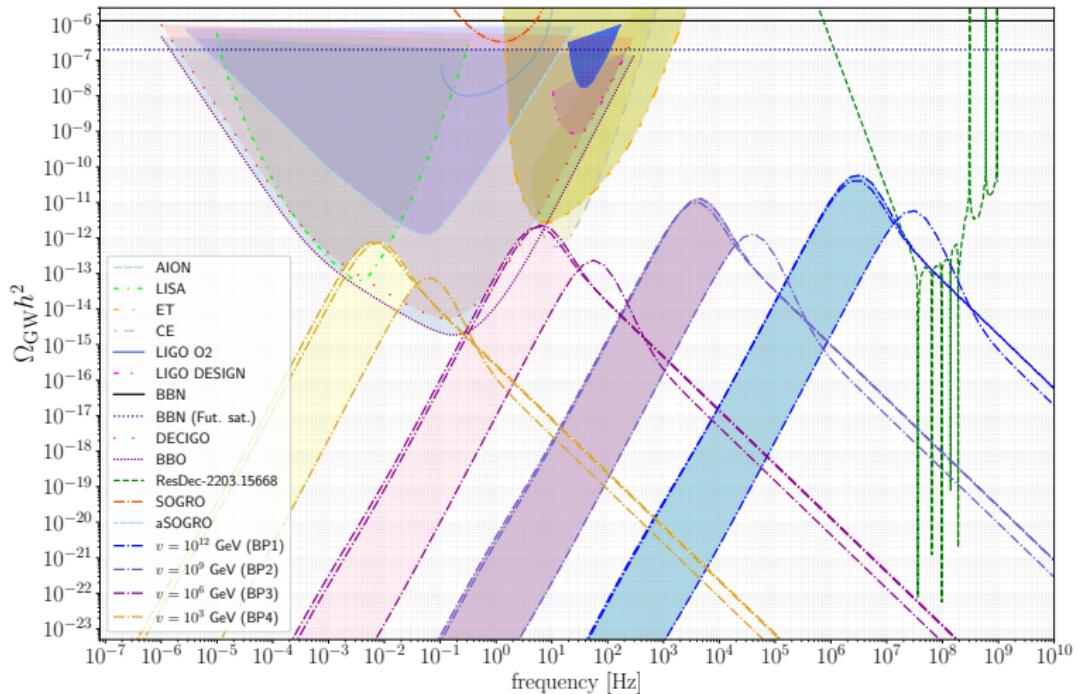
$$m_\chi^2(h_c) = -\mu^2 + \lambda h_c^2$$

$$m_g^2(h_c) = g^2 h_c^2$$

$$m_f^2(h_c) = \frac{y^2}{2} h_c^2$$

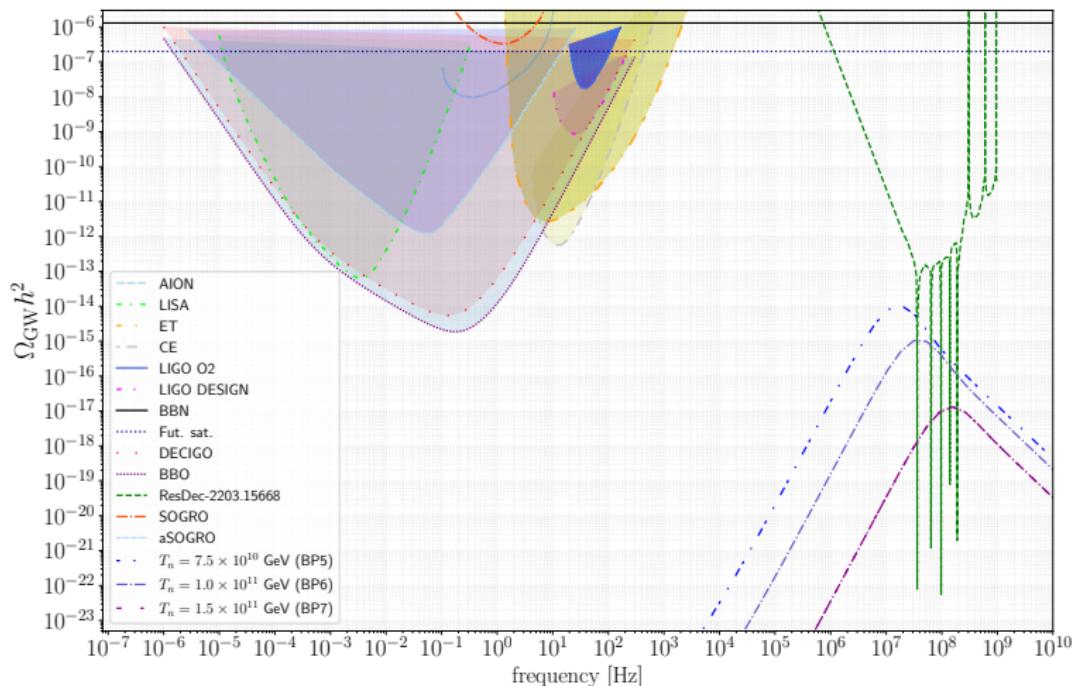
Gravitational Wave Spectrum I

No fermions, $g = 0.95$, $\lambda = 10^{-3}$
Bubble wall velocity $v_w = 1, v_{\text{det}}, 0.1$

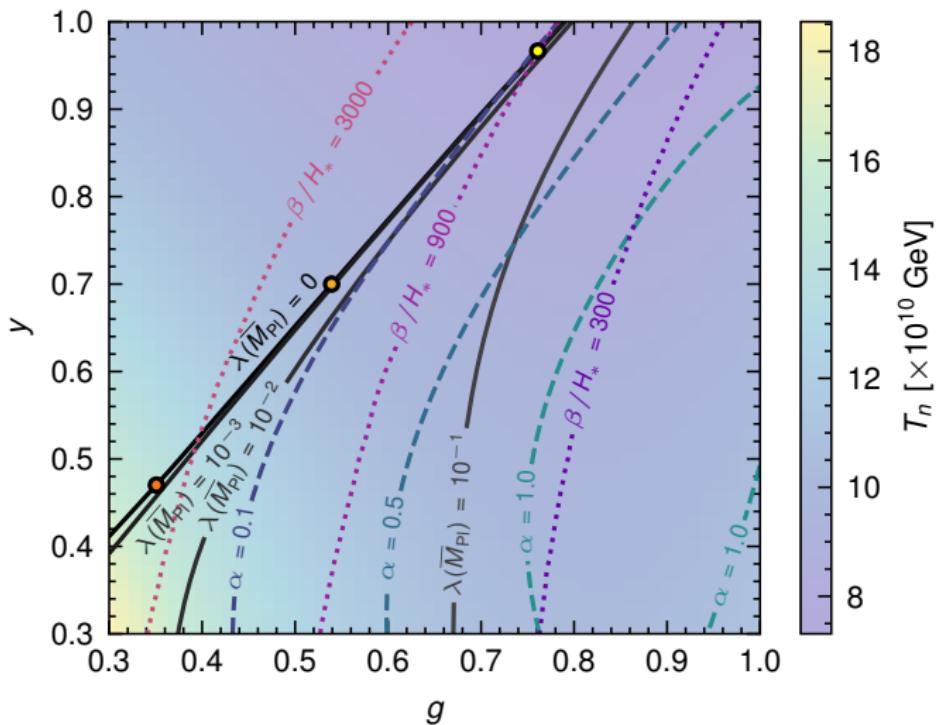


Gravitational Wave Spectrum II

1 fermion pair, $\nu_w = \nu_{\text{det}}$, $\nu = 10^{12} \text{ GeV}$, $\lambda = 10^{-3}$
 $(g, y) \simeq (0.76, 0.97), (0.54, 0.70), (0.35, 0.47)$



Parameter Space with 1 Fermion Pair



Inflationary Observables

CMB measurements (pre-ACT)

Scalar power spectrum amplitude $A_s = (2.098 \pm 0.023) \times 10^{-9}$

Scalar spectral index $n_s = 0.9649 \pm 0.0042$

Tensor-to-scalar power ratio $r < 0.036$ (95% CL)

Planck, A&A 641 (2020); BICEP/Keck, PRL 127 (2021)

Inflationary Observables

CMB measurements (pre-ACT)

Scalar power spectrum amplitude $A_s = (2.098 \pm 0.023) \times 10^{-9}$

Scalar spectral index $n_s = 0.9649 \pm 0.0042$

Tensor-to-scalar power ratio $r < 0.036$ (95% CL)

Planck, A&A 641 (2020); BICEP/Keck, PRL 127 (2021)

Model results (60 e-foldings)

- $n_s \simeq 0.965$
- $r \simeq 0.003$
- $A_s \simeq 5.1 \frac{\lambda(M_P)}{\xi^2} \quad \Rightarrow \quad \xi \simeq 5 \times 10^4 \sqrt{\lambda(M_P)}$

Inflationary Observables

Model results (60 e-foldings)

- $n_s \simeq 0.965$
- $r \simeq 0.003$
- $A_s \simeq 5.1 \frac{\lambda(M_P)}{\xi^2} \quad \Rightarrow \quad \xi \simeq 5 \times 10^4 \sqrt{\lambda(M_P)}$

	$g(M_P)$	$\lambda(M_P)$	$y(M_P)$	n_ψ	ξ
BP1	0.98	0.37	-	0	3×10^4
BP2	0.99	0.57	-	0	4×10^4
BP3	1.00	0.92	-	0	5×10^4
BP4	1.02	2.22	-	0	7×10^4
BP5	0.79	$\simeq 0$	1.10	1	$\mathcal{O}(1)$
BP6	0.55	$\simeq 0$	0.74	1	$\mathcal{O}(1)$
BP7	0.35	$\simeq 0$	0.48	1	$\mathcal{O}(1)$

Inflationary Observables

Model results (60 e-foldings)

- $n_s \simeq 0.965$
- $r \simeq 0.003$
- $A_s \simeq 5.1 \frac{\lambda(M_P)}{\xi^2} \quad \Rightarrow \quad \xi \simeq 5 \times 10^4 \sqrt{\lambda(M_P)}$

	$g(M_P)$	$\lambda(M_P)$	$y(M_P)$	n_ψ	ξ
BP1	0.98	0.37	-	0	3×10^4
BP2	0.99	0.57	-	0	4×10^4
BP3	1.00	0.92	-	0	5×10^4
BP4	1.02	2.22	-	0	7×10^4
BP5	0.79	$\simeq 0$	1.10	1	$\mathcal{O}(1)$
BP6	0.55	$\simeq 0$	0.74	1	$\mathcal{O}(1)$
BP7	0.35	$\simeq 0$	0.48	1	$\mathcal{O}(1)$

- With fermions: renormalization group running allows $\xi \sim 1$
- Analogous to **Critical Higgs Inflation**

Hamada et al., PRL 112 (2014), PRD 91 (2015)

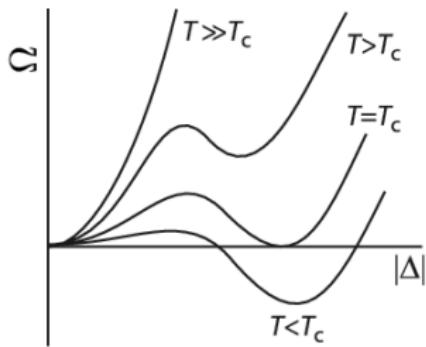
Conclusions and Outlook

- Unified framework for inflation and first-order phase transition
- Dark $U(1)_X$ and non-minimal coupling to gravity
 \rightsquigarrow gravitational waves and inflation from same scalar
- Peak GW frequencies between 10^{-2} Hz and 10^8 Hz

Conclusions and Outlook

- Unified framework for inflation and first-order phase transition
- Dark $U(1)_X$ and non-minimal coupling to gravity
 \rightsquigarrow gravitational waves and inflation from same scalar
- Peak GW frequencies between 10^{-2} Hz and 10^8 Hz
- Future improvements
 - Calculation of bubble wall velocity using WallGo
Ekstedt et al., JHEP 04 (2025)
 - Inflaton coupling to SM \rightsquigarrow reheating
 - Dark Matter
 - Different gauge groups

First-Order Phase Transitions



Kinnunen et al., Rep. Prog. Phys. 81 (2018)

- ① High temperature: potential minimum at $\phi = 0$
- ② $T < T_c$: deeper minimum at $\phi = v$, separated by **barrier**
- ③ Tunneling \rightsquigarrow bubbles of true vacuum
- ④ **GW** from expanding bubbles and surrounding plasma

Parameters Determining the GW Signal

- T_n : temperature at which probability of nucleating 1 bubble per horizon volume is ~ 1
- $\alpha = \frac{1}{\rho_{\text{rad}}} \left[\Delta V(h_c, T) - T \frac{dV(h_c, T)}{dT} \right] \Big|_{T=T_n}$
- $\frac{\beta}{H(T_n)} = T_n \left. \frac{d}{dT} \frac{S_3}{T} \right|_{T_n} \simeq \left(\frac{T_n}{T-T_n} \right) \left(\frac{S_3(T)}{T} - \frac{S_3(T_n)}{T_n} \right)$
(T : arbitrary reference temperature not too far from T_n)
- Expressions modified for large supercooling

SM Extensions with First-Order Phase Transitions

- **SM: no FO electroweak PT for Higgs mass $\gtrsim 70$ GeV**
Kajantie et al., PRL **77** (1996); Karsch et al., NP Proc. Suppl. **53** (1997)
Csikor et al., PRL **82** (1999)
- **Most famous: SUSY with light stop**
Carena et al., PLB **380** (1996); Espinosa, NPB **475** (1996)
Delepine et al., PLB **386** (1996); Cline & Kainulainen, NPB **482** (1996); ...
- **Minimal from model building perspective: 2HDM**
Dorsch et al., JHEP **10** (2013); Basler et al., JHEP **02** (2017)
Andersen et al., PRL **121** (2018); ...
- **Minimal from EFT perspective: higher-dimensional operators**
Zhang, PRD **47** (1993); Grojean et al., PRD **71** (2005)
Bödeker et al., JHEP **02** (2005); Chala et al., JHEP **07** (2018); ...
- **Many recent works** Review: Roshan & White, arXiv:2401.04388