

PATHS THROUGH THE DARK: COMPARING APPROACHES FOR FOPTS

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In collaboration with Eric Madge, Maura E. Ramírez-Quezada and Pedro Schwaller

GW detection opens a window to study physics BSM

GWs down to the nHz regime

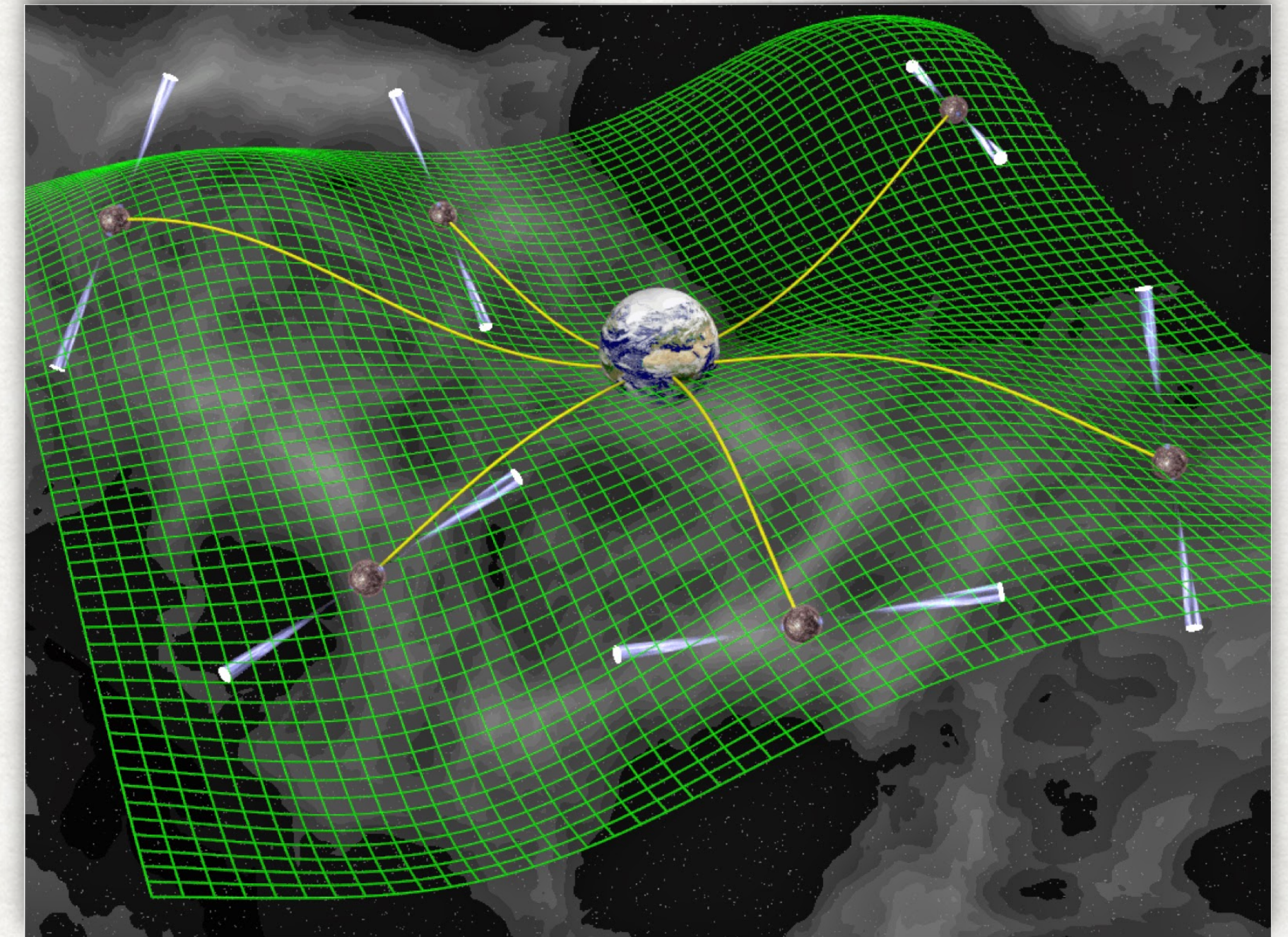
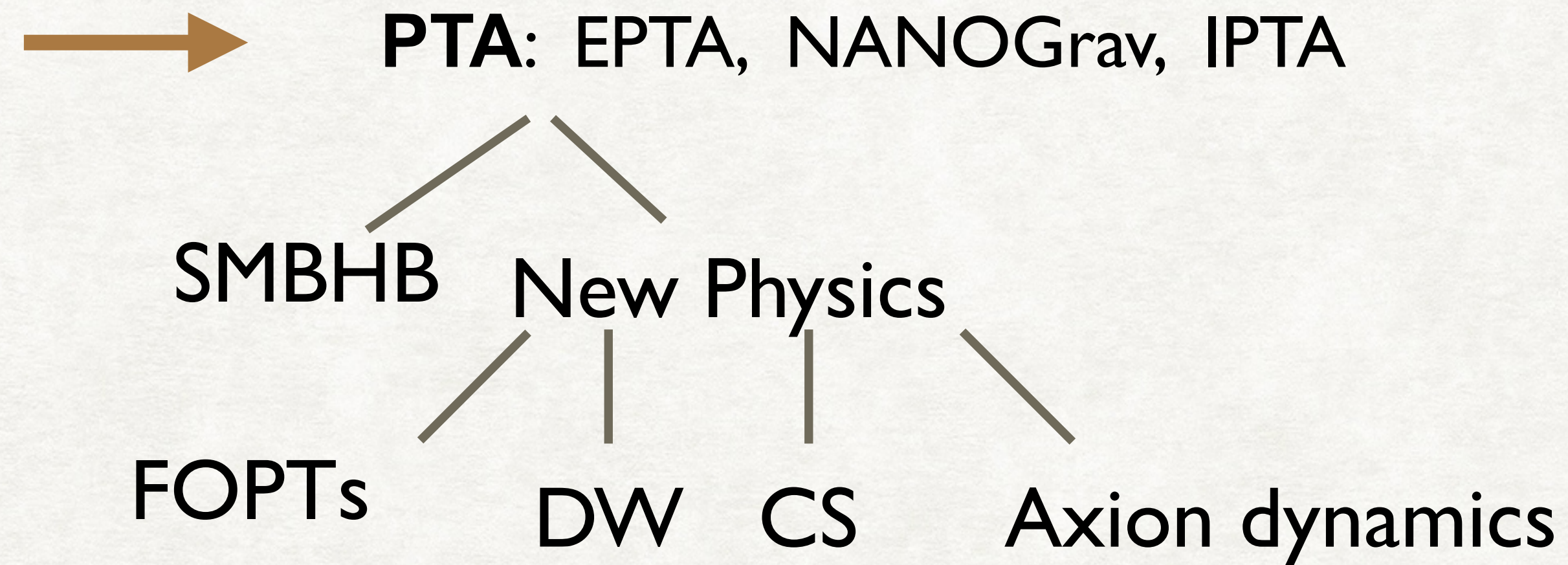


Image credit: David Champion

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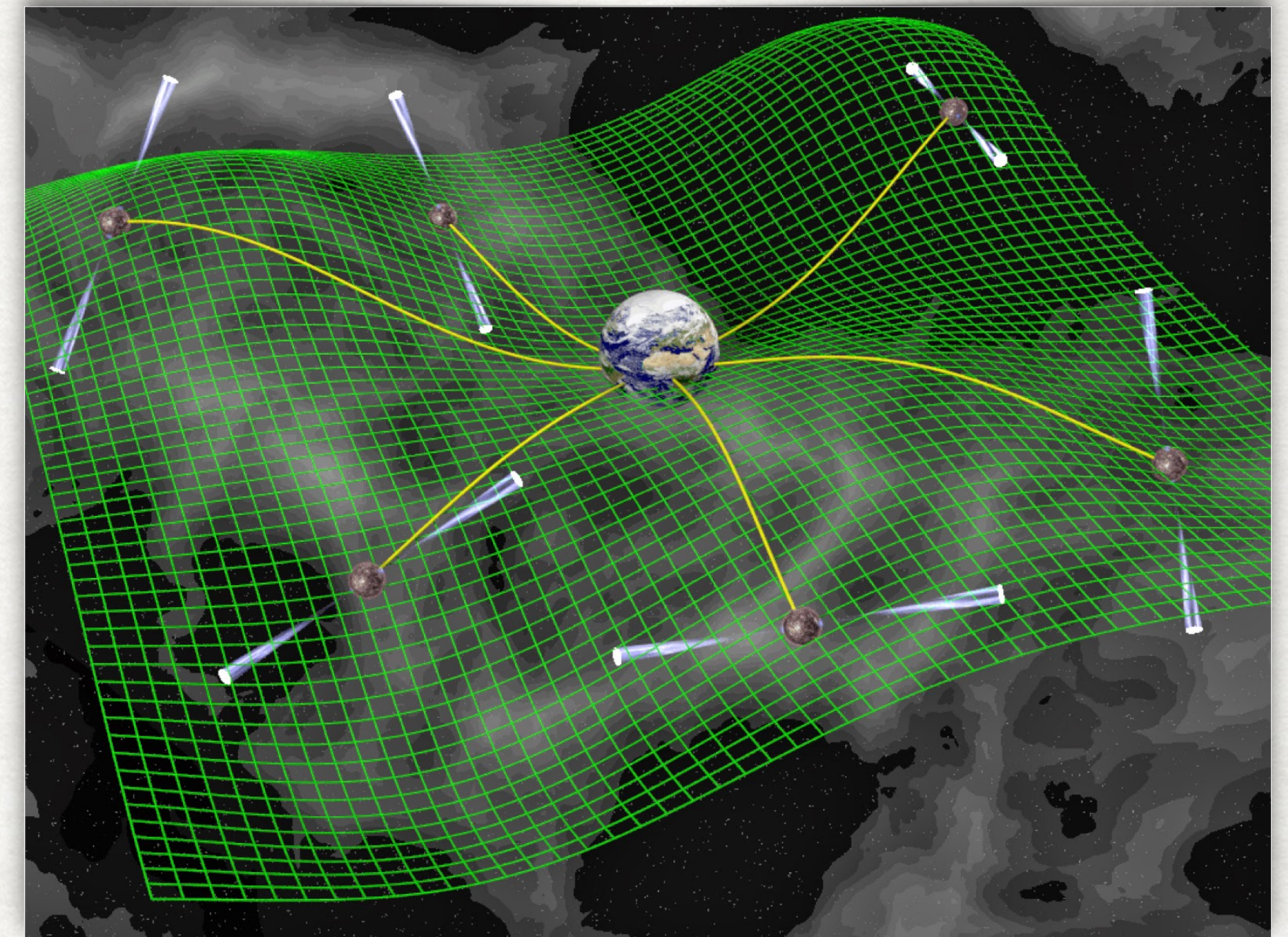
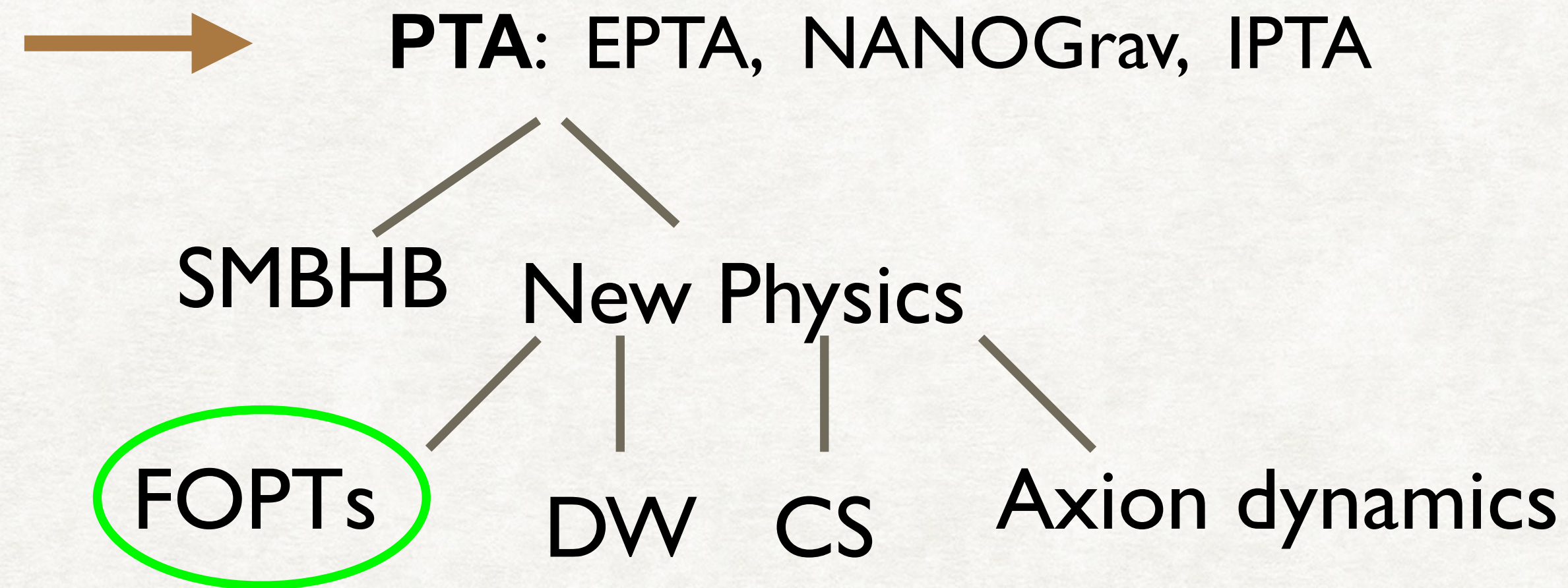
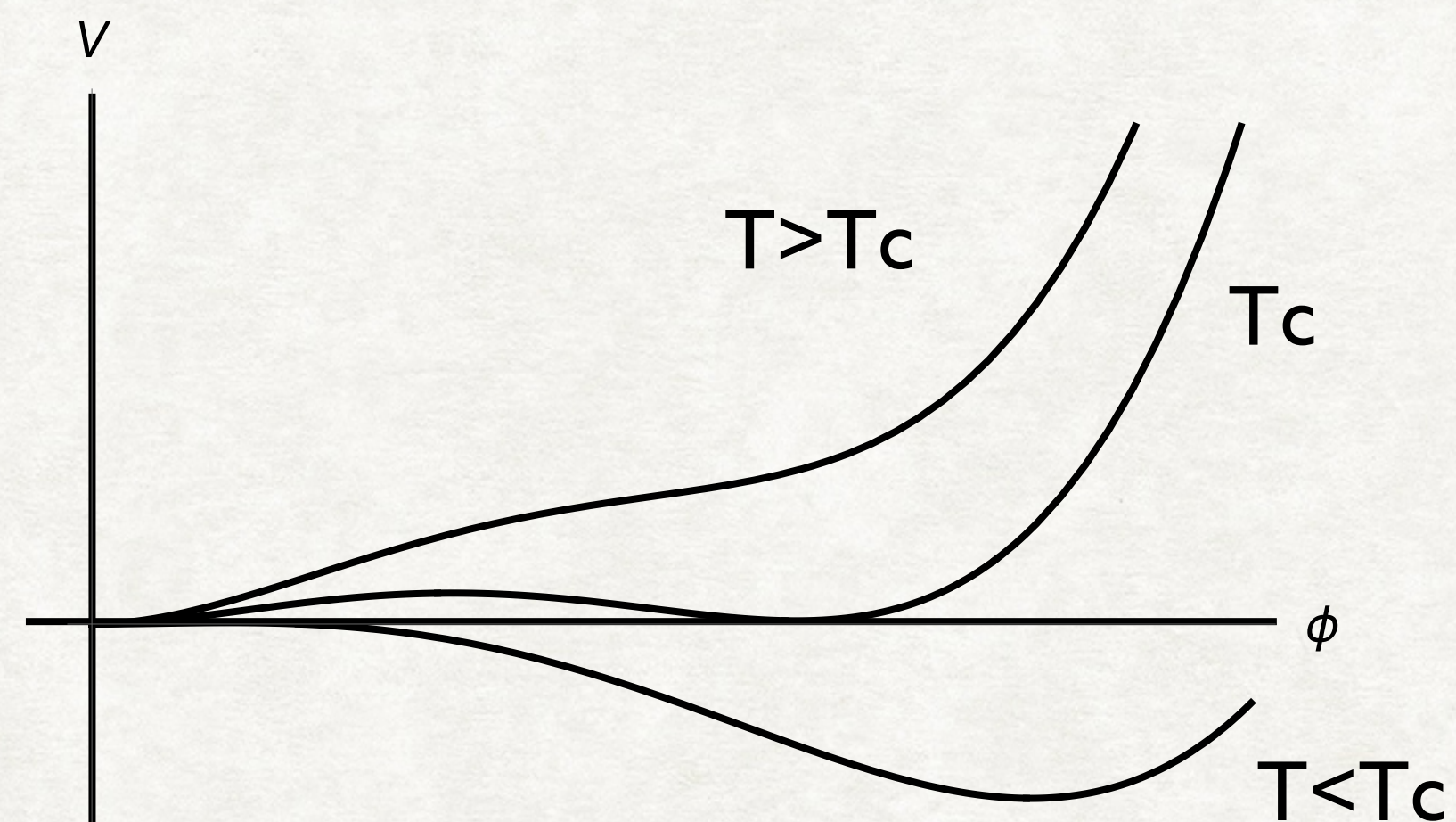


Image credit: David Champion



Supercooling: Delay of the FOPT below T_c

→ large release of latent heat α

PREVIOUS WORK

➔ Supercooled FOPTs

Model: classically scale invariance + radiative breaking of conformal symmetry: **Coleman-Weinberg (CW)**.

- Zero-T CW at 1-Loop

$$V_{CW} = \sum_i N_i \frac{m_i^4(\phi)}{64\pi^2} \left[\log \left(\frac{m_i^2(\phi)}{\mu_R^2} \right) - c_i \right]$$

- Finite Temperature:

$$V_1(\phi, T) = \frac{m^2(T)}{2} \phi^2 - \frac{\delta(T)}{3} \phi^3 + \frac{\lambda(T)}{4} \phi^4$$

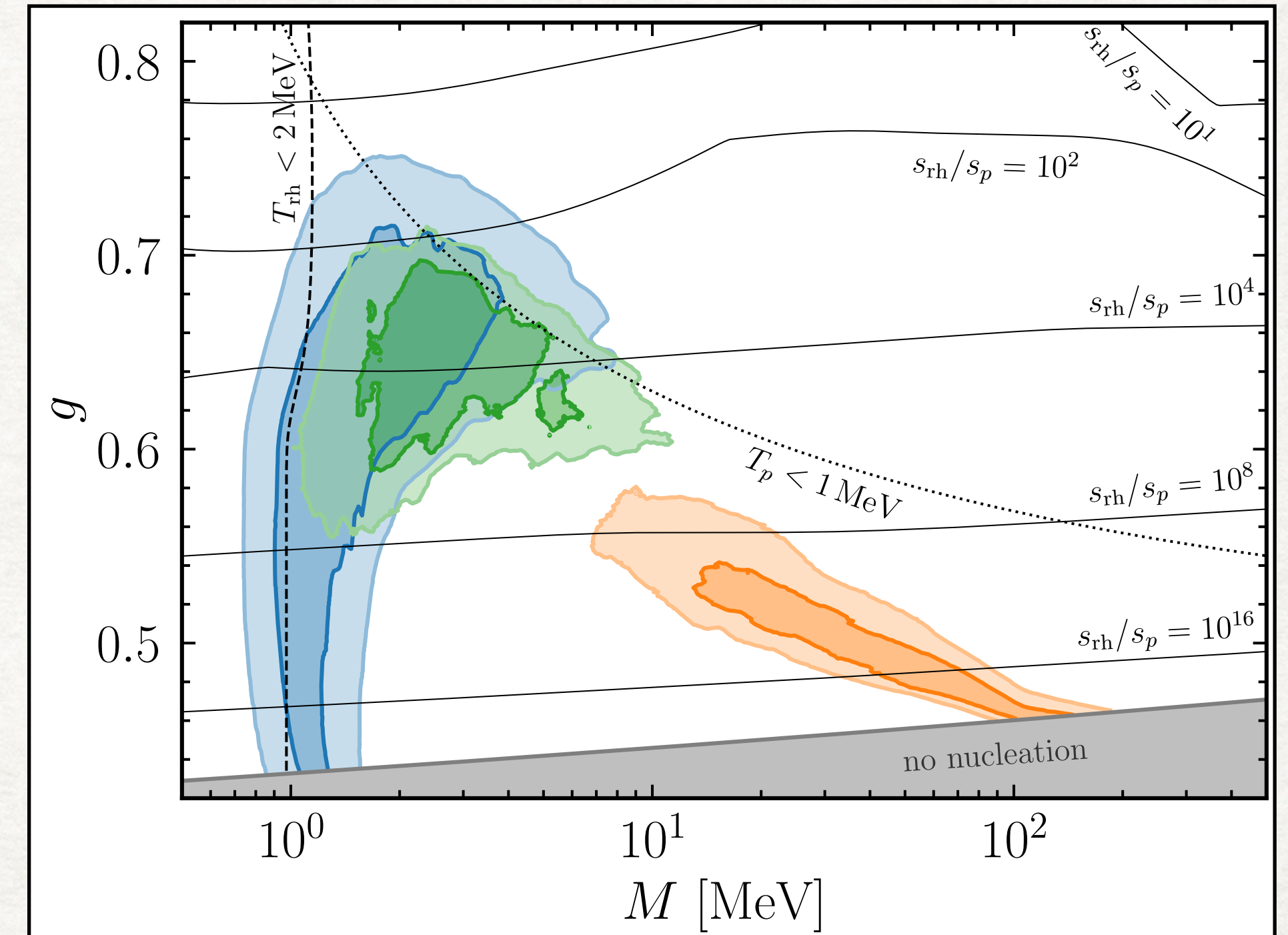
$\log(\phi)$ cancels

Neglect Daisy contributions !

$$\kappa \equiv \frac{\lambda(T)m^2(T)}{\delta^2(T)} = \frac{\log(T/M)}{6}$$

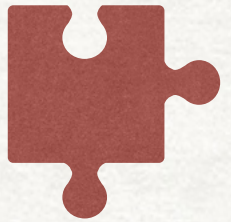
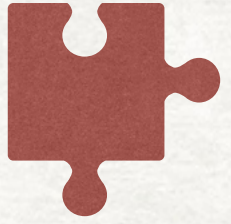
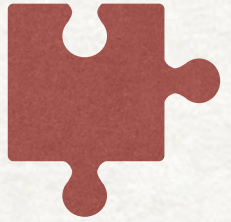
➔ **OPA**

For more details look <https://arxiv.org/abs/2212.08085> by N. Levi, T. Opferkuch, D. Redigo



From <https://arxiv.org/pdf/2306.14856> by E. Madge, E. Morgante, **CPI**, N. Ramblerg, W. Ratzinger, S. Schenk, P. Schwaller

What we wanna explore now?

-  OPA vs. HT expansion and full form of the 4D potential adding Daisy resummation.
-  Daisy resummation vs. Dimensional reduction (DR) EFTs
-  Effects of 2-Loop corrections using DR.

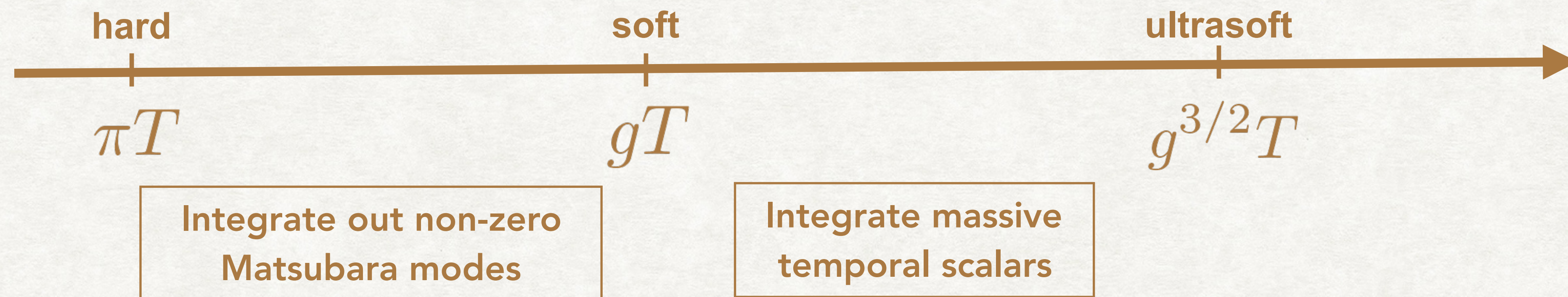
Perturbation theory breakdown $\longrightarrow$ IR divergences

DAISY RESUMATION

$$V_{\text{ring}}(\phi, T) = -\frac{T}{12\pi} \sum_i n_i \left[(m_i^2(\phi, T))^{3/2} - (m_i^2(\phi))^{3/2} \right]$$



DIMENSIONAL REDUCTION (DR)



SET UP

Dark photon: Extend SM with an additional U(1).

Supercooling

1-Loop potencial in the $\overline{MS}$ scheme:

$$V_{1\text{-loop}}(\phi) = -\frac{m^2}{2}\phi^2 + \frac{\lambda}{4}\phi^4 + \frac{3(g^2\phi^2)^2}{64\pi^2} \left[\log \frac{g^2\phi^2}{\mu^2} - \frac{5}{6} \right] + \frac{(3\lambda\phi^2 - m^2)^2}{64\pi^2} \left[\log \frac{|3\lambda\phi^2 - m^2|}{\mu^2} - \frac{3}{2} \right] + \frac{(\lambda\phi^2 - m^2)^2}{64\pi^2} \left[\log \frac{|\lambda\phi^2 - m^2|}{\mu^2} - \frac{3}{2} \right]$$

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Breaking parameter $\epsilon = m^2/v^2$

$$\lambda = \frac{g^4}{16\pi^2} \left(1 - 3 \log \frac{g^2 v^2}{\mu^2} \right) + \epsilon + \frac{\lambda}{16\pi^2} \left[3(3\lambda - \epsilon) \left(1 - \log \frac{|3\lambda - \epsilon|}{\mu^2/v^2} \right) + (\lambda - \epsilon) \left(1 - \log \frac{|\lambda - \epsilon|}{\mu^2/v^2} \right) \right]$$

$$\text{If } \epsilon \ll \frac{g^4}{16\pi^2} \quad \lambda \sim \frac{g^4}{16\pi^2}$$

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But DRalgo power counting

$$g^2 \sim \lambda \sim m^2$$

MATCHING 3D AND 4D THEORY

couplings

$$g_{3d}^2 = g^2 T - \frac{g^4 L_b T}{48\pi^2}$$

$$\lambda_{3d} = \lambda T + \frac{T}{16\pi^2} [g^4 (2 - 3L_b) + 6g^2 \lambda L_b - 10\lambda^2 L_b]$$

$$\lambda_A = 0 + \frac{g^4 T}{\pi^2}$$

$$\lambda_{A\Phi} = 2g^2 T + \frac{g^2 T}{24\pi^2} [24\lambda - g^2 (L_b - 4)]$$

masses

$$m_{3d}^2 = m^2 - \frac{T^2(3g^2 + 4\lambda)}{12} + \frac{L_b(3g^2 - 4\lambda)m^2}{16\pi^2} + \frac{g^4(8 + 216c_+ + 39L_b)T^2}{576\pi^2} + \frac{\lambda^2(12c_+ + 5L_b)T^2}{24\pi^2} - \frac{g^2\lambda(1 + 12c_+ + 3L_b)T^2}{24\pi^2} - \frac{8g_{3d}^4 - 16g_{3d}^2\lambda_{3d} + 16\lambda_{3d}^2 + \lambda_{A\Phi}^2}{32\pi^2} \log \frac{\mu_3}{\mu_R}$$

$$\mu_D^2 = \frac{g^2 T^2}{3} + \frac{g^2}{144\pi^2} [12\lambda T^2 + g^2(7 - L_b)T^2 - 36m^2]$$

constants

$$c_+ = \frac{1}{2} (\gamma_E - L_b - 12 \log A)$$

$$L_b = \log \frac{\mu_R^2 e^{2\gamma_E}}{16\pi^2 T^2}$$

DIFFERENT APPROACHES

$$\phi_{3d} \rightarrow \phi/\sqrt{T}$$

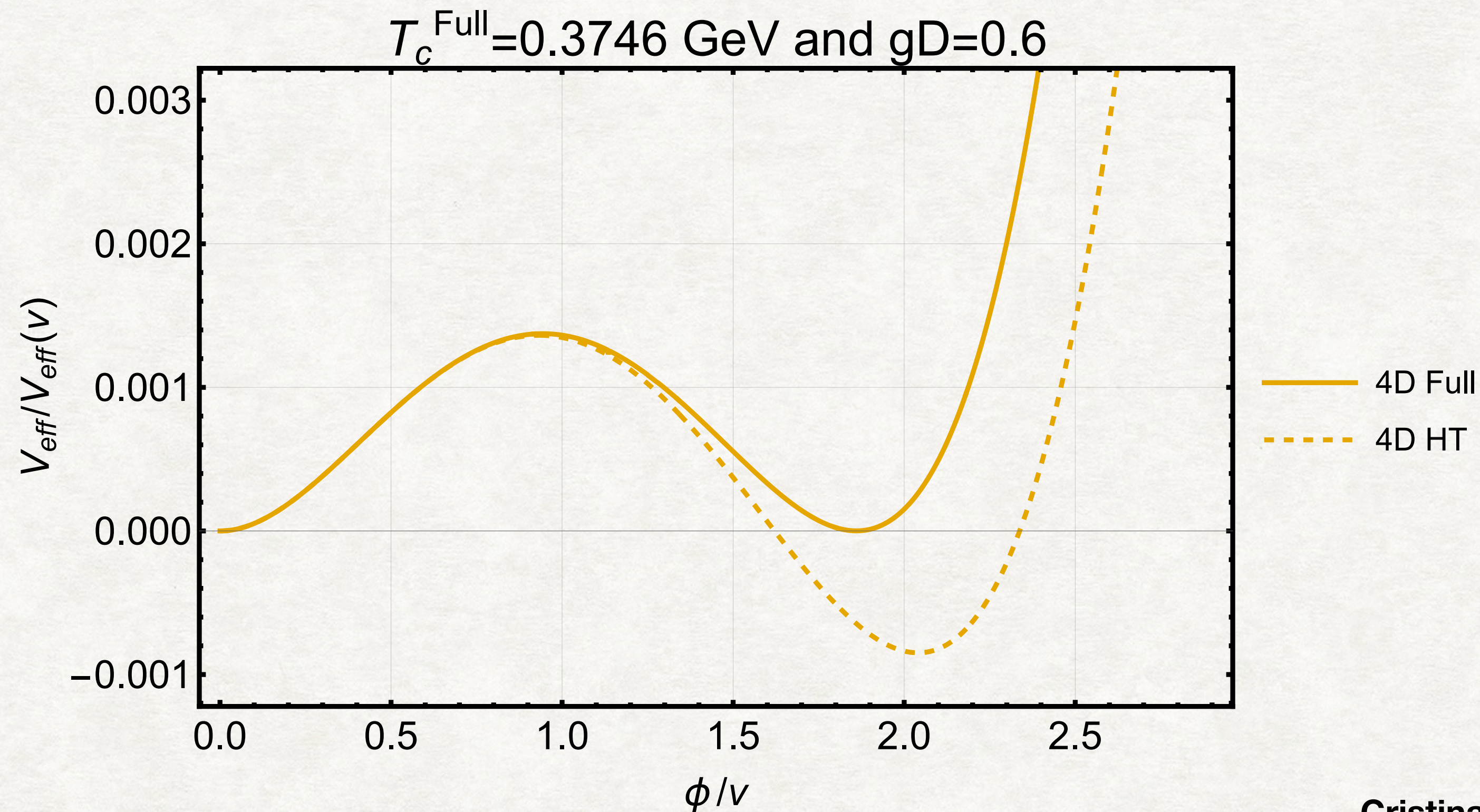
4D: $V_{eff}(\phi, T) = V_{tree}(\phi) + V_{CW}(\phi) + V_T(\phi, T) + V_{daisy}(\phi, T)$

Model parameters

$$\mu_R = v = 1$$

$$\lambda = 10^{-10}$$

$$m^2 = \lambda v$$



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3D soft:

NLO (1-Loop):

- masses at LO +
couplings NLO

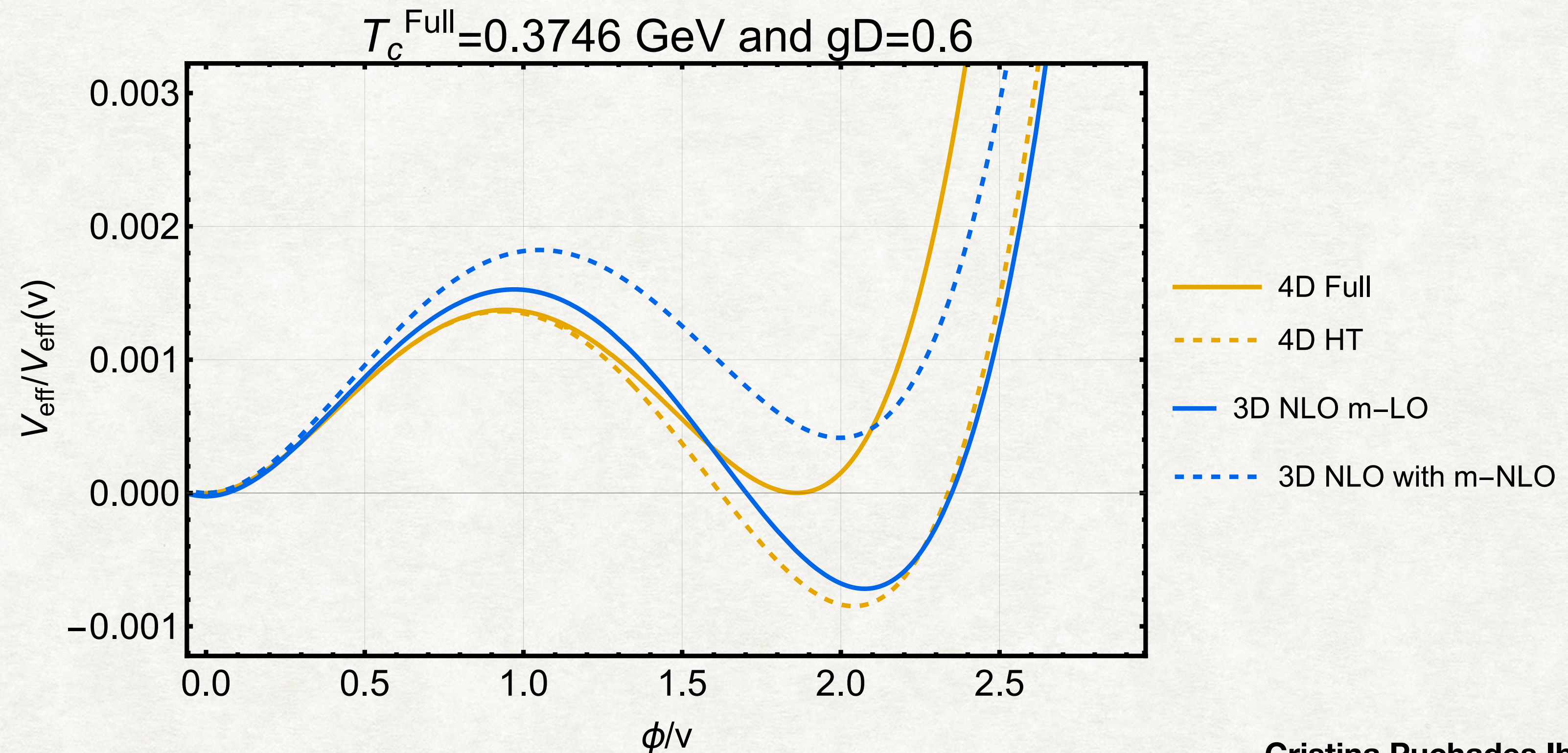
- masses at NLO with 3d
couplings NLO+couplings
NLO

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3D soft:

NLO (1-Loop):

NNLO (2-Loop):

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- masses at NLO with 3d
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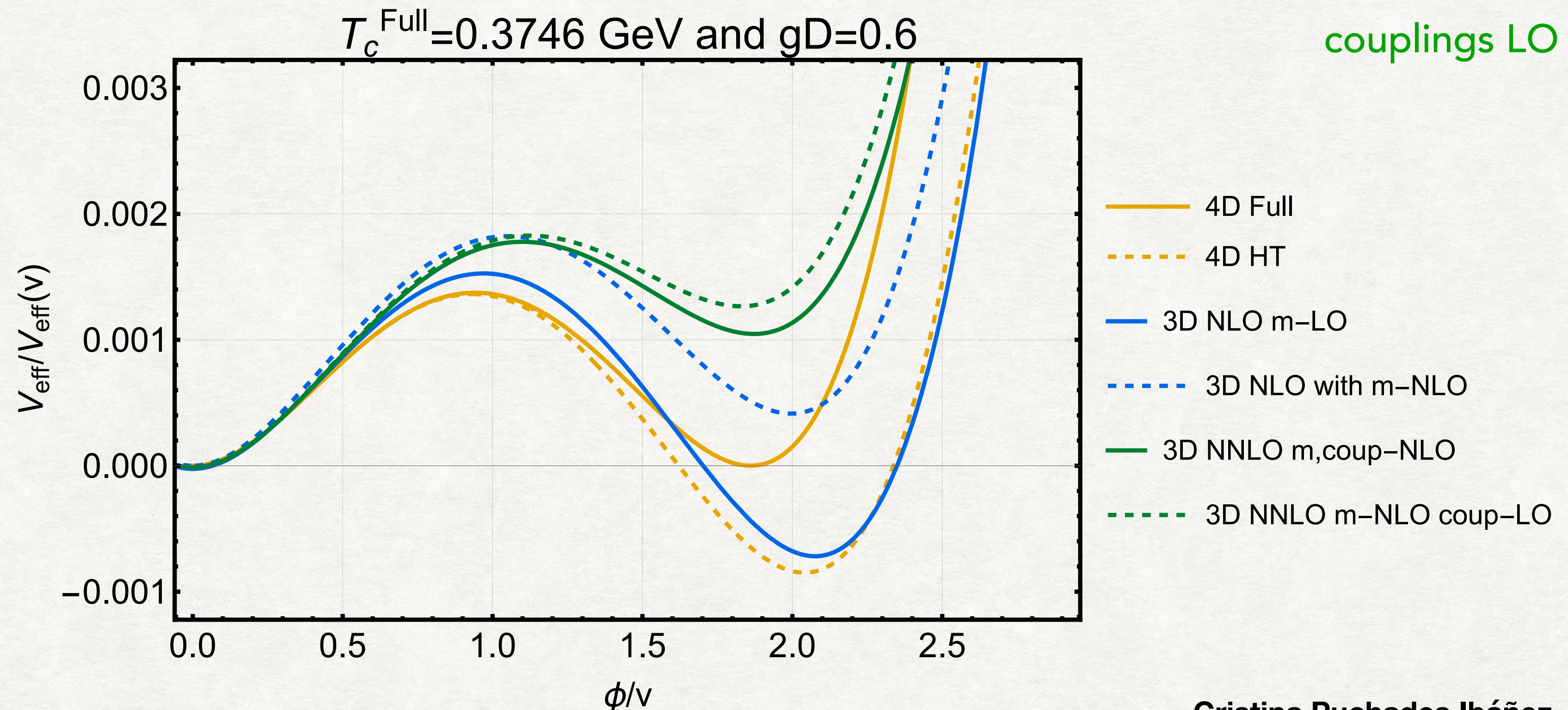
- V-LO masses NLO
with couplings at LO +
V-NLO+V-NNLO
masses LO and
couplings LO

Model parameters

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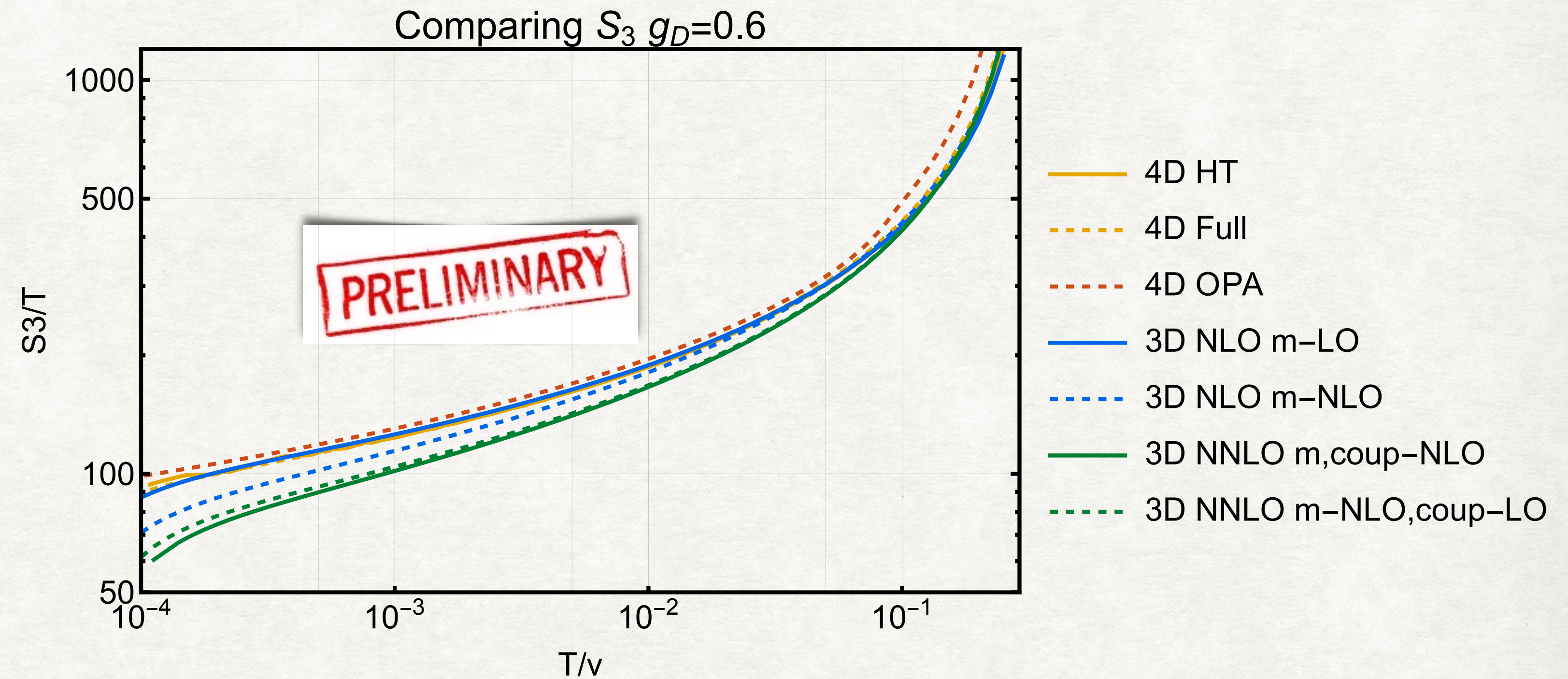
FOPT PARAMETERS

Espinosa:

$$\frac{S_3}{T} = \frac{32\pi\sqrt{2}}{3} \int_0^{\phi_0} d\phi \frac{[V_{\text{eff}}(\phi, T) - V_t(\phi)]^{3/2}}{\left(\frac{dV_t}{d\phi}\right)^2}$$

$$\alpha \simeq \frac{\Delta V_{CW}^{NLO}}{\rho_*} \quad \rho_* = \frac{\pi^2}{30} g_* T_*^4$$

$$\frac{\Gamma(T_n)}{H(T_n)^4} = 1 \quad \frac{\beta}{H_*} = T_* \frac{d}{dT} \left(\frac{S_3}{T} \right) \Big|_{T=T_*}$$



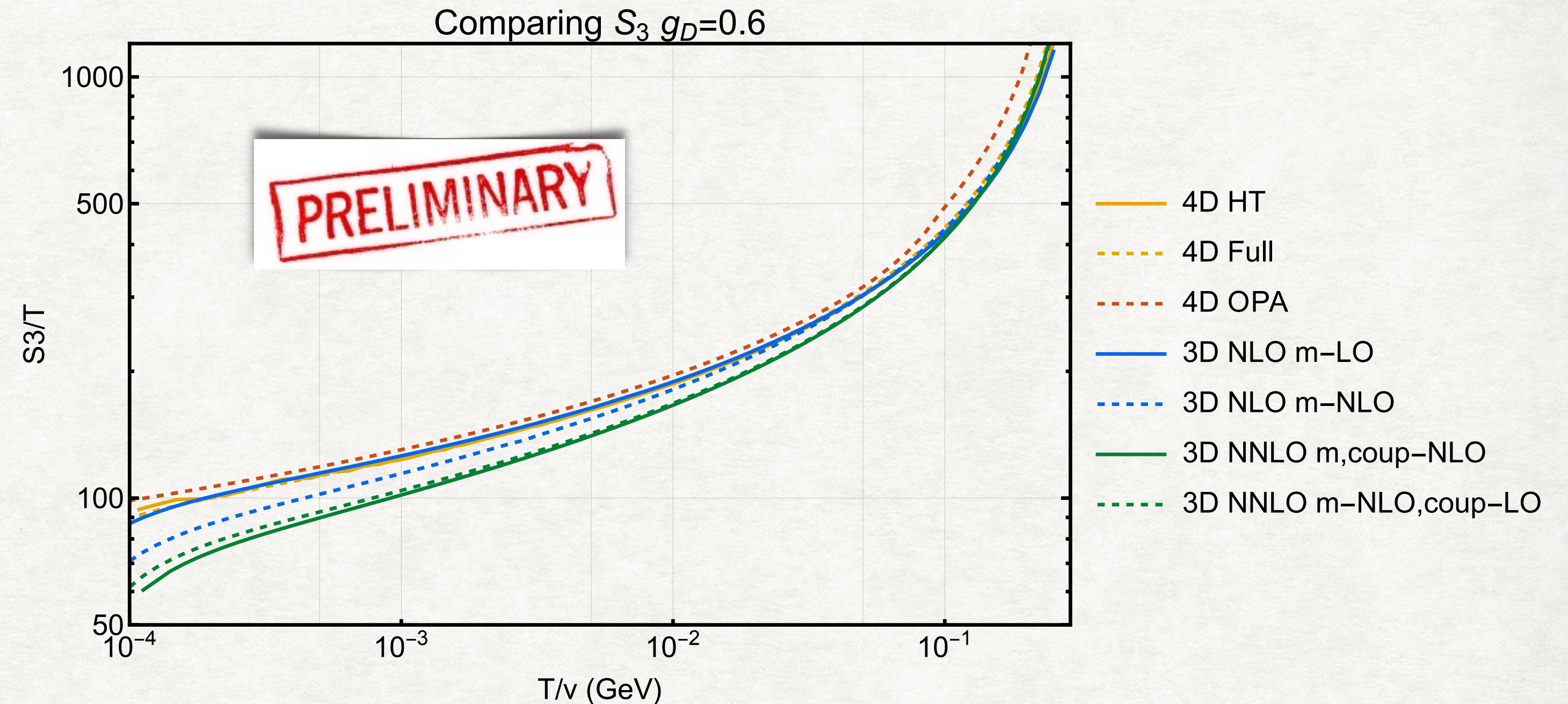
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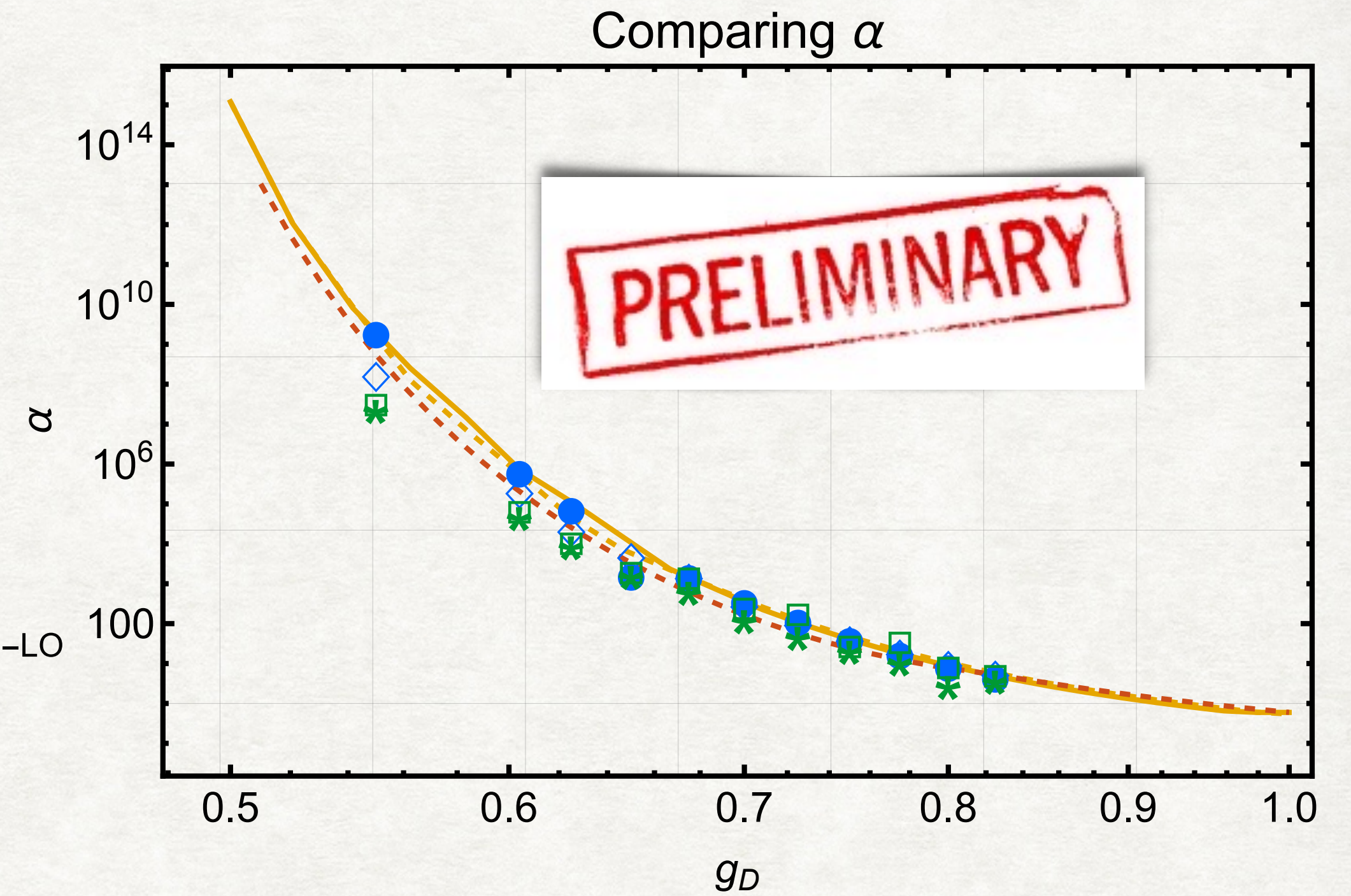
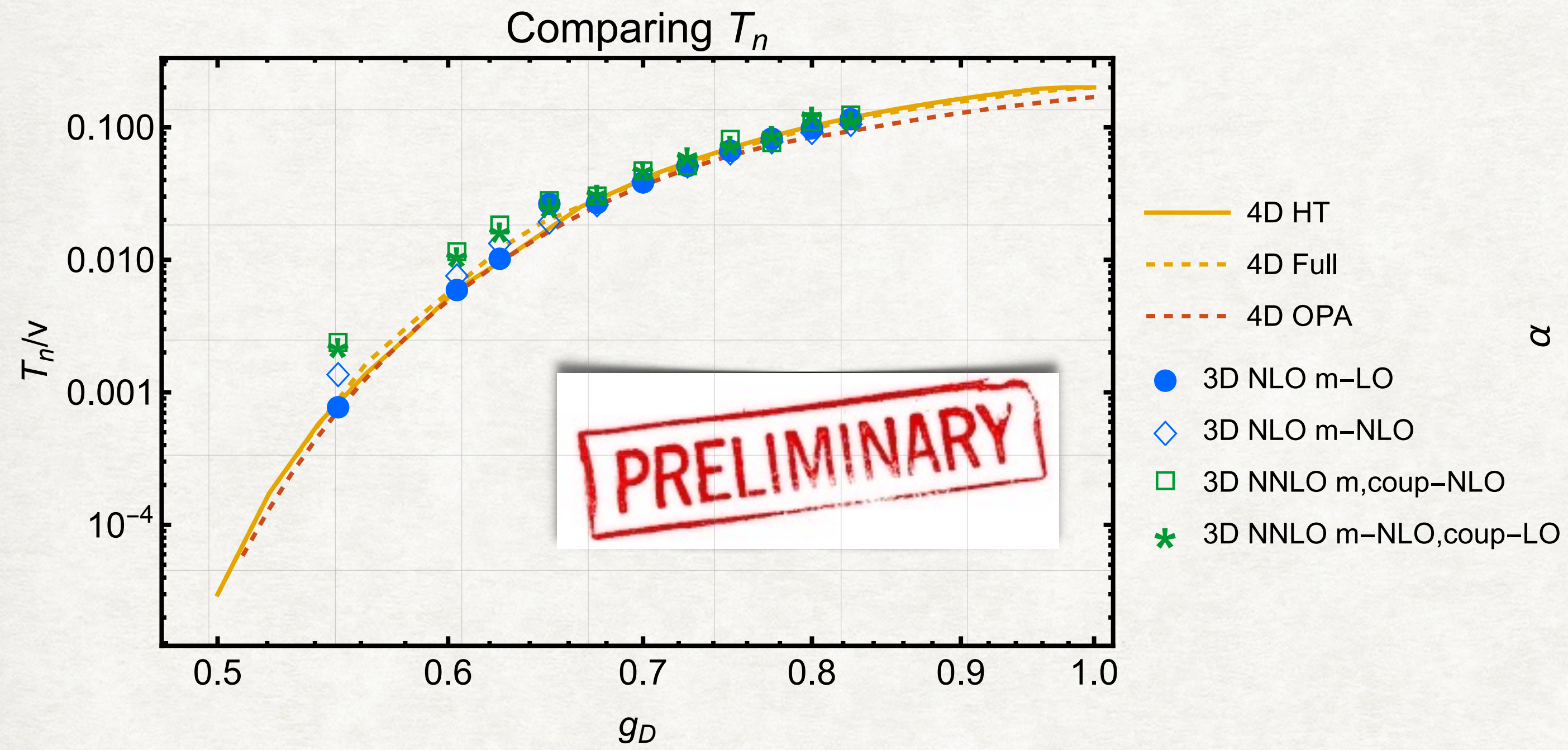
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For more details look <https://arxiv.org/pdf/2312.12413>

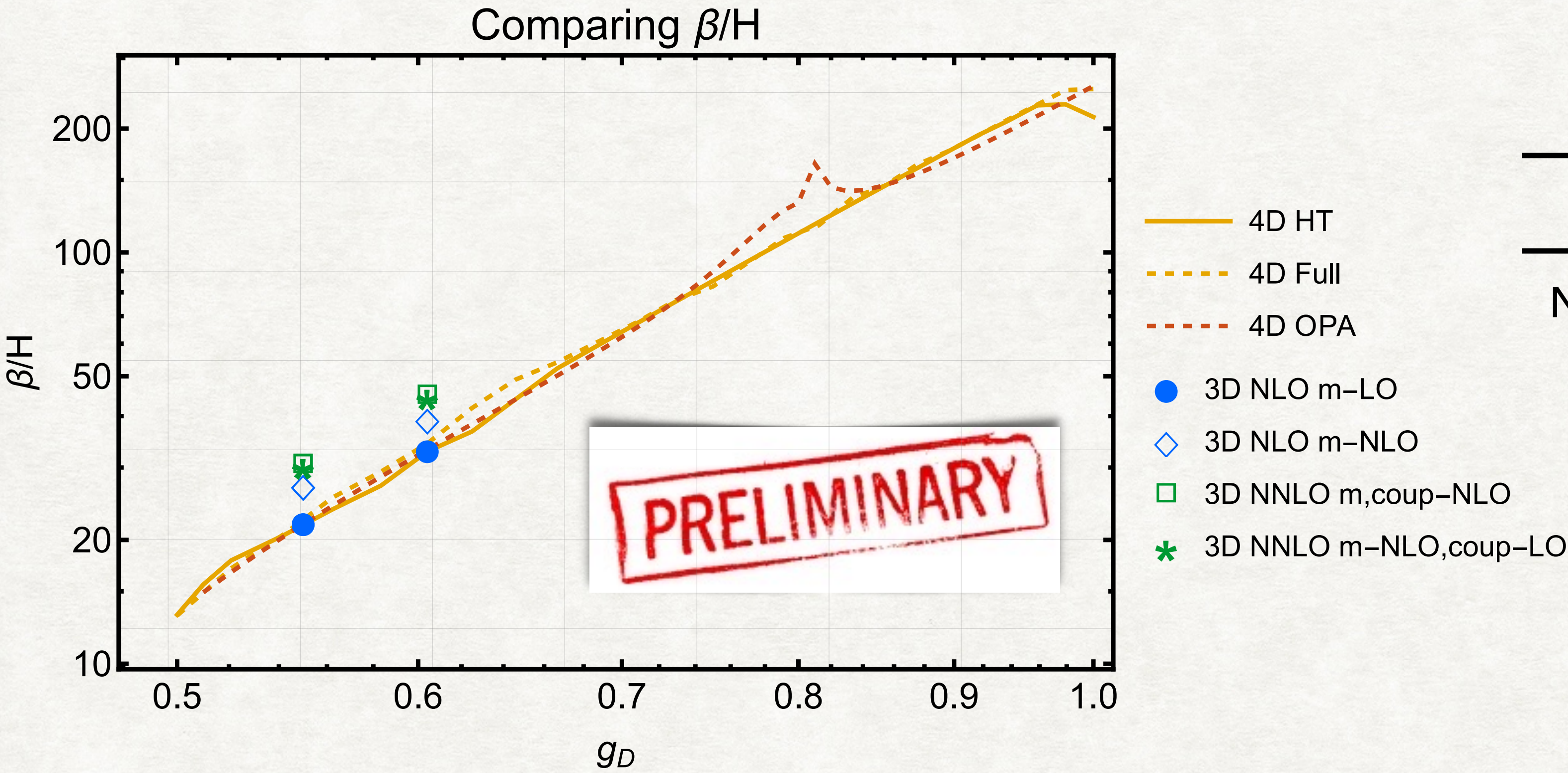
M. Kierkla, B. Siezewska, T. V. I.
Tenkanen and J. van de Vis

RESULTS



- Numeric problems.
- The LO with NLO masses and the NNLOs make a difference for small couplings.
- LO with LO masses is very close to the 4D theory: DR includes daisy

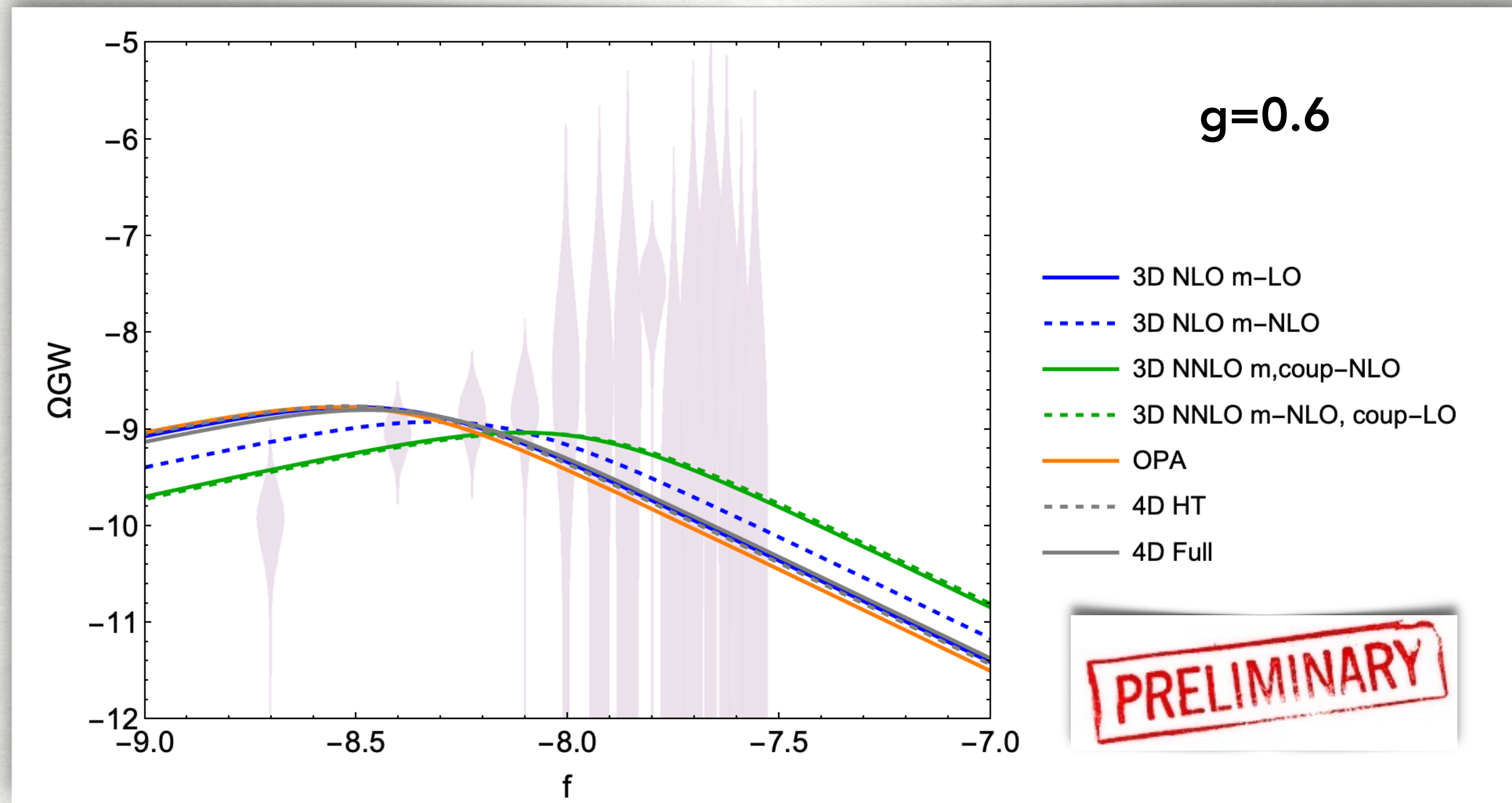
RESULTS



	$g = 0.6$	$g = 0.55$
NLO	$\sim 18\%$	$\sim 31\%$
NNLO	$\sim 32\%$	$\sim 45\%$

- “artificial” peak
- Due to numeric problems we have only 2 points.
- Smaller the coupling more impact have DR corrections

CONCLUSIONS AND IMPROVEMENTS



Improve numerics.

Explore the effect of the running.

Consistent power counting

Effect of V-NLO with LO masses negligible.

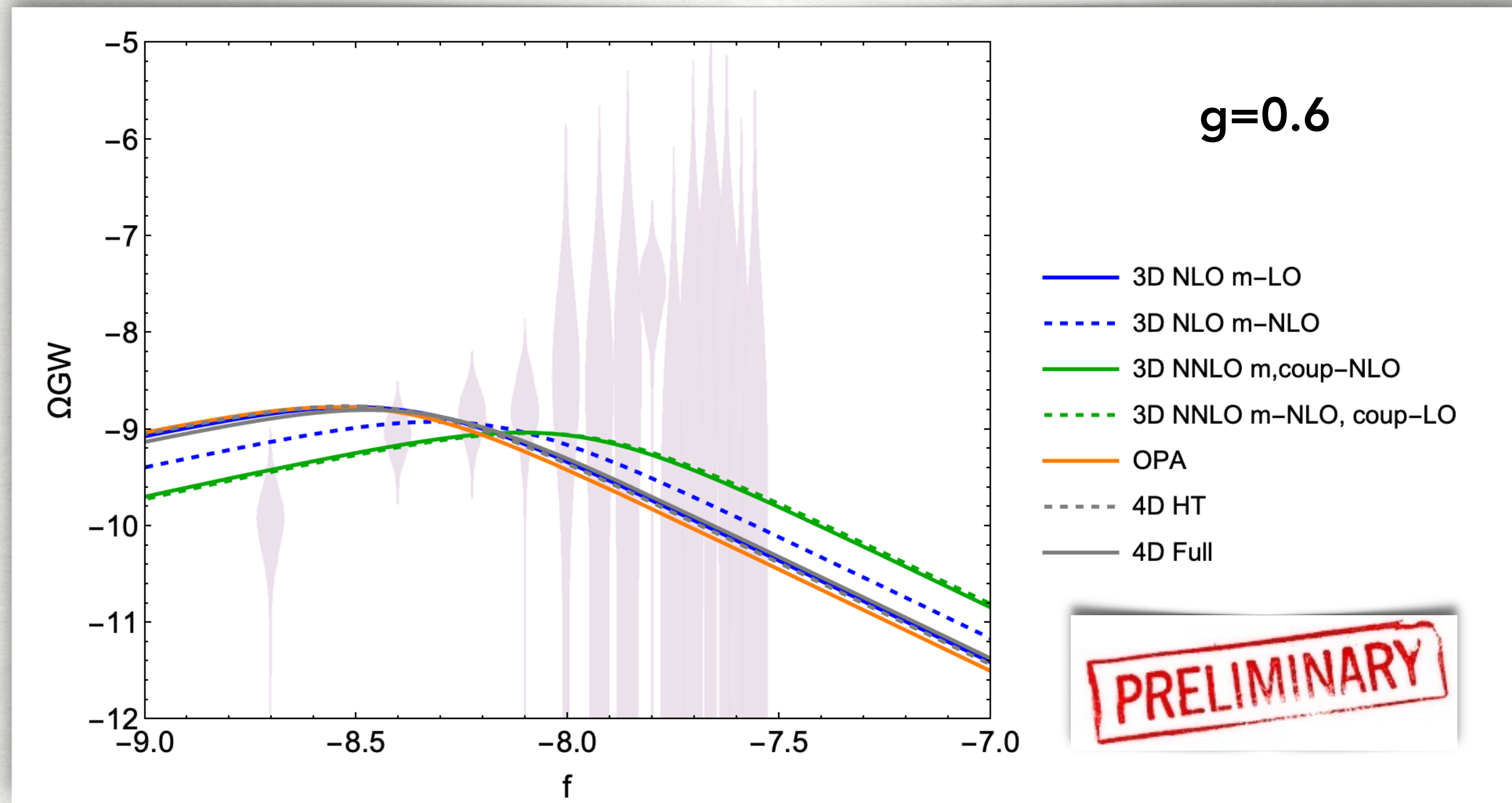
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GW spectrum in the relativistic limit from

<https://arxiv.org/pdf/2403.05615>

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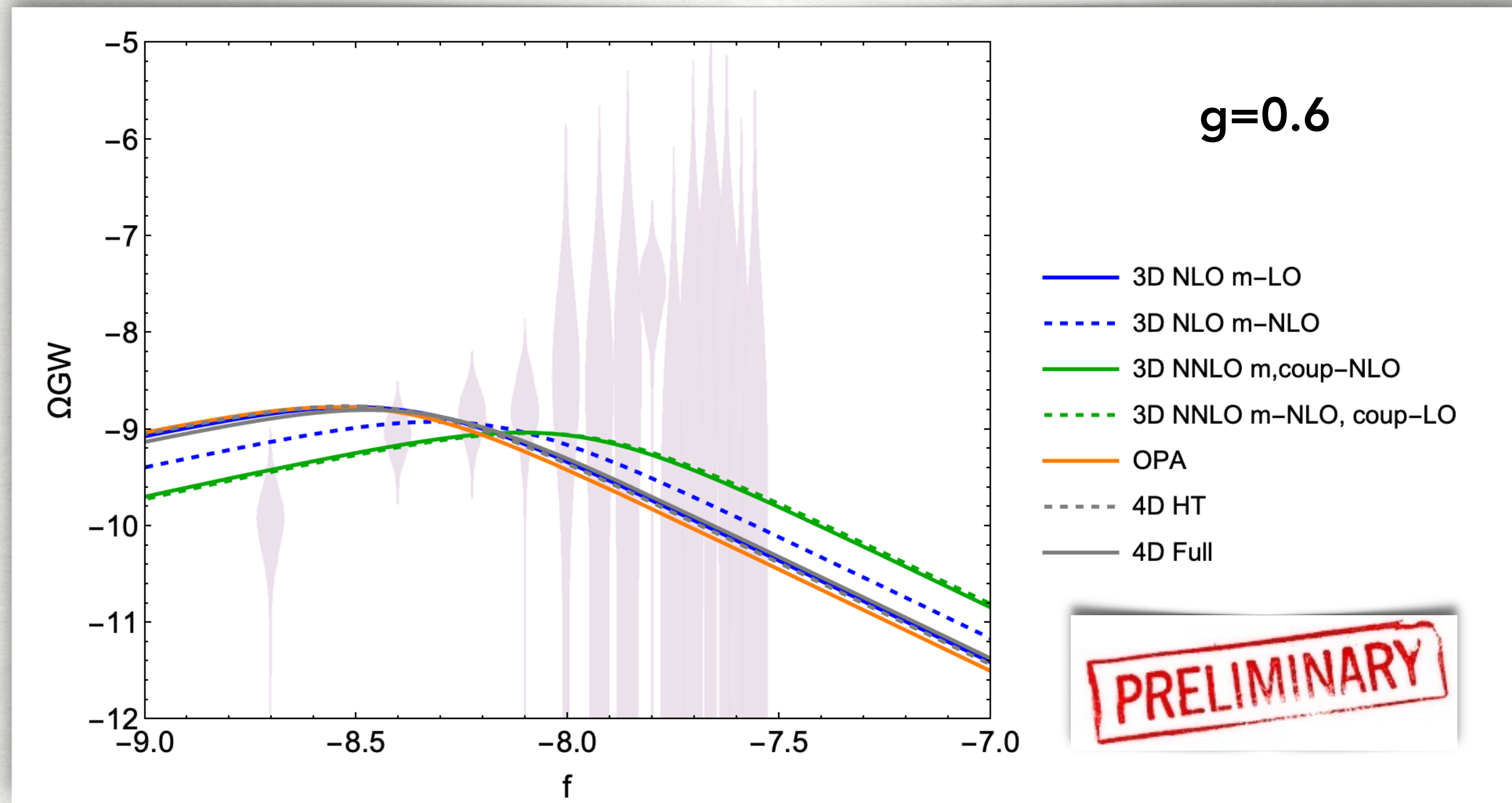
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